

Allright. There's some code here, and it does stuff.

```

rrEchelon.py  Takes a matrix and returns its reduced row echelon form.
getSyms.py   Uses a couple of generators to find the 48 symmetries of a
              cubic crystal. It dumps them out to S.pkl.
S.pkl        List of numpy matrices, each one corresponding to a point
              group symmetry of a cubic crystal
fcc?NN.py    Little code that spits out the constraints on the components
              of the ?NN force constant tensor for an FCC crystal.
fcc.py       Spits out constraints for 1-8NN. Add more NN by adding
              their one of their vectors to the list, V.

mat2vec.py   This is the interesting one.

```

Some pair of atoms in a crystal is 'connected' by some vector. The connection is made mathematically with a force constant tensor. We'll talk 3D from here on out. Without knowing any better, we write:

$$F = \begin{pmatrix} F_{11} & F_{12} & F_{13} \\ F_{12} & F_{22} & F_{23} \\ F_{13} & F_{32} & F_{33} \end{pmatrix} \quad (1)$$

And you can see that we have 9 degrees of freedom (DOF). It seems likely that some of the magical symmetries of the crystal will reduce the DOF. We can look into this by taking one of the 3×3 representations of the point group symmetries, S_s , and applying it to the F , requiring that F remain unchanged. That looks like so:

$$S_s = \begin{pmatrix} S_{11}^s & S_{12}^s & S_{13}^s \\ S_{21}^s & S_{22}^s & S_{23}^s \\ S_{31}^s & S_{32}^s & S_{33}^s \end{pmatrix} \quad (2)$$

$$S_s^T F S_s = F \quad (3)$$

That equation, is rather intimidating, so we rewrite it thusly:

$$S_s^T F - F S_s^{-1} = 0 \quad (4)$$

This looks a lot like the oh so familiar Lyapunov Equation:

$$AX + XB = C \quad (5)$$

where A , B , C and X are all square matrices of dimension N . So let us consider it. It is clear that all the terms on the left-hand side of the equation are linear in the components of X , thus, their sum must be linear in the components of X . This means that we may rewrite the equation as follows:

$$Mx + b = 0 \quad (6)$$

where M is a $N^2 \times N^2$ matrix, and y and b are N^2 -vectors. The order in which you choose to map the components of the matrix X into the vector x is arbitrary; however, once you choose an order, you must be consistent. Let's just show one example... you could write:

$$x = \begin{pmatrix} x_{11} \\ x_{12} \\ x_{13} \\ x_{21} \\ x_{22} \\ x_{23} \\ x_{31} \\ x_{32} \\ x_{33} \end{pmatrix} \quad (7)$$

Which would imply the same ordering for b , where the components would be given by $-C_{ij}$.

Let's construct the matrix $Z \equiv AX + XB$. For the 3x3 case, we have:

$$Z = \begin{pmatrix} A_{11}X_{11} + A_{12}X_{21} + A_{13}X_{31} & A_{11}X_{12} + A_{12}X_{22} + A_{13}X_{32} & A_{11}X_{13} + A_{12}X_{23} + A_{13}X_{33} \\ A_{21}X_{11} + A_{22}X_{21} + A_{23}X_{31} & A_{21}X_{12} + A_{22}X_{22} + A_{23}X_{32} & A_{21}X_{13} + A_{22}X_{23} + A_{23}X_{33} \\ A_{31}X_{11} + A_{32}X_{21} + A_{33}X_{31} & A_{31}X_{12} + A_{32}X_{22} + A_{33}X_{32} & A_{31}X_{13} + A_{32}X_{23} + A_{33}X_{33} \end{pmatrix} + \begin{pmatrix} X_{11}B_{11} + X_{12}B_{21} + X_{13}B_{31} & X_{11}B_{12} + X_{12}B_{22} + X_{13}B_{32} & X_{11}B_{13} + X_{12}B_{23} + X_{13}B_{33} \\ X_{21}B_{11} + X_{22}B_{21} + X_{23}B_{31} & X_{21}B_{12} + X_{22}B_{22} + X_{23}B_{32} & X_{21}B_{13} + X_{22}B_{23} + X_{23}B_{33} \\ X_{31}B_{11} + X_{32}B_{21} + X_{33}B_{31} & X_{31}B_{12} + X_{32}B_{22} + X_{33}B_{32} & X_{31}B_{13} + X_{32}B_{23} + X_{33}B_{33} \end{pmatrix} \quad (8)$$

