

**Electronic Structure and Photochemical Reactivity  
of Binuclear Metal Complexes**

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*To Mom, Dad, and Gunilla*

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## Abstract

A valence bond (VB) "weak coupling" model of the electronic structure for  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  is developed and generalized to the class of dimeric systems in which the metals are nonbonded, in a formal sense, and can be viewed as weakly coupled. With the VB model, the energies and widths of the previously observed optical absorption bands can be rationalized; in addition, plausible assignments are made for bands that were not interpreted satisfactorily or not observed in earlier work. The VB model does not change any of the molecular orbital-based interpretations of the thermal chemistry, photochemistry, or photophysics of these systems.

Photophysical characterization of the  $^1,^3(\text{d}\sigma^*\text{p}\sigma)$  excited state of  $\text{Ir}_2(\text{TMB})_4^{2+}$  finds a system quite comparable to other binuclear  $\text{d}^8$  complexes. Both fluorescence ( $\lambda_{\text{max}}$  735 nm,  $\tau \sim 70 \pm 30$  ps) and phosphorescence ( $\lambda_{\text{max}}$  1080 nm,  $\tau = 210 \pm 20$  ns) are observed.

The relatively long lifetime of the  $^3(\text{d}\sigma^*\text{p}\sigma)$  excited state of  $\text{Ir}_2(\text{TMB})_4^{2+}$  suggests that it should be able to participate in bimolecular photochemical reactions. The diradical-like structure of the excited state, an electron (or oxidizing hole) localized on the exterior of the  $\text{M}_2$  unit (the  $\text{d}\sigma^*$  orbital) and an electron localized in the interior of the dimer cage (the  $\text{p}\sigma$  orbital), implies that one-electron chemistry will be observed. Reactions of the ground state follow two-electron pathways, similar to those observed for mononuclear  $\text{d}^8$  complexes.

The  $^3(\text{d}\sigma^*\text{p}\sigma)$  excited state of  $\text{Ir}_2(\text{TMB})_4^{2+}$  is found to be a powerful reductant,  $E^0(\text{Ir}_2(\text{TMB})_4^{3+}/^3(\text{Ir}_2(\text{TMB})_4^{2+})^*) \sim -1.0$  V (SSCE). Excited-state electron-transfer quenching by pyridinium acceptors is observed to follow classical Marcus theory for outer-sphere electron transfer. No "inverted" behavior is found. The bimolecular electron-transfer reaction is highly nonadiabatic,  $\kappa \sim 0.0001$ , because of the large donor-acceptor separation,  $\sim 8$  Å. The results for  $\text{Ir}_2(\text{TMB})_4^{2+}$  are discussed in comparison to those for  $[\text{Ir}(\mu\text{-pz})\text{COD}]_2$ .

$\text{Ir}_2(\text{TMB})_4^{2+}$  is found to react photochemically with alkyl halides. Although the  $^3(\text{d}\sigma^*\text{p}\sigma)$  excited state is a good reductant, outer-sphere electron transfer seems unlikely ( $E^0(\text{RX}/\text{RX}^{\bullet-}) < -1.5 \text{ V (SSCE)}$ ). An  $\text{S}_{\text{RN}}1$  pathway has been suggested to explain the alkyl halide photoreduction reaction observed for metal complexes with  $E^0(\text{M}_2^{+}/^3\text{M}_2^{\bullet}) < -1.5 \text{ V (SSCE)}$ ; however, atom transfer to the  $^3(\text{d}\sigma^*\text{p}\sigma)$  excited state is the favored reaction mechanism for the alkyl halide photoreduction reaction of  $\text{Ir}_2(\text{TMB})_4^{2+}$ . The generality of this reaction is discussed.

While there is some ambiguity as to the primary photoprocess for alkyl halide photoreactivity,  $^3(\text{d}\sigma^*\text{p}\sigma)$  excited-state hydrogen-atom transfer has been established as the mechanism of the reaction of  $\text{Ir}_2(\text{TMB})_4^{2+}$  and a number of organic substrates. The atom-transfer reactivity of the  $^3(\text{d}\sigma^*\text{p}\sigma)$  excited state is attributed to the presence of a hole in the  $\text{d}\sigma^*$  orbital, analogous to the  $^3\text{n}\pi^*$  state of organic ketones. Interaction of the oxidizing hole with the electron pair of the C-H bond is the presumed pathway.

Electrochemical oxidation of  $\text{Rh}_2(\text{TMB})_4^{2+}$  generates the  $\text{d}^8\text{-d}^7$  species  $\text{Rh}_2(\text{TMB})_4^{3+}$ . This complex reacts with 1,4-cyclohexadiene to abstract a hydrogen atom mimicking the initial step of the  $^3(\text{d}\sigma^*\text{p}\sigma)$  photoreaction. The importance of this result is discussed in terms of energy storage systems and extension of the range of hydrocarbon oxidations with binuclear  $\text{d}^8$  complexes.

The  $\text{d}^7\text{-d}^7$  dihydride product obtained from the photoreaction of  $\text{Ir}_2(\text{TMB})_4^{2+}$  and 1,4-cyclohexadiene is isolated and characterized. In addition to NMR, UV-Vis, IR, and Raman spectra, the complex is characterized crystallographically. The reactivity of this complex is also discussed.

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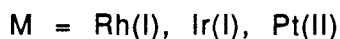
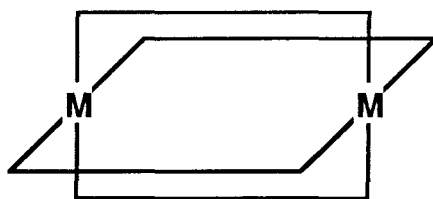
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# **Chapter 1**

## **Introduction**

Binuclear metal complexes are attractive systems for the activation of organic and inorganic substrates because they offer the possibility of cooperative involvement of their active sites: multiple binding sites for substrates; multielectron redox capabilities; and for heteronuclear complexes, distinct metal centers that can interact differently with a given substrate.<sup>1-5</sup> Numerous studies have focused on the thermal reactivity of binuclear complexes with a variety of substrates, with less attention having been given to the activation of substrates by the electronic excited states of these complexes.<sup>6-8</sup> Since a photoexcited molecule is best envisioned as a new species possessing physical and chemical properties distinct from its corresponding ground state molecule, photoexcited binuclear complexes may possess reactivity very different from the ground state.<sup>9,10</sup> It is therefore of interest to investigate the possibility of the conversion of light to net chemical energy through a photochemical process.<sup>11-13</sup> Binuclear  $d^8$  complexes with the general structure shown in Figure 1.1 have been



**Figure 1.1.** General Structure for a face-to-face binuclear  $d^8$  complex. Bidentate ligands are diisocyanoalkanes or  $\text{P}_2\text{O}_5\text{H}_2^{4-}$  (e.g.,  $\text{Rh}_2\text{b}_4^{2+}$  ( $b = 1,3$ -diisocyanopropane),  $\text{Rh}_2(\text{TMB})_4^{2+}$  ( $\text{TMB} = 2,5$ -diisocyano-2,5-dimethylhexane),  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$ ). For A-frame  $d^8$  complexes two bridging ligands (e.g.,  $[\text{Ir}(\mu\text{-pz})\text{COD}]_2$ ,  $\text{pz} = \text{pyrazole}$ ,  $\text{COD} = 1,5$ -cyclooctadiene).

studied for several years because of their unique spectral properties, thermal reactivity, and photochemical reactivity. The work in this area grew out of studies of the spectroscopic properties of Rh(I) isocyanide complexes.

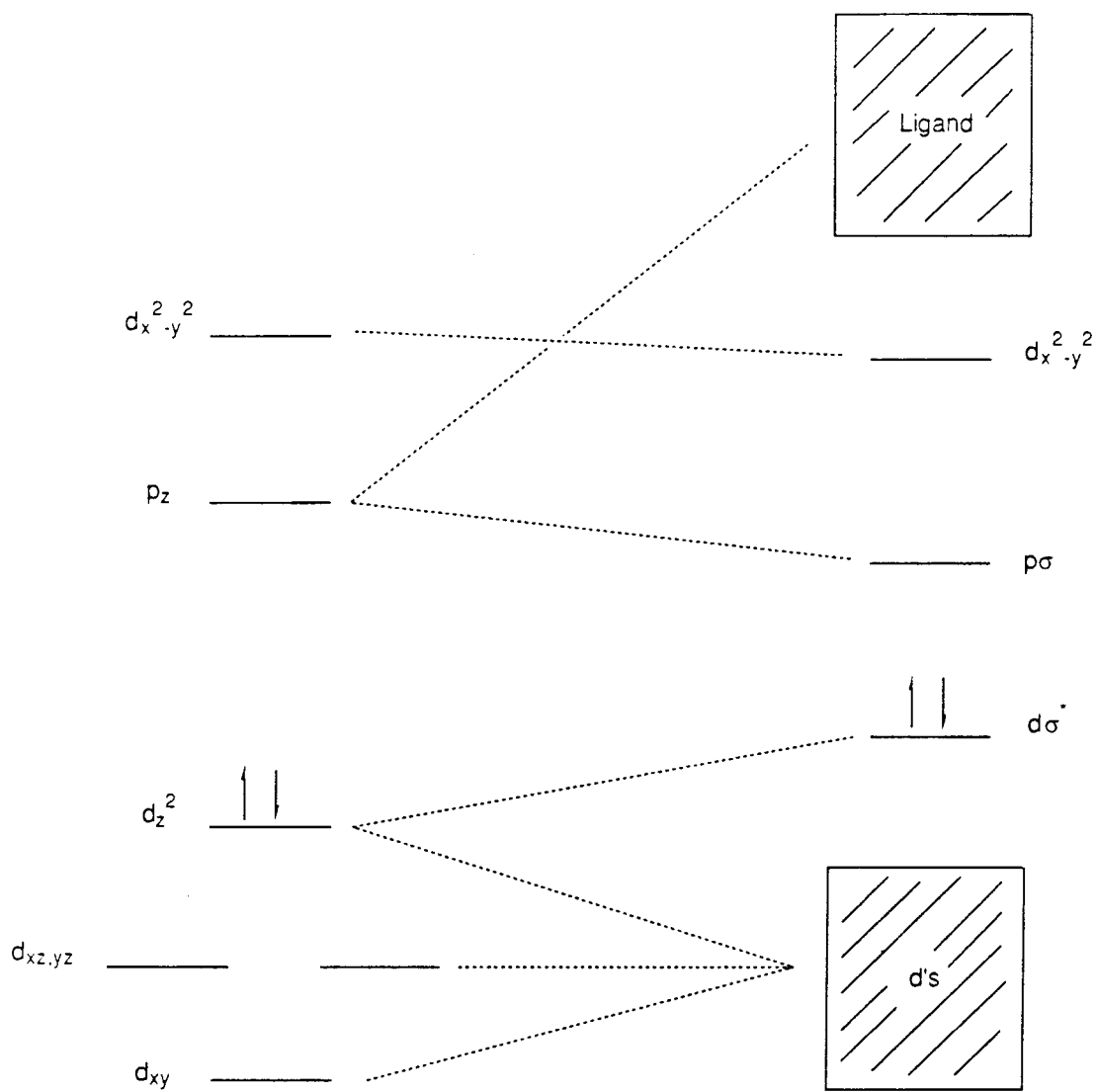
For years Rh(I) and Ir(I) isocyanide complexes have been suspected of existing as oligomers owing to the intense color, uncharacteristic of  $d^8$  complexes possessing  $\pi$ -acceptor ligands.<sup>14,15</sup> It was not until Mann demonstrated that the bands in the visible and near-infrared region of the spectrum for  $\text{Rh}(\text{CNC}_6\text{H}_5)_4^+$  were not linearly related to the concentration of monomer but to some root (e.g., square, cubic, etc.) of monomer concentration that the extent of oligomerization was realized.<sup>16</sup> The existence of discrete species was later confirmed by x-ray diffraction analysis.<sup>17,18</sup>

Stacking of square planar  $d^8$  complexes has been observed in the solid state.<sup>19</sup> Electrostatic binding (attraction of oppositely charged, square planar units) was thought to be the important interaction. Oligomerization of similarly charged, square planar units could be understood to occur only from some favorable metal-metal interaction. Metal-metal interactions for the stacked linear chain complexes have been discussed in reference to the d/p band structure and the observed electrical conductivities. The dark colors of the linear chain salts and the solid-state emission are discussed in terms of  $d \rightarrow p$  transitions.<sup>20</sup>

Binuclear Rh species were later synthesized using a series of diisocyanoalkanes, ligands chosen to favor formation of dimeric complexes.<sup>21,22</sup> Characterization of these complexes found them to be analogous to the oligomeric Rh species studied earlier.

The electronic structure of binuclear  $d^8$  complexes has been understood in terms of a simple molecular orbital (MO) model (Figure 1.2).<sup>16</sup> Starting from a monomer orbital scheme, two square planar units are brought together in a face-to-face orientation. The orbitals perpendicular to the molecular plane,  $d_{z^2}$  and  $p_z$ , interact strongly, yielding  $d\sigma/\sigma^*$  and  $p\sigma/\sigma^*$  orbitals. To a first approximation the ground state is expected to be nonbonding since both the  $d\sigma$  and  $d\sigma^*$  orbitals are filled. However,

**Figure 1.2.** Molecular orbital diagram for the interaction of two square planar  $d^8$  metal ions.



extensive spectroscopic studies of binuclear complexes have established that the metal-metal interaction in the  $^1A_{1g}(d\sigma^*)^2$  ground state is weakly bonding.<sup>23,24</sup>

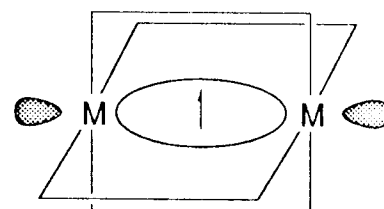
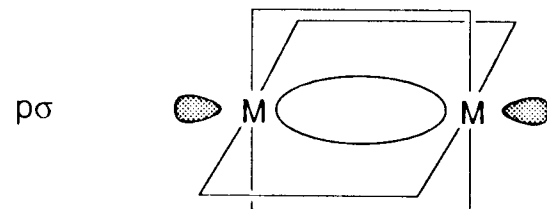
Perturbational mixing of orbitals of the same symmetry will stabilize the lower set and destabilize the upper ones; stabilization of the filled lower set is viewed as a source of metal-metal bonding.

As suggested by the MO model and well established for numerous other studies, the lowest energy transition is  $d\sigma^* \rightarrow p\sigma$ .<sup>22,25</sup> This excitation results in formation of a metal-metal single bond in the excited state. The spectral and photophysical properties of all binuclear  $d^8$  complexes are found to be dominated by the large contraction along the metal-metal coordinate that occurs upon excitation.<sup>25-27</sup> The  $d\sigma^* \rightarrow p\sigma$  transition is metal-localized and can be viewed as movement of an electron from an orbital localized on the exterior of the  $M_2$  unit (the  $d\sigma^*$  orbital) to an orbital localized in the interior of the dimer cage (the  $p\sigma$  orbital). The excitation results in hole formation on a metal center at an open coordination site (Figure 1.3). The diradical picture of the excited state is important in discussing the photochemistry and the electrochemistry of these systems.

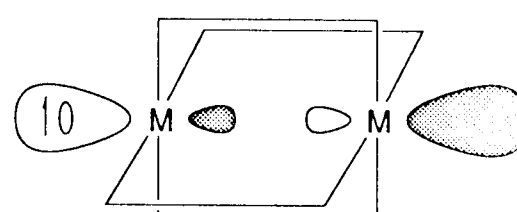
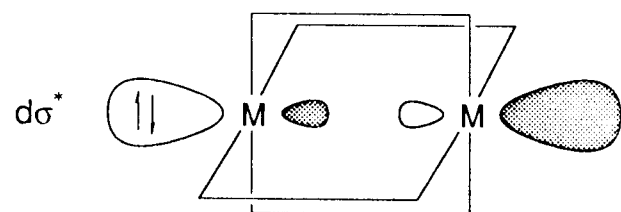
Additional studies of the binuclear Rh(I) complexes revealed that the lowest energy singlet and triplet excited states derived from the  $(d\sigma)^2(d\sigma^*)^1(p\sigma)^1$  electronic configuration are luminescent at ambient temperatures in fluid solution.<sup>22-24,28</sup> (For stacked linear chain complexes luminescence has been observed in the solid state and discussed in terms of a  $dp$  excited state.)<sup>20</sup> The  $^1(d\sigma^*p\sigma)$  excited-state lifetime, in most cases, is less than 1 ns. The  $^3(d\sigma^*p\sigma)$  excited-state lifetime is found to range from 30 ns to 10  $\mu$ s. Ground-state and excited-state Raman studies found that the metal-metal vibrational frequency increases upon excitation, an increase in the metal-metal bond strength.<sup>26</sup> The increased metal-metal interaction is also manifest in the low-temperature absorption spectra.<sup>25,27</sup> The band at 670 nm in the spectrum of  $Rh_2b_4^{2+}$



**Figure 1.3.** Pictorial representation of the  $M_2$ -localized hole in a  $^3(d\sigma^*p\sigma)$  state.



$h\nu$



assigned to the  $^1A_{1g} \rightarrow ^3E_u$  transition shows a vibrational progression in a frequency of  $\sim 150\text{ cm}^{-1}$ , consistent with the Rh-Rh stretching frequency of  $144\text{ cm}^{-1}$  obtained from the Raman studies. (The ground-state Rh-Rh vibrational frequency is  $79\text{ cm}^{-1}$ .) Analysis of the band shape suggests that the Rh-Rh distance decreases  $0.3\text{ \AA}$  in the  $^3(d\sigma^*p\sigma)$  excited state, consistent with the increased metal-metal bonding interaction in the excited state.

The MO picture also yields insight into the chemistry of these systems. For monomeric  $d^8$  complexes, the predominant mode of reaction is oxidative addition.<sup>29</sup> The chemistry is controlled by the  $(d_z^2)^2$  electron configuration, and while most reactions are classified as oxidative addition, the pathways followed are electron transfer, nucleophilic attack by the metal center, or acid/base. For binuclear  $d^8$  complexes, the chemistry of the ground state is expected to be controlled by the  $(d\sigma^*)^2$  electron configuration. Similar to the monomeric species, the binuclear complexes, undergo oxidative addition reactions. Many complexes have been found to undergo two-electron oxidative addition reactions with dihalides and alkyl halides.<sup>21,30,31</sup> The pathway for these reactions are not well understood. In analogy to the monomeric complexes, the dimers are expected to undergo electron transfer, nucleophilic attack by the metal, and acid/base reactions. The transfer of two  $d\sigma^*$  electrons results in the formation of a metal-metal single bond. The presence of this bond has been verified by x-ray diffraction.<sup>32,33</sup> Other manifestations of the metal-metal bond have been observed spectroscopically.<sup>34</sup>

The lifetime of the  $^3(d\sigma^*p\sigma)$  excited state, normally in between 100 ns and 10  $\mu\text{s}$ , makes binuclear  $d^8$  complexes attractive for bimolecular photoprocesses. An electronic excited state of a metal complex is both a stronger reductant and a stronger oxidant than the ground state.<sup>10</sup> Excitation of an electron from a low-energy to a high-energy orbital reduces the ionization potential and increases the electron affinity of a molecule. Therefore, binuclear  $d^8$  complexes with relatively long-lived excited states

can participate in intermolecular electron-transfer reactions that are uphill for the corresponding ground-state species. Such excited-state electron-transfer reactions often play key roles in multistep schemes for the conversion of light to chemical energy.<sup>9</sup>

While electron-transfer processes are common in inorganic photochemistry, excited-state atom transfer is limited to a small class of inorganic complexes. For  $\text{UO}_2^{2+}$ , the diradical excited state ( $\bullet\text{U}-\text{O}\bullet$ ) is active in alcohol oxidation.<sup>35</sup> The primary photoprocess is hydrogen atom abstraction by the oxygen-centered radical. Photoaddition to a metal center via atom transfer has been observed for binuclear metal complexes such as  $\text{Re}_2(\text{CO})_{10}$ .<sup>36-39</sup> The primary photoprocess is metal-metal bond homolysis. The photogenerated metal radical undergoes thermal atom-abstraction reactions. Until recently, atom transfer to a metal-localized excited state had not been observed.

Atom transfer to a metal complex is facilitated if localized electron or hole generation occurs at one or more open coordination sites. Binuclear  $d^8$  complexes have been found to undergo photochemical atom transfer to one of the metal centers.<sup>40-43</sup> These complexes possess open coordination sites in addition to an electronic structure that localizes the electron (hole) necessary for atom transfer to the metal center.

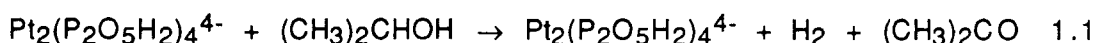
The initial interest in the binuclear  $d^8$  complexes was stimulated by observations of their photochemical electron-transfer reactivity.<sup>41</sup> From spectroscopic and electrochemical studies, the  $^3(d\sigma^*p\sigma)$  excited state is predicted to be a powerful reductant, with  $E^0(\text{M}_2^+/\text{}^3\text{M}_2^*)$  estimated to range from -0.8 to -2.0 V vs SSCE in  $\text{CH}_3\text{CN}$ . That this state is a powerful reductant has been confirmed by investigation of the electron-transfer quenching of  $^3\text{M}_2^*$  by a series of pyridinium acceptors with varying reduction potentials. For several binuclear complexes, the excited-state reduction potential cannot be calculated accurately because of the irreversibility of the

ground-state electrochemistry, but it can be estimated from bimolecular electron-transfer quenching experiments.

For systems that are powerful excited-state reductants, photoreduction of alkyl halides is observed.<sup>41,44</sup> This reaction was initially interpreted to be an outer-sphere electron transfer to form an alkyl-halide radical anion, which rapidly decomposes to yield  $R\cdot$  and  $X^-$ . Subsequent thermal reactions give the observed products, an  $S_{RN}1$  mechanism (Figure 1.4). While such a mechanism,  $S_{RN}1$ , appears plausible for a metal complex with  $E^0(M_2^{+}/^3M_2^*) < -1.5$  V (SSCE), it seems unlikely for complexes with  $E^0(M_2^{+}/^3M_2^*) > -1.0$  V (SSCE). Reduction potentials for alkyl halides of interest are generally more negative than -1.5 V (SSCE).<sup>45</sup> An alternative pathway to outer-sphere electron transfer, which yields similar photoredox products with alkyl halides, is excited-state atom transfer (Figure 1.5).

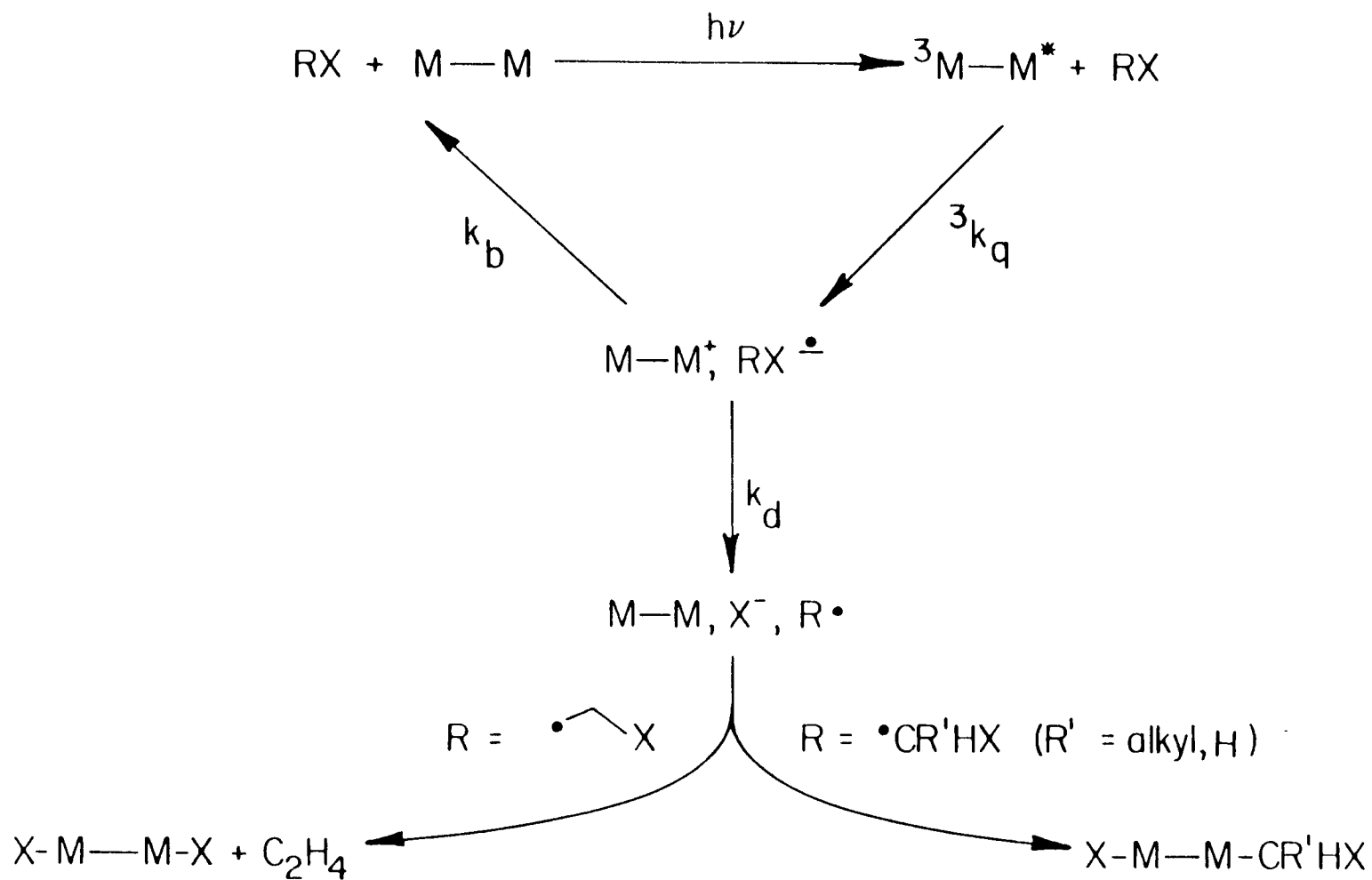
Although the primary photoprocess for alkyl-halide photoreduction may not be atom transfer in all cases,  $^3(d\sigma^*p\sigma)$  excited-state hydrogen atom transfer has been established as the mechanism of the reaction between  $Pt_2(P_2O_5H_2)_4^{4-}$  and a number of organic and organometallic substrates.<sup>40-43</sup> Initial work in this area focused on the catalytic conversion of isopropanol to acetone (Equation 1.1).<sup>46</sup>

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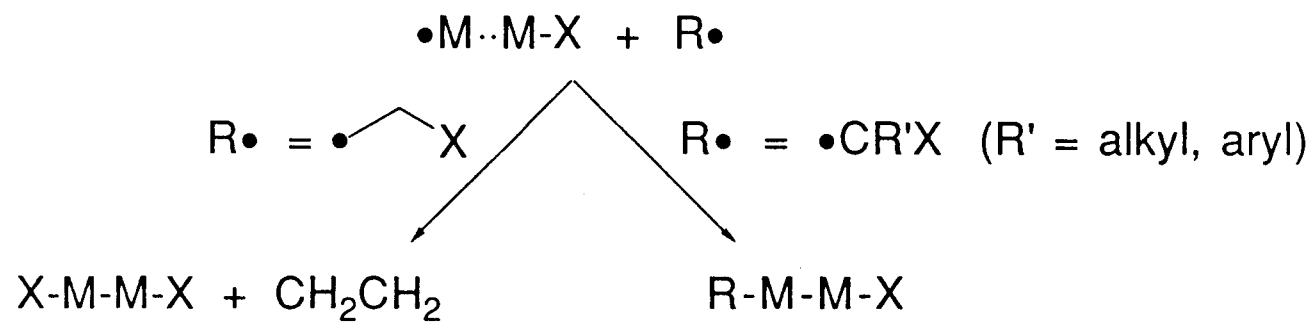
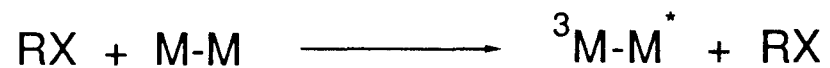
From detailed studies of this system, it was concluded that the primary photoprocess is abstraction of the  $\alpha$ -hydrogen by the  $^3Pt_2^*$  to form a monohydride species (directly observed by transient absorption spectroscopy for a number of substrates) and the organic radical (Equation 1.2), with the final photoproduct being  $Pt_2H_2$  and acetone (Equation 1.3). The  $Pt_2H_2$  complex has been characterized by NMR, UV-Vis, and IR (but has not been successfully isolated).<sup>47</sup>

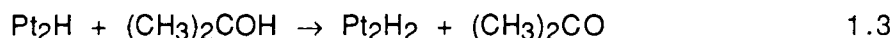
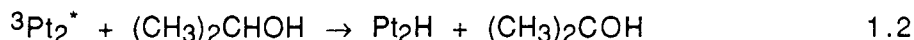
**Figure 1.4.**  $S_{RN}1$  mechanistic scheme for halocarbon photooxidative addition to binuclear  $d^8$  complexes.



**Figure 1.5.** Atom-transfer mechanism for halocarbon photooxidative addition.







To further extend the understanding of the chemistry and spectroscopy of binuclear  $d^8$  complexes, an Ir(I) analog of the Rh(I) species was prepared,  $\text{Ir}_2(\text{TMB})_4^{2+}$ .<sup>30</sup> This complex was found to compare well to the Rh(I) dimers with the expected red shift of the  $d\sigma^* \rightarrow p\sigma$  transition. The  $^3(d\sigma^*p\sigma)$  excited state was found to have a temperature-independent lifetime of approximately 200 ns.

The crystal structure of the  $\text{B}(\text{C}_6\text{H}_5)_4^-$  salt of  $\text{Ir}_2(\text{TMB})_4^{2+}$  is described in Chapter 2. The optical absorption spectrum is presented and discussed in reference to a large number of binuclear  $d^8$  complexes. Photophysical characterization of the emissive  $^1,^3(d\sigma^*p\sigma)$  excited states is reported along with Raman vibrational data for the ground state and the  $^3(d\sigma^*p\sigma)$  excited state.

A detailed study of the electronic spectrum of  $\text{Ir}_2(\text{TMB})_4^{2+}$  is presented in Chapter 3. A change in the metal from rhodium to iridium should perturb the electronic structure, thereby altering the characteristic  $d^8$ - $d^8$  electronic spectrum and allowing key aspects of the MO model to be tested. Electronic absorption and magnetic circular dichroism spectra of  $\text{Rh}_2(\text{TMB})_4^{2+}$  and  $\text{Ir}_2(\text{TMB})_4^{2+}$  are reported along with polarized single-crystal absorption spectra of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2 \cdot \text{CH}_3\text{C}_6\text{H}_5$ .

Earlier work found that  $\text{Rh}_2\text{b}_4^{2+}$ ,  $\text{Rh}_2(\text{TMB})_4^{2+}$ ,  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$ , and  $[\text{Ir}(\mu\text{-pz})\text{COD}]_2$  react with dihalides to yield the corresponding  $d^7$ - $d^7$  dimer,  $\text{X-M-M-X}$ .<sup>21,40,31</sup>  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$  and  $[\text{Ir}(\mu\text{-pz})\text{COD}]_2$  were also found to oxidatively add  $\text{CH}_3\text{I}$  to yield  $\text{CH}_3\text{-M-M-I}$ . It was found that  $\text{Ir}_2(\text{TMB})_4^{2+}$  oxidatively adds dihalides ( $\text{Cl}_2$ ,  $\text{Br}_2$ ,  $\text{I}_2$ ),  $\text{CH}_3\text{I}$ ,  $\text{HCl}$ , and  $\text{CH}_2(\text{CN})_2$  to yield the corresponding  $d^7$ - $d^7$  complex,  $\text{X-M-M-Y}$  ( $\text{X} = \text{Y}$ ,  $\text{Cl}_2$ ,  $\text{Br}_2$ ,  $\text{I}_2$ ;  $\text{X} = \text{I}$ ,  $\text{Y} = \text{CH}_3$ ;  $\text{X} = \text{Cl}$ ,  $\text{Y} = \text{H}$ ;  $\text{X} = \text{CH}(\text{CN})_2$ ,  $\text{Y} = \text{H}$ ).<sup>30</sup> Because of the greater reactivity expected for Ir and the interest in the

photochemical reactivity of the excited states of  $\text{Ir}_2(\text{TMB})_4^{2+}$ , extension of the thermal chemistry was pursued. Thermal reactions with nonpolar and low-polarity reagents ( $\text{H}_2$ ,  $\text{R}_3\text{CH}$ ,  $\text{R}_3\text{SiH}$ ,  $\text{R}_3\text{SnH}$ ), and electrophilic reagents ( $\text{X}_2$ ,  $\text{RX}$ ,  $\text{HX}$ ,  $\text{CH}_3\text{COCl}$ ) are discussed in Chapter 4.

The relatively long lifetime of the  $^3(\text{d}\sigma^*\text{p}\sigma)$  excited state of  $\text{Ir}_2(\text{TMB})_4^{2+}$  in fluid solution suggests that it should be able to participate in bimolecular photochemical reactions. In light of the favorable absorption and emission properties of this complex, it is of interest to examine the possibility of the conversion of visible light to net chemical energy through a photochemical process, either electron transfer or atom transfer. The excited-state electron-transfer reactivity of  $\text{Ir}_2(\text{TMB})_4^{2+}$  with pyridinium acceptors is reported and analyzed in the context of classical Marcus theory for outer-sphere electron transfer in Chapter 4. The photochemical reactions of  $\text{Ir}_2(\text{TMB})_4^{2+}$  with alkyl halides and hydrogen-atom donors are reported.

Hydrogen-atom transfer has been established as an important reaction pathway for the triplet  $\text{d}\sigma^*\text{p}\sigma$  excited state of binuclear  $\text{d}^8$  complexes. Substrates that serve as hydrogen-atom donors include hydrocarbons (e.g., cyclohexene and 1,4-cyclohexadiene), alcohols with  $\alpha(\text{C-H})$  bonds, and triorganosilanes, -germanes, and -stannanes. Irradiation of the metal complex in the presence of these substrates produces complexes of the type  $\text{M}_2\text{H}_2$ . Earlier work has characterized the product formed with  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$  as the binuclear platinum(III) dihydride,  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{H}_2^{4-}$  ( $\text{Pt}_2\text{H}_2$ ). In Chapter 5, the characterization of the dihydride of  $\text{Ir}_2(\text{TMB})_4^{2+}$  ( $\text{Ir}_2\text{H}_2$ ) is reported. In addition to NMR, UV-Vis, IR, and Raman spectra, the complex has been characterized crystallographically. The reactivity of  $\text{Ir}_2\text{H}_2$  is also reported.

Theoretical studies have also contributed to the study of the electronic structure and the spectral properties of metal dimers.<sup>48,49</sup> Results of *ab initio* calculations of  $\text{Re}_2\text{Cl}_8^{2-}$  designed to provide accurate bond energies and torsion barriers as well as

accurate shapes for the potential curves are reported in Chapter 6. These studies use the generalized valence bond (GVB) approach in which electron correlations are included for all eight electrons available for the quadruple bond, while solving self-consistently for all orbitals.

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**Chapter 2**  
**Structure, Optical Absorption Spectrum, and Photophysics of**  
 **$\text{Ir}_2(\text{TMB})_4^{2+}$**

## Introduction

Extensive studies of binuclear  $d^8$  complexes have revealed a class of compounds that are spectroscopically unique and chemically very interesting. The ground state of these complexes, while formally nonbonding, does possess a weak metal-metal bond. The electronic excited states have a formal metal-metal single bond. This arises from excitation of an electron from the  $d\sigma^*$  orbital to the  $p\sigma$  orbital. Such a nonbonding-to-bonding transition is unique among inorganic complexes. In studying binuclear  $d^8$  complexes, one wishes to understand how the metal-metal interaction in the ground state and excited states influence the spectroscopy and chemistry.

Many of the binuclear  $d^8$  complexes are found to be emissive. Excited states formed by absorption of a photon rather than merely by undergoing thermal deactivation to the ground state decay by emission of a photon. The emissive states have been shown to be the  $^1,^3(d\sigma^*p\sigma)$  excited states. The  $^1(d\sigma^*p\sigma)$  excited state lifetime, in most cases, is less than 1 ns. The  $^3(d\sigma^*p\sigma)$  lifetime is found to range from 30 ns to 10  $\mu$ s. The structure and properties of these excited states have been the focus of much work. One would like to characterize these states much as one would characterize the ground state.

To further extend the understanding of the chemistry and spectroscopy of binuclear  $d^8$  complexes, a study of an Ir(I) complex,  $\text{Ir}_2(\text{TMB})_4^{2+}$ , was undertaken. The synthesis and preliminary characterization of  $\text{Ir}_2(\text{TMB})_4^{2+}$  has been reported.<sup>1</sup> The crystal structure of the  $\text{B}(\text{C}_6\text{H}_5)_4^-$  salt of  $\text{Ir}_2(\text{TMB})_4^{2+}$  is presented. The optical absorption spectrum of  $\text{Ir}_2(\text{TMB})_4^{2+}$  is discussed in reference to a large number of binuclear  $d^8$  complexes. It is shown that the energy of the  $^1(d\sigma^* \rightarrow p\sigma)$  transition is dictated by excited-state properties, not by ground-state interactions as prescribed by the MO model. Photophysical characterization of the emissive  $^1,^3(d\sigma^*p\sigma)$  excited states is reported along with Raman vibrational data for the ground state and the  $^3(d\sigma^*p\sigma)$  excited state.

## Experimental

### Synthesis

All synthetic procedures were carried out with standard Schlenk techniques. All solvents were from freshly opened bottles with no further purification. All solvents were Schlenk-degassed prior to use. Standard procedures were used to prepare 2,5-diisocyano-2,5-dimethylhexane (TMB)<sup>2</sup> and  $[\text{Ir}(\text{COD})\text{Cl}]_2$  (COD = 1,4-cyclooctadiene)<sup>3</sup>. All other chemicals were of reagent grade or comparable quality and were used as received. The  $^1\text{H}$  NMR spectra were obtained on a 400-MHz JNM-GX400 FT NMR spectrometer. The IR spectra were measured on a Beckman IR 4240.

#### $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$

To a filtered acetonitrile solution of  $[\text{Ir}(\text{COD})\text{Cl}]_2$  (0.1 g in 10 ml  $\text{CH}_3\text{CN}$ ), a filtered solution of TMB (0.1g in a minimum of  $\text{CH}_3\text{CN}$ ) was added with stirring. The solution immediately turned dark blue. The mixture was allowed to stir for approximately 15 minutes, at which time the solvent and COD were removed under vacuum. The resulting metallic-brown solid was dissolved in methanol. A filtered methanol solution of  $\text{NaB}(\text{C}_6\text{H}_5)_4$  (0.11g, 10 ml) was added to the stirred methanol solution of  $[\text{Ir}_2(\text{TMB})_4]\text{Cl}_2$ . Immediate precipitation of a blue solid occurred. The solution was filtered and the blue powder washed with cold methanol. The solid could be recrystallized from acetone or acetonitrile. Yields were quite variable. Calculated for  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$ : C, 62.91; H, 6.24; N, 6.67. Found: C, 62.61; H, 6.43; N, 6.40. NMR:  $\delta$  ( $\text{CD}_3\text{CN}$ , 20 °C); 1.49 (broad singlet,  $\text{CH}_3$ , 48H), 1.86 (broad singlet,  $\text{CH}_2$ , 16H), 6.83(triplet-of-triplets, 8H), 6.98 (triplet, 16H), 7.25 (multiplet, 16H). IR:  $\nu$ ,  $\text{cm}^{-1}$  (Nujol mull, NaCl plates); 2140 (vs,  $\text{N}\equiv\text{C}$ ). Free ligand IR:  $\nu$ ,  $\text{cm}^{-1}$  (Nujol mull, NaCl plates); 2120 (vs,  $\text{N}\equiv\text{C}$ ).

## X-ray Data Collection and Reduction for $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2 \cdot 2\text{CH}_3\text{CN}$

Crystal data are given in Table 2.1. Oscillation and Weissenberg photographs showed monoclinic symmetry. Bright-blue crystals, freshly grown from acetonitrile, were cut to size with a razor blade. An irregular chunk was glued to a glass fiber and then quickly covered with epoxy glue to retard decomposition. It was centered on a Nonius CAD-4 diffractometer equipped with graphite-monochromated  $\text{MoK}\alpha$  radiation, and the cell dimensions plus an orientation matrix were obtained from the setting angles of 25 reflections with  $22^\circ < 2\theta < 25^\circ$ . Two crystals were used to measure 6332 reflections with an  $\omega$  scan. A scale factor relating the two data sets was calculated from the three check reflections, which were measured every 10,000 seconds of x-ray exposure. These showed a linear decay of about 9% in  $I$  for each crystal, and the data were corrected for this decay. Backgrounds were measured for each reflection at each end of the scan; an average background as a function of  $2\theta$  was calculated and used to correct the measured scan counts. Absences in the data of  $0k0$ ,  $k=2n+1$  and  $h0l$ ,  $h+l=2n+1$  identify the space group as  $P2_1/n$ , #14. After deleting space-group absences and merging equivalent reflections, 5562 reflections used in the structure solution and refinement remained, of which 5116 had  $F_o^2 > 0$  and 4333 had  $F_o^2 > 3\sigma(F_o^2)$ . The data were corrected for Lorentz and polarization factors but not for absorption; the epoxy glue on the crystals prevented their measurement and identifying the crystal faces. Variances of the individual reflections were assigned, based on counting statistics plus an additional term,  $0.014I^2$ . Variances for the merged reflections were obtained by standard propagation of error plus another additional term,  $0.014I^2$ . Scattering factors were taken from Reference 4.

The coordinates of the iridium atoms were obtained from a Patterson map and the remaining nonhydrogen atoms were found with successive structure factor-Fourier calculations. After three cycles of full-matrix least squares, minimizing  $\sum w(F_o^2 -$

**Table 2.1.** Crystal Data for  $[\text{Ir}_2(\text{C}_{10}\text{H}_{16}\text{N}_2)_4](\text{B}(\text{C}_6\text{H}_5)_4)_2 \cdot 2\text{CH}_3\text{CN}$ .

Formula: $\text{Ir}_2\text{C}_{92}\text{H}_{110}\text{N}_{10}\text{B}_2$	Formula Wt.: 1761.99
$a = 17.053(2) \text{ \AA}$	Space Group $P2_1/n$ (#14)
$b = 32.481(11) \text{ \AA}$	
$c = 17.616(5) \text{ \AA}$	T: 22 °C
$\beta = 115.59(2)^\circ$	
$V = 8800(5) \text{ \AA}^3$	$\lambda \text{ MoK}\alpha = 0.71073 \text{ \AA}$
$\mu = 32.56 \text{ cm}^{-1}$	$d_{\text{calc}} = 1.329 \text{ g cm}^{-3}$
Crystal irregular, max. dim. ~0.4 mm	
$F(000) = 3548 \text{ electrons}$	

$F_o^2$ ), hydrogen atoms were introduced at calculated positions, assuming a C-H distance of 0.95 Å and staggered geometries. Each hydrogen atom was given an isotropic thermal parameter 10% greater than that of the carbon atom to which it was bonded. The least-squared converged (iridium atoms and methyl carbon atoms of the ligands anisotropic, the remaining atoms isotropic: 106 nonhydrogen atoms, 515 parameters, plus 104 hydrogen atoms) with  $R(=\Sigma(F_o-F_c)/\Sigma F_o)$  of 0.062 for all reflections with  $F_o^2 > 0$  and 0.050 for the reflections with  $F_o^2 > 3\sigma(F_o^2)$ . The goodness-of-fit,  $[\Sigma w(F_o^2 - F_c^2)^2 / (n-p)]^{1/2}$ , was 2.64. The final parameters are given in Table 2.2 and selected distances and angles in Table 2.3. The structure factors, hydrogen parameters and further distances and angles are given in Appendix 1.

A final difference Fourier map showed excursions of +1.00 and -1.12 electrons /Å<sup>3</sup> (e/Å<sup>3</sup>) in the vicinity of the iridium atoms, two peaks of 1.15 and 0.97 e/Å<sup>3</sup> along the Ir-Ir vector, 3.2 Å from the Ir atoms, and several peaks of about 0.75 e/Å<sup>3</sup> in the vicinity of the acetonitrile molecules. The general noise level of the map was  $\pm 0.4$  e/Å<sup>3</sup>. The large excursions near the iridium atoms, as well as the two axial peaks, are caused by the inability to correct for absorption. The peaks near the acetonitrile molecules probably represent different orientations of partial solvent molecules, but they could not be interpreted in any chemically reasonable way. The rather high goodness-of-fit is the result of both of these factors.

The lack of an absorption correction has caused some of the carbon atoms to be mispositioned, particularly those bonded to iridium. The wide range of Ir-C distances ( $>0.1$  Å) is surely not real but is a result of the least squares moving the carbon atoms to minimize the peaks around the iridium atoms caused by the absorption. This is also the reason that two of the carbon atoms, C11 and C31, have thermal parameters that are too small: the least squares is adjusting them so as to compensate as much as possible for the absorption effects.

Table 2.2. Final Parameters (x,y,z x 10<sup>4</sup>).

Atom	x	y	z	B
Ir1	1252(0.4)	648(0.2)	1031(0.4)	338(2)*
Ir2	1457(0.4)	1628(0.2)	1084(0.4)	330(2)*
C1	-20(11)	660(5)	415(10)	3.6(4)
N1	-760(9)	660(4)	82(8)	3.8(3)
C2	-1695(10)	738(5)	-391(10)	2.9(4)
C3	-2182(10)	468(5)	-41(12)	674(62)*
C4	-1937(10)	611(5)	-1293(9)	621(59)*
C5	-1829(10)	1197(5)	-287(10)	3.7(4)
C6	-1689(10)	1338(5)	570(10)	3.4(4)
C7	-1284(10)	1767(5)	826(10)	3.4(4)
C8	-1790(10)	2105(5)	198(12)	724(63)*
C9	-1227(11)	1885(5)	1708(10)	650(56)*
N2	-389(8)	1757(4)	898(8)	3.3(3)
C10	306(10)	1704(5)	968(10)	3.5(4)
C11	1398(8)	629(4)	-28(9)	1.1(3)
N3	1518(8)	629(4)	-588(9)	3.9(3)
C12	1621(12)	663(6)	-1387(11)	4.0(4)
C13	2533(13)	598(7)	-1201(13)	962(69)*
C14	1050(15)	320(6)	-1978(12)	1141(88)*
C15	1269(11)	1087(6)	-1770(11)	4.9(4)
C16	348(11)	1183(6)	-1987(11)	5.1(5)
C17	160(11)	1616(6)	-1802(12)	4.2(4)
C18	455(11)	1949(5)	-2206(10)	593(59)*
C19	-773(10)	1676(6)	-1978(12)	866(68)*
N4	635(8)	1670(4)	-889(10)	4.3(3)
C20	965(10)	1657(5)	-180(12)	3.9(4)
C21	2495(11)	596(5)	1679(11)	4.1(4)
N5	3224(9)	562(4)	2058(8)	4.2(3)
C22	4201(11)	542(6)	2566(12)	4.6(5)
C23	4344(12)	485(7)	3497(11)	1060(84)*
C24	4542(11)	174(5)	2280(12)	839(70)*
C25	4561(11)	953(5)	2490(11)	4.8(4)
C26	4458(11)	1066(6)	1596(12)	5.5(5)
C27	4216(11)	1499(5)	1326(11)	4.2(4)
C28	4860(11)	1814(6)	1945(13)	989(73)*
C29	4140(13)	1567(6)	431(12)	971(64)*
N6	3346(9)	1577(4)	1291(8)	3.8(3)
C30	2633(10)	1599(5)	1206(10)	3.5(4)
C31	1060(8)	641(4)	2079(9)	0.9(3)
N7	929(8)	645(4)	2632(9)	3.3(3)
C32	761(10)	705(5)	3382(10)	2.5(4)
C33	1177(11)	350(5)	3973(10)	671(55)*
C34	-183(10)	714(6)	3110(11)	800(58)*
C35	1149(10)	1116(5)	3764(10)	4.3(4)

Atom	$x$	$y$	$z$	$B$
C36	2113(10)	1162(5)	4093(10)	4.7(4)
C37	2455(11)	1576(6)	3953(11)	4.4(4)
C38	3429(10)	1586(6)	4335(11)	887(70)*
C39	2041(13)	1924(6)	4206(12)	1021(83)*
N8	2163(8)	1619(4)	3021(10)	4.9(4)
C40	1866(10)	1632(5)	2322(11)	3.7(4)
B1	8189(11)	1149(6)	4764(12)	3.0(4)
C41	7796(10)	944(5)	3829(10)	3.1(4)
C42	7949(11)	1122(6)	3176(12)	4.8(4)
C43	7651(12)	940(6)	2366(12)	5.6(5)
C44	7254(11)	571(6)	2211(12)	5.5(5)
C45	7098(12)	376(6)	2834(13)	6.7(5)
C46	7371(11)	571(6)	3621(11)	4.8(4)
C51	7949(9)	1623(5)	4734(10)	3.2(4)
C52	8273(11)	1854(6)	5486(12)	5.8(5)
C53	8073(13)	2269(7)	5549(13)	7.4(6)
C54	7495(13)	2458(6)	4827(14)	7.4(5)
C55	7111(12)	2245(6)	4061(13)	6.9(5)
C56	7362(11)	1839(6)	4031(11)	5.1(4)
C61	7755(10)	932(5)	5329(10)	3.0(4)
C62	6868(12)	963(5)	5068(11)	5.1(5)
C63	6432(13)	825(6)	5539(14)	7.3(5)
C64	6926(12)	606(6)	6272(12)	5.9(5)
C65	7808(12)	562(6)	6576(12)	5.7(5)
C66	8221(10)	721(5)	6092(11)	4.5(4)
C71	9245(9)	1062(5)	5174(9)	3.0(4)
C72	9897(12)	1367(6)	5464(11)	5.6(5)
C73	10795(12)	1257(6)	5844(12)	6.3(5)
C74	10997(12)	865(6)	5934(12)	5.6(5)
C75	10424(12)	553(6)	5657(12)	5.9(5)
C76	9535(10)	660(5)	5269(10)	4.1(4)
B2	9643(11)	3681(5)	2539(11)	2.2(4)
C81	10406(9)	3388(5)	2499(9)	3.1(4)
C82	10348(10)	2973(5)	2411(10)	4.0(4)
C83	11047(14)	2737(7)	2423(13)	7.5(6)
C84	11804(12)	2904(6)	2494(12)	5.9(5)
C85	11883(11)	3314(6)	2570(11)	5.4(5)
C86	11197(11)	3553(5)	2611(10)	4.2(4)
C91	8837(9)	3393(5)	2464(10)	3.0(4)
C92	8576(10)	3340(5)	3108(11)	4.7(4)
C93	7879(11)	3086(5)	3004(12)	5.0(5)
C94	7426(11)	2861(5)	2293(12)	4.4(4)
C95	7650(11)	2920(5)	1650(11)	4.8(4)
C96	8330(10)	3176(5)	1734(11)	3.8(4)



Atom	$x$	$y$	$z$	$B$
C101	10039(9)	3924(5)	3435(10)	3.0(4)
C102	10591(10)	3715(5)	4174(11)	4.3(4)
C103	10848(11)	3908(6)	4978(12)	5.7(5)
C104	10581(12)	4290(6)	5002(13)	6.2(5)
C105	10050(14)	4518(7)	4331(15)	7.8(6)
C106	9775(11)	4310(6)	3513(12)	5.1(5)
C111	9277(9)	4017(5)	1764(10)	2.7(4)
C112	9691(10)	4137(5)	1256(10)	3.5(4)
C113	9359(11)	4449(5)	666(11)	4.5(4)
C114	8628(10)	4653(5)	526(10)	3.8(4)
C115	8171(10)	4536(5)	971(11)	4.1(4)
C116	8492(10)	4216(5)	1556(10)	3.2(4)
N9	5818(18)	2289(9)	265(18)	14.7(9)
C200	5212(22)	2497(10)	-279(21)	12.2(9)
C201	4554(19)	2675(9)	-915(19)	12.2(9)
N10	-174(18)	2691(9)	-956(20)	16.0(9)
C203	340(26)	2825(12)	-271(28)	16.3(12)
C204	881(19)	2990(9)	389(20)	12.0(9)

\*  $U_{eq} \times 10^4$

$$U_{eq} = \frac{1}{3} \sum_i \sum_j [U_{ij}(a_i^* a_j^*)(\bar{a}_i \bar{a}_j)]$$

**Table 2.3.** Selected Distances and Angles.

Atom	Atom	Distance (Å)	
Ir1	Ir2	3.199(1)	
Ir1	C1	1.962(17)	
Ir1	C11	1.986(14)	
Ir1	C21	1.932(18)	
Ir1	C31	2.008(14)	
Ir2	C10	1.900(17)	
Ir2	C20	2.015(18)	
Ir2	C30	1.924(17)	
Ir2	C40	1.981(17)	
C1	N1	1.14(2)	
C11	N3	1.09(2)	
C21	N5	1.13(2)	
C31	N7	1.09(2)	
C10	N2	1.15(2)	
C20	N4	1.13(2)	
C30	N6	1.16(2)	
C40	N8	1.11(2)	
Atom	Atom	Atom	Angle(Deg.)
Ir1	C1	N1	177.6(15)
Ir1	C11	N3	176.2(14)
Ir1	C21	N5	179.2(16)
Ir1	C31	N7	177.4(13)
Ir2	C10	N2	178.8(15)
Ir2	C20	N4	175.3(16)
Ir2	C30	N6	178.9(15)
Ir2	C40	N8	173.6(16)

## **Spectroscopic Measurements**

All solvents were dried and degassed by standard methods.<sup>5,6</sup> Absorption spectra were recorded with a Cary 17 spectrophotometer. Spectra were obtained of solutions prepared on a high-vacuum line in a cell consisting of a 10 ml Pyrex bulb, a 1 cm pathlength quartz cell, and a Teflon vacuum valve.

## **Electronic Emission Spectroscopy**

Electronic emission spectra were measured using an emission spectrophotometer constructed at CalTech that has been described previously.<sup>7</sup> Spectra recorded at ambient temperatures were obtained for solutions that were prepared on a high-vacuum line in a cell consisting of a 10 ml Pyrex bulb, a 1 cm pathlength quartz cell, and a Teflon vacuum valve. For the measurements at 77 K, solutions were prepared in the same manner in a glass NMR tube attached to a Teflon vacuum valve. Solvent was bulb-to-bulb distilled into the cell or tube containing the compound from the appropriate solvent storage flask. For the 77 K emission measurements, the NMR tubes were held in a liquid-nitrogen-filled quartz finger dewar.

## **Emission Lifetime Measurements**

Emission lifetime measurements were conducted with a Nd:YAG pulsed laser system that has been described previously using 532 nm excitation.<sup>8</sup> The solutions for the measurements at ambient temperatures were prepared by the procedure used for the emission measurements. The solutions for 77 K lifetime measurements were prepared in NMR tubes by the same procedure. All emission intensity decays exhibited first-order kinetics over at least 3 half-lives. For the measurements at 77 K, the NMR tubes were held in a liquid-nitrogen-filled quartz finger dewar. The laser was operated at a low-power output to minimize local sample heating.

**Resonance Raman Spectroscopy**

Resonance Raman spectra were measured by Dr. Thomas Loehr at the Oregon Graduate Center and by Dr. R.F. Dallinger at Los Alamos National Laboratory. Time-resolved resonance Raman spectra were measured by Dr. R.F. Dallinger at Los Alamos National Laboratory.

**Additional Measurements**

Elemental analyses were obtained at the CalTech Analytical Laboratory.

## Results and Discussion

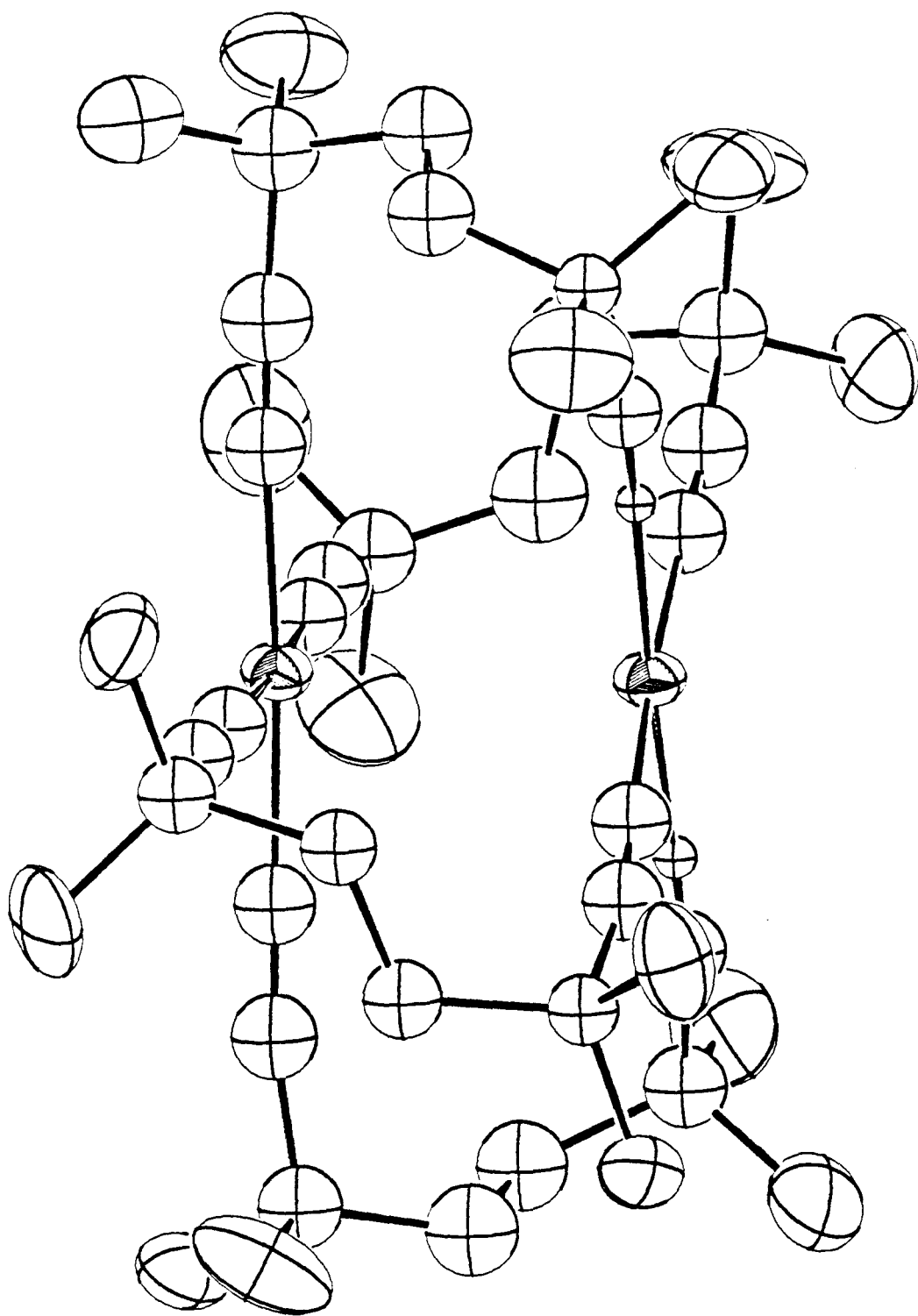
### Crystal Structure

The structure of  $\text{Ir}_2(\text{TMB})_4^{2+}$  is shown in Figure 2.1. The Ir-Ir distance of 3.119(1) Å is shorter than the metal-metal separation found for the analogous Rh(I) dimers ( $\text{Rh}_2(\text{CNC}_6\text{H}_5)_8^{2+}$ , 3.193(0) Å;  $\text{Rh}_2\text{b}_4^{2+}$ , 3.242(1) Å;  $\text{Rh}_2(\text{TMB})_4^{2+}$ , 3.262(1) Å).<sup>2,9</sup> The shorter metal-metal separation may be indicative of a stronger metal-metal interaction (metal-metal bond) in  $\text{Ir}_2(\text{TMB})_4^{2+}$  in comparison to the Rh complexes.

The most striking feature of the structure is the twist of the two  $\text{Ir}(\text{CN})_4$  units (Figure 2.2). The four TMB ligands are arranged as a four-bladed propeller about the Ir-Ir axis, giving the two  $\text{Ir}(\text{CN})_4$  units a partially staggered conformation. The torsional angle, N-M-M-N, is 28(4)° (Table 2.4), similar to the torsion angle observed for  $\text{Rh}_2(\text{TMB})_4^{2+}$ , 31°. The pitch of the blade is 27(3)° as measured from  $\text{C}_\beta\text{-Ir-Ir-C}_\beta$ . A best-fit plane of the  $\text{Ir}(\text{CN})_4$  units shows a slight distortion (Table 2.5). Each Ir experiences a small pyramidal distortion with the metal displaced toward the center of the cation. For  $\text{Rh}_2\text{b}_4^{2+}$ , a similar pyramidal distortion is observed with all four CN groups displaced toward the center of the dimer.<sup>2</sup> A tetrahedral distortion about the Rh atom is observed for  $\text{Rh}_2(\text{TMB})_4^{2+}$  and  $\text{Rh}_2(\text{CNC}_6\text{H}_5)_8^{2+}$ .<sup>2,9</sup> The average bond lengths and angle for the  $\text{Ir}(\text{CN})_4$  units are unexceptional and compare well with other Ir(I) and Rh(I) isocyanide complexes.<sup>2,10,11</sup> The overall structure of the TMB ligand is quite similar to that observed in  $\text{Rh}_2(\text{TMB})_4(\text{PF}_6)_2 \cdot 2\text{CH}_3\text{CN}$ .<sup>2</sup>

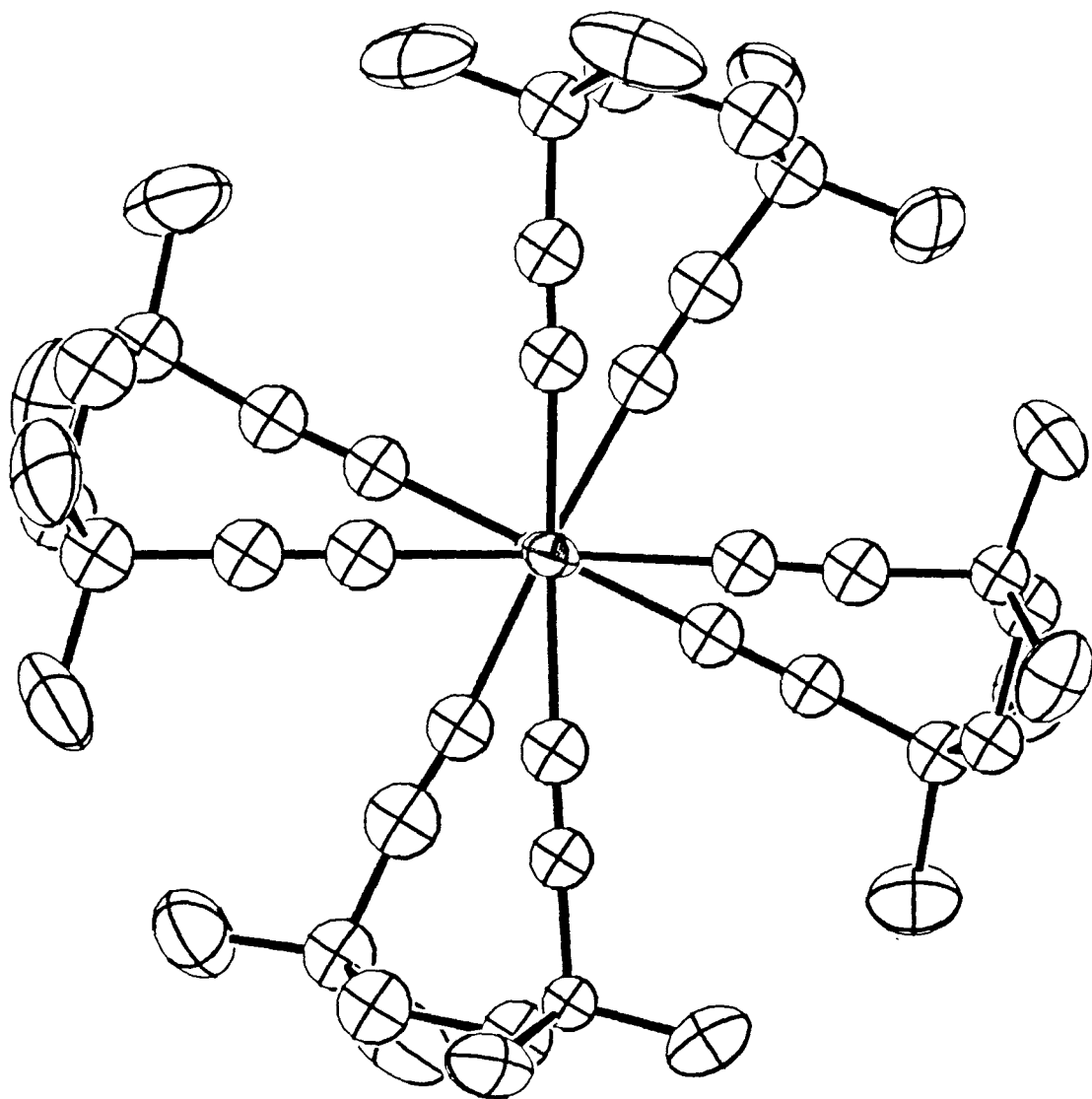
One of the interesting features of this structure is the arrangement of the cations in the unit cell. An ORTEP drawing of the unit cell, with the counter ions and solvent molecules removed for clarity, is given in Figure 2.3. The Ir-Ir vectors are aligned almost parallel to the b axis. Such an arrangement makes this crystal morphology ideally suited for polarized single-crystal spectroscopy.<sup>12,13</sup> However, the

**Figure 2.1.** View of the structure of  $\text{Ir}_2(\text{TMB})_4^{2+}$ .



**Figure 2.2.** View of the structure of  $\text{Ir}_2(\text{TMB})_4^{2+}$  viewed down the Ir-Ir vector.





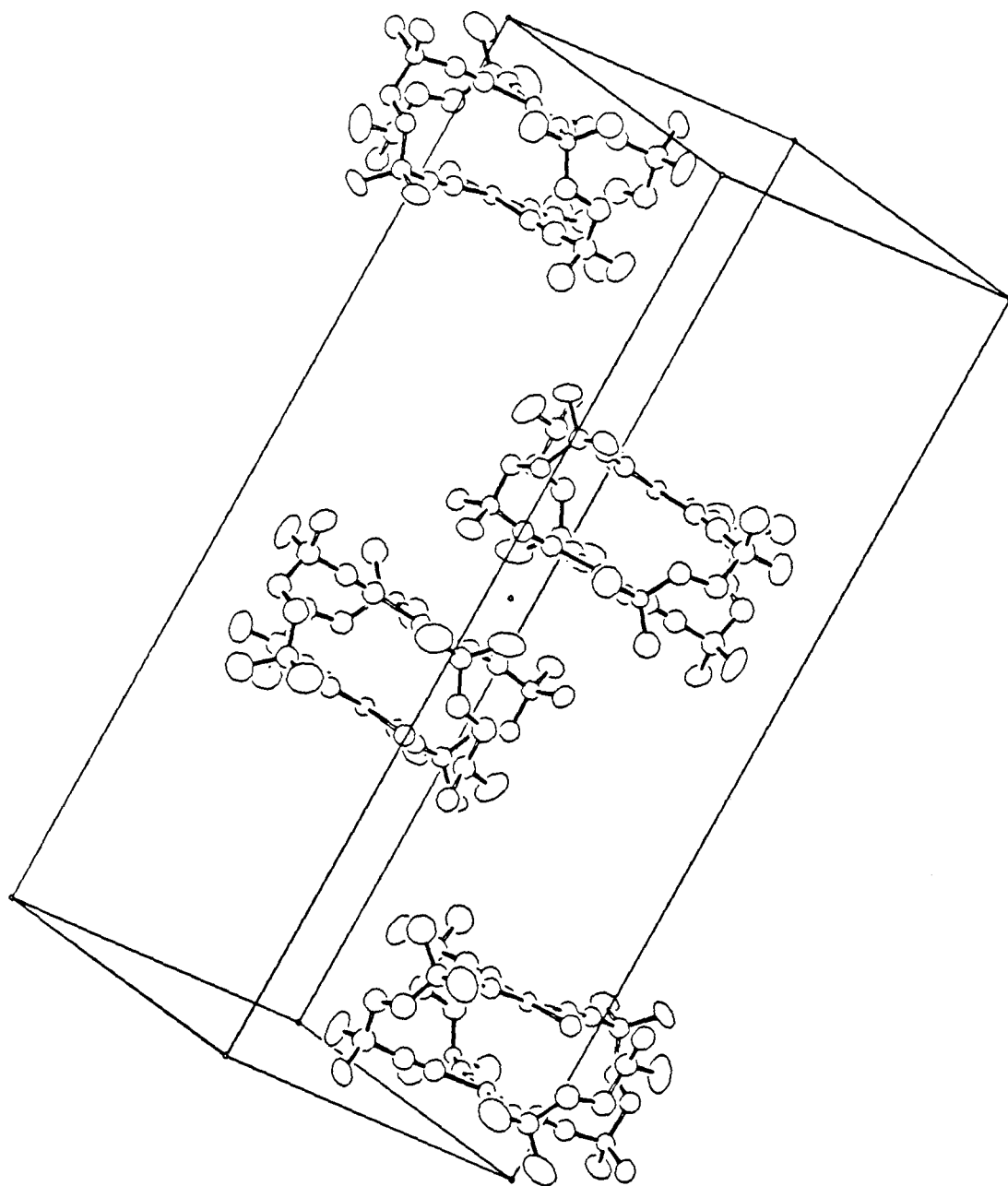
**Table 2.4.** Torsion Angles.

Atom	Atom	Atom	Atom	Angle (Deg.)
N1	lr1	lr2	N2	23.5(4)
N3	lr1	lr2	N4	31.6(4)
N5	lr1	lr2	N6	26.1(4)
N7	lr1	lr2	N8	29.7(4)
C2	lr1	lr2	C7	24.5(3)
C12	lr1	lr2	C17	32.6(3)
C22	lr1	lr2	C27	29.3(3)
C32	lr1	lr2	C37	27.3(3)

**Table 2.5.** Best plane fit of Ir(CN)<sub>4</sub> units.

Atom	Deviation (Å)	Atom	Deviation (Å)
Ir1	0.068	Ir2	-0.074
C1	0.004	C10	-0.005
N1	-0.055	N2	0.061
C11	0.018	C20	-0.033
N3	0.031	N4	-0.026
C21	0.001	C30	0.011
N5	-0.053	N6	0.048
C31	0.027	C40	-0.022
N7	0.028	N8	-0.034

**Figure 2.3.** View of the unit cell with  $\text{B}(\text{C}_6\text{H}_5)_4^-$  and  $\text{CH}_3\text{CN}$  molecules removed.



macroscopic features of the crystals (large chunky blue-black crystals, far too thick for spectroscopy) make such a study impractical.

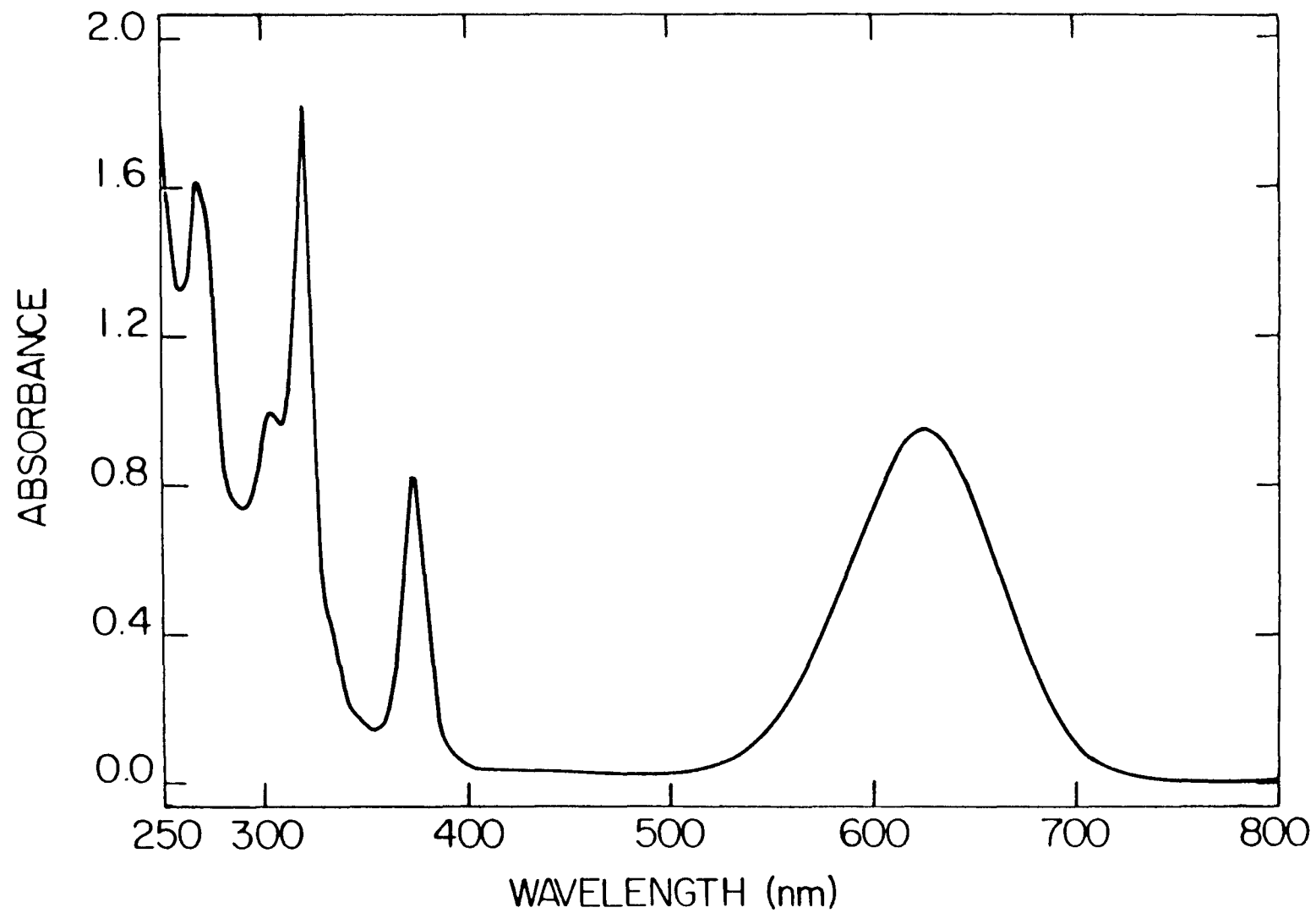
## Electronic Spectroscopy

A brief discussion of the spectroscopy and photophysics of  $\text{Ir}_2(\text{TMB})_4^{2+}$  has already been presented.<sup>1</sup> A detailed analysis of the optical absorption spectrum of  $\text{Ir}_2(\text{TMB})_4^{2+}$  is presented in Chapter 3.

The electronic absorption spectrum of  $\text{Ir}_2(\text{TMB})_4^{2+}$  is shown in Figure 2.4. In analogy to  $\text{Rh}_2(\text{TMB})_4^{2+}$ ,<sup>14</sup> the intense absorption centered at 625 nm ( $\epsilon = 11200 \text{ M}^{-1} \text{ cm}^{-1}$ ) is assigned to the  $^1(d\sigma^* \rightarrow p\sigma)$  transition ( $^1A_{1g} \rightarrow ^1A_{2u}$ ). The intense features below 300 nm, having maxima at 318 nm ( $\epsilon = 22000 \text{ M}^{-1} \text{ cm}^{-1}$ ) and 372 nm ( $\epsilon = 9750 \text{ M}^{-1} \text{ cm}^{-1}$ ), are assigned to  $d_{xz,yz} \rightarrow p_z$  transitions ( $^1A_{1g} \rightarrow ^1E_{1u}, ^3E_{1u}$ ). To lower energy of the 625 nm band in both single crystal and concentrated solution, a feature centered at 820 nm ( $\epsilon = 150 \text{ M}^{-1} \text{ cm}^{-1}$ ) is observed (Figure 2.5). This feature is assigned to the  $^3(d\sigma^* \rightarrow p\sigma)$  transition.

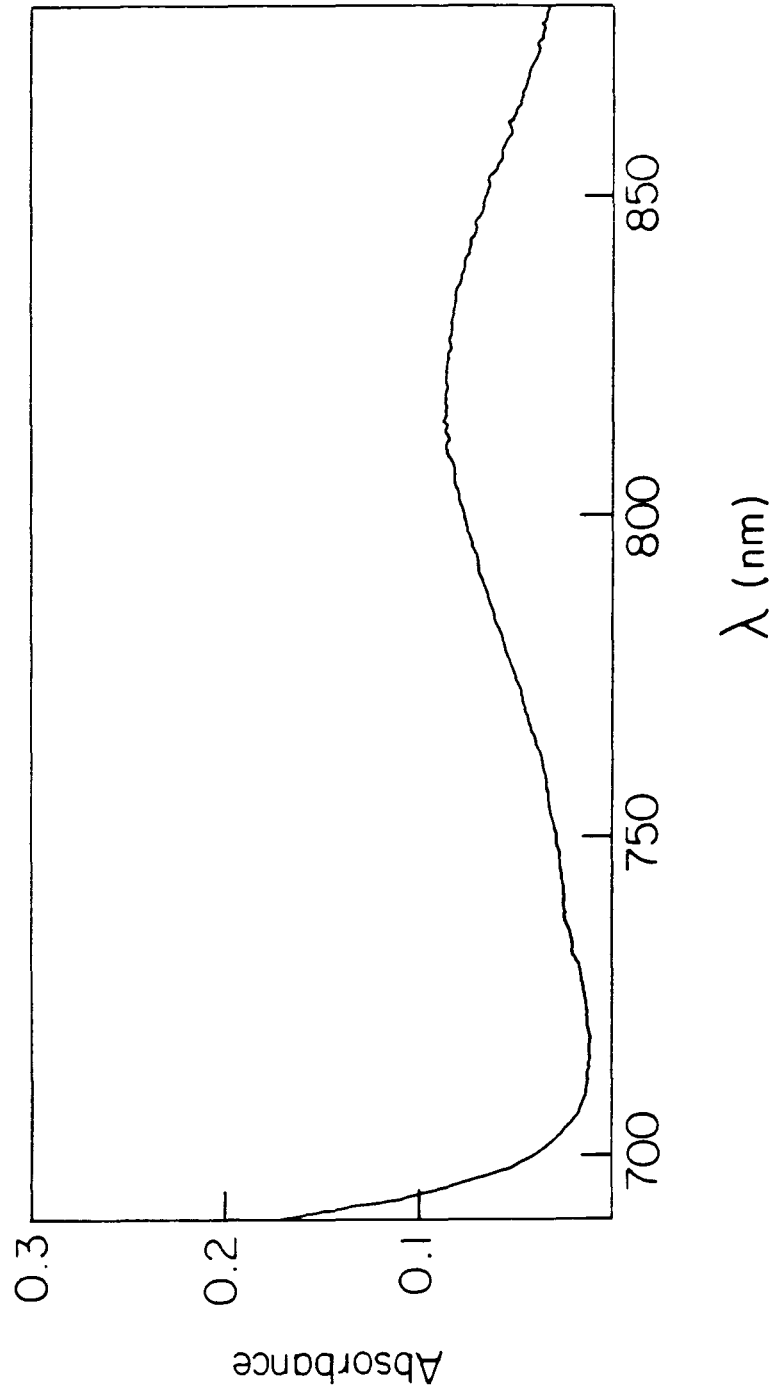
The electronic absorption spectrum of  $\text{Ir}_2(\text{TMB})_4^{2+}$  is summarized and compared to the spectra of  $\text{Ir}(\text{CN-}t\text{-butyl})_4^+$  and  $\text{Rh}_2(\text{TMB})_4^{2+}$  in Table 2.6.<sup>15</sup> The large red shift of the  $d_z^2 \rightarrow p_z$  transition in going to the dimer has been understood to be a consequence of the smaller d-p gap inferred from the MO model. The greater width of the  $A_{2u}(^1A_{2u})$  and  $E_u(^3A_{2u})$  bands is a result of the greater metal-metal bonding in the excited state. Promotion of an electron from the  $d\sigma^*$  orbital to the  $p\sigma$  orbital results in a large molecular distortion, contraction along the metal-metal coordinate. While the  $^1(d\sigma^* \rightarrow p\sigma)$  transition is shifted to the red of the monomer  $d_z^2 \rightarrow p_z$  transition, the  $d\pi \rightarrow p\sigma$  transitions of the dimer have energies and widths almost coincident with the monomer transitions; this observation is general to  $d^8$ - $d^8$  complexes. Explanations for the dichotomy of transition types have been presented but none is very satisfying.<sup>14</sup> The VB model, as discussed in Chapter 3, does accommodate both transition types. In comparison

**Figure 2.4.** Electronic absorption spectrum of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in  $\text{CH}_3\text{CN}$  at 25 °C.





**Figure 2.5.** Electronic absorption spectrum of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in a glassy matrix of  $\text{CH}_3\text{CH}_2\text{CN}/2\text{-MeTHF}$  (1:2 volume ratio) at 77 K.



**Table 2.6.** Absorption maxima for  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$ ,  $[\text{Ir}(\text{CN-}t\text{-butyl})_4]\text{BF}_4$ , and  $[\text{Rh}_2(\text{TMB})_4](\text{PF}_6)_2$  in  $\text{CH}_3\text{CN}$  at 25 °C.

$^1\text{A}_{1g} \rightarrow$	$\text{Ir}_2(\text{TMB})_4^{2+}$	$\text{Ir}(\text{CN-}t\text{-butyl})_4^+$	$\text{Rh}_2(\text{TMB})_4^{2+}$
$^3\text{A}_{2u}$	820 nm	489 nm	a
$^1\text{A}_{2u}$	625	423	520
$^3\text{E}_u(\text{E}_u)$	372	373	342
$^3\text{E}_u(\text{A}_{2u})$	333	a	a
$^1\text{E}_u$	318	308	316
$^3\text{B}_{1u}$	305	291	290

a. Not observed.

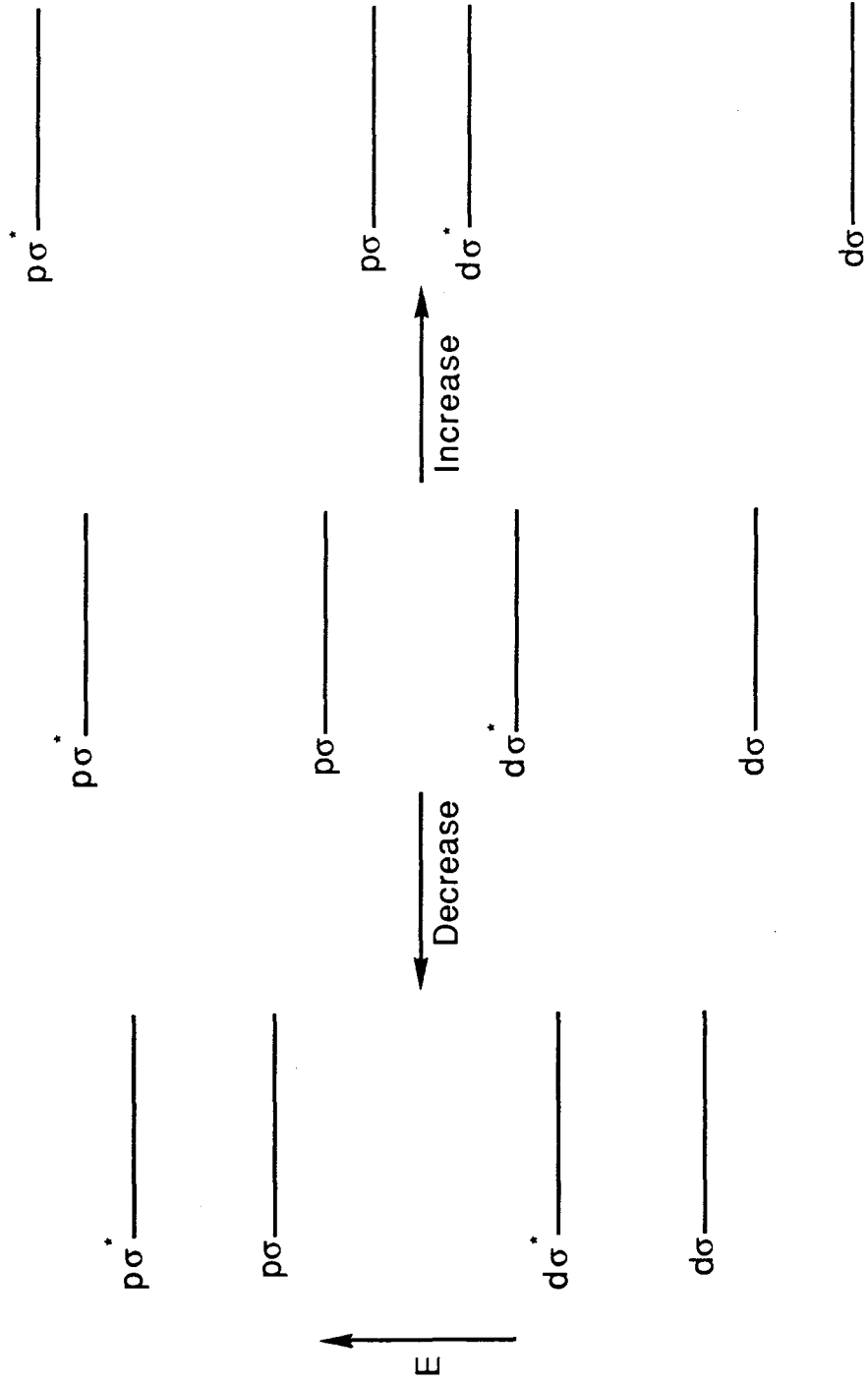
to  $\text{Rh}_2(\text{TMB})_4^{2+}$ , the  $^1(d\sigma^* \rightarrow p\sigma)$  and  $^3(d\pi \rightarrow p\sigma)$  for  $\text{Ir}_2(\text{TMB})_4^{2+}$  are red-shifted some  $3000\text{ cm}^{-1}$ . The widths of all transitions are quite comparable, indicating that a similar molecular distortion is occurring for both complexes upon excitation.

Earlier works have correlated the energy of the d-p transition for binuclear  $d^8$  complexes with the ground-state metal-metal separation.<sup>16</sup> Within the MO formalism, variation in the metal-metal separation will yield variation in the metal-metal interaction in the ground state. An increase in the metal-metal interaction in the ground state will serve to decrease the d-p gap, thereby shifting the  $d\sigma^* \rightarrow p\sigma$  transition to the red. A blue shift in the transition energy is predicted for a decrease in the metal-metal interaction (Figure 2.6). This analysis has been extended to include variations in the metal.<sup>16</sup> A change in metal, Rh to Ir, generally yields a large red shift in the  $^1(d\sigma^* \rightarrow p\sigma)$  transition. The variation in transition energy with metal has been presented as another example of the shift in energy of the  $d_z^2 \rightarrow p_z$  transition with change in metal-metal interaction, with the metal-metal interaction increased upon going from Rh to Ir.

The premise of the arguments for the shift in the  $^1(d\sigma^* \rightarrow p\sigma)$  transition energy is that the ground-state metal-metal interaction, as described by the MO model, controls the electronic transitions; that is, the excited states are well described within the ground-state MO formalism.

As will be discussed in detail in Chapter 3, the electronic structure of binuclear  $d^8$  complexes is more accurately described by a VB "weak-coupling" model. In this model, the ground state and various excited states are best described as weakly coupled monomers with the excitation localized on one center. For some states, it will be advantageous to increase the metal-metal one-electron interaction (shorten the metal-metal distance); this shortening of the metal-metal distance will result in an increased mixing of the covalent state with the higher-energy ionic state. This mixing will result in a lowering of the state energy relative to the weak coupling limit. Thus, unlike the ground-state MO model, the transition energies within the VB model are controlled by

**Figure 2.6.** MO diagram showing the effects on the energy of the  $d\sigma$  and  $p\sigma$  levels for a binuclear complex with an increase or a decrease in the metal-metal interaction.



excited-state properties, the amount of stabilization gained by increasing the metal-metal interaction.

Table 2.7 lists the  $^1(d\sigma^* \rightarrow p\sigma)$  transition energies for several binuclear isocyanide complexes. The crystallographically determined metal-metal separations are tabulated along with the  $^1(d_z^2 \rightarrow p_z)$  transition energies for the appropriate monomer model. For the series  $Rh_2(TMB)_4^{2+}$ ,  $Rh_2b_4^{2+}$ , and  $Rh_2(CNC_6H_5)_8^{2+}$ , a decrease in the metal-metal separation and coincident bathochromic shift of the  $d\sigma^* \rightarrow p\sigma$  transition is observed. However, when considering the shift from the corresponding  $d_z^2 \rightarrow p_z$  monomer excitation, a comparable red shift is observed, a varying reference point. For the series of Rh(I) arylisocyanide complexes ( $L = CNC_6H_5$ ,  $CN-p-C_6H_4F$ , and  $CN-p-C_6H_4NO_2$ ), a systematic increase in the Rh-Rh distance is observed with a systematic **red shift** of the transition energy. A result contrary to the earlier reasoning, an increase in the metal-metal separation yields a decrease in metal-metal interaction and results in a blue shift of the  $^1(d\sigma^* \rightarrow p\sigma)$  transition. Comparing  $[Rh(2,6-diMe-4-BrC_6H_2NC)_3Cl]_2^{2+}$  to  $Rh_2(CNC_6H_5)_8^{2+}$ , a blue shift of only  $640\text{ cm}^{-1}$  is observed for the  $^1(d\sigma^* \rightarrow p\sigma)$  transition for an increase of  $0.258\text{ \AA}$  in the Rh-Rh separation, a surprisingly small change in energy for such a large variation in metal-metal distance. If one, instead, compares the energy of the  $d\sigma^* \rightarrow p\sigma$  excitation to the monomer  $d_z^2 \rightarrow p_z$  transition, a shift of  $\sim 6500\text{ cm}^{-1}$  is observed for all Rh(I) isocyanide complexes except  $Rh_2b_4^{2+}$ . Thus, comparisons of transition energies for a series of metal dimers can be misleading in that the absolute energies of the transitions may not be comparable. The true reference point for each complex (the monomer excitation) may vary through the series. For the binuclear Rh(I) complexes, while the metal-metal separation varies dramatically in the ground state, the shift in the energy of the  $d_z^2 \rightarrow p_z$  transition is quite constant. The amount of stabilization in the  $^1(d\sigma^*p\sigma)$  excited state is comparable through the series of Rh(I) complexes. For those systems for which Rh(II)-Rh(II) species are known (Table 2.8), it is observed that a comparable, short metal-metal

**Table 2.7.**  $^1(d_z^2 \rightarrow p_z)$  transition energies for  $d^8$  complexes.

Metal Complex	$d_z^2 \rightarrow p_z$ , nm <sup>a</sup>	$\Delta E$ , cm <sup>-1</sup> <sup>b</sup>	M-M, Å <sup>c</sup>	Ref.
Rh(CN-ethyl) <sub>4</sub> <sup>+</sup>	380			17
Rh <sub>2</sub> b <sub>4</sub> <sup>2+</sup>	553	8200	3.242(1)	2,18
Rh(CN- <i>t</i> -butyl) <sub>4</sub> <sup>+</sup>	386			9
Rh <sub>2</sub> (TMB) <sub>4</sub> <sup>2+</sup>	520	6700	3.262(1)	2
Ir(CN- <i>t</i> -butyl) <sub>4</sub> <sup>+</sup>	423			15
Ir <sub>2</sub> (TMB) <sub>4</sub> <sup>2+</sup>	625	7600	3.199(1)	
Rh(CNC <sub>6</sub> H <sub>5</sub> ) <sub>4</sub> <sup>+</sup>	411			9
Rh <sub>2</sub> (CNC <sub>6</sub> H <sub>5</sub> ) <sub>8</sub> <sup>2+</sup>	568	6700	3.193(0)	9
Rh(CN- <i>p</i> -FC <sub>6</sub> H <sub>4</sub> ) <sub>4</sub> <sup>+</sup>	d			
[Rh(CN- <i>p</i> -FC <sub>6</sub> H <sub>4</sub> ) <sub>4</sub> ] <sub>2</sub> <sup>2+</sup>	597		3.207(2)	19
Rh(CN- <i>p</i> -NO <sub>2</sub> C <sub>6</sub> H <sub>4</sub> ) <sub>4</sub> <sup>+</sup>	d			
[Rh(CN- <i>p</i> -NO <sub>2</sub> C <sub>6</sub> H <sub>4</sub> ) <sub>4</sub> ] <sub>2</sub> <sup>2+</sup>	606		3.25(1)	19
Rh(2,6-diMe-4-BrC <sub>6</sub> H <sub>2</sub> NC) <sub>3</sub> Cl <sup>+</sup>	408			20
[Rh(2,6-diMe-4-BrC <sub>6</sub> H <sub>2</sub> NC) <sub>3</sub> Cl] <sub>2</sub> <sup>2+</sup>	548	6300	3.451(2)	20

a. For mononuclear  $d^8$  complexes,  $\lambda_{\max}$  for the  $^1(d_z^2 \rightarrow p_z)$  transition. For binuclear  $d^8$  complexes,  $\lambda_{\max}$  for the  $^1(d\sigma^* \rightarrow p\sigma)$  transition.

b. Red shift of the  $^1(d_z^2 \rightarrow p_z)$  transition in going from monomer to dimer.

$$\Delta E = E(^1(d_z^2 \rightarrow p_z)) - E(^1(d\sigma^* \rightarrow p\sigma)).$$

c. Crystallographically determined metal-metal separation.

d. Not known.



**Table 2.8.** Crystallographically determined metal-metal distances for Rh(I) and Rh(II) complexes.

d <sup>8</sup> -d <sup>8</sup> complex	M-M, Å	Ref.	d <sup>7</sup> -d <sup>7</sup> complex	M-M, Å	Ref.
Rh <sub>2</sub> b <sub>4</sub> <sup>2+</sup>	3.242	2	Rh <sub>2</sub> b <sub>4</sub> Cl <sub>2</sub> <sup>2+</sup>	2.837	2 1
Rh <sub>2</sub> (TMB) <sub>4</sub> <sup>2+</sup>	3.262	2	Rh <sub>2</sub> (TMB) <sub>4</sub> Cl <sub>2</sub> <sup>2+</sup>	2.770	2 2
Rh <sub>2</sub> (CNC <sub>6</sub> H <sub>5</sub> ) <sub>8</sub> <sup>2+</sup>	3.193	8			
Rh <sub>2</sub> (CN-p-CH <sub>3</sub> C <sub>6</sub> H <sub>4</sub> ) <sub>8</sub> <sup>2+</sup>	a		Rh <sub>2</sub> (CN-p-CH <sub>3</sub> C <sub>6</sub> H <sub>4</sub> ) <sub>8</sub> I <sub>2</sub> <sup>2+</sup>	2.785	2 3

a. Not known.

separation can be accommodated for all the complexes. This observation implies that the ground-state potential surface is quite soft and can accommodate