We see that Z_{12} depends on the 1st row of A , the 2nd column of X , the 1st row of X , and the 2nd column of B . More generally:

$$Z_{ij} = \sum_k A_{ik}X_{kj} + B_{kj}X_{ik} \quad (9)$$

Because of our choice of mapping from X to x , the Z_{ij} tells you what row in M you are dealing with. The indices on the variables X_{kj} or X_{ik} tell you which column in M . All that is left is to add the A_{ik} and B_{kj} into the slots in M , so given. This is difficult to show explitily, however, this is the best I got:

$$M = \begin{pmatrix} M_{11} & M_{12} & \dots & M_{19} \\ M_{21} & M_{22} & \dots & M_{29} \\ \vdots & \vdots & \ddots & \vdots \\ M_{91} & M_{92} & \dots & M_{99} \end{pmatrix} \leftarrow \begin{pmatrix} \{Z_{11}, X_{11}\} & \{Z_{11}, X_{12}\} & \dots & \{Z_{11}, X_{33}\} \\ \{Z_{21}, X_{11}\} & \{Z_{21}, X_{12}\} & \dots & \{Z_{21}, X_{33}\} \\ \vdots & \vdots & \ddots & \vdots \\ \{Z_{33}, X_{11}\} & \{Z_{33}, X_{12}\} & \dots & \{Z_{33}, X_{33}\} \end{pmatrix} \quad (10)$$

The left and center of the equation are M and the components of M . The thing on the right is supposed to indicate that any time you have Z_{ij} on the left in Eq. 9, and an X_{mn} next to one of the coefficients A_{pq} or B_{rt} on the right, you add that coefficient at the slot marked $\{Z_{ij}, X_{mn}\}$ in M . The problem $AX + XB = C$ has now been reduced to $Mx + b = 0$, which linear algebra tells us how to solve.

For our particular case, we wish to find the constraints on the components of F . Simply take $A \rightarrow S_s^T$, $B \rightarrow S_s^{-1}$, $X \rightarrow F$ in Eq. 5, and $C = 0$ and then put M into reduced row echelon form. Reading off the rows gives you the constraints on the components of F . For expmpample, you may end up with something that looks like this:

$$\begin{array}{cccccccccc} xx & xy & xz & yx & yy & yz & zx & zy & zz \\ [0. & 1. & 0. & 0. & 0. & 0. & 0. & 0. & 0.] \\ [0. & 0. & 1. & 0. & 0. & 0. & 0. & 0. & 0.] \\ [0. & 0. & 0. & 1. & 0. & 0. & 0. & 0. & 0.] \\ [0. & 0. & 0. & 0. & 1. & 0. & 0. & 0. & -1.] \\ [0. & 0. & 0. & 0. & 0. & 1. & 0. & 0. & 0.] \\ [0. & 0. & 0. & 0. & 0. & 0. & 1. & 0. & 0.] \\ [0. & 0. & 0. & 0. & 0. & 0. & 0. & 1. & 0.] \end{array}$$

Which we can rewrite like so:

$$\begin{array}{rcl} xy & & = 0 \\ & xz & = 0 \\ & & yx & = 0 \\ & & & yy & -zz & = 0 \\ & & & & yz & = 0 \\ & & & & & zx & = 0 \\ & & & & & & zy & = 0 \end{array}$$

We see, then, that there are two DOF. Both $yy = zz$, and xx may be varied independently This particular force constant tensor is axially symmetric.

If there are n symmetry elements, S_s , that can transform your bond vector back onto itself, you simply stack up the n M_s , and find the reduced row echelon form of M :

$$M = \begin{pmatrix} M_1 \\ M_2 \\ \vdots \\ M_n \end{pmatrix} \quad (11)$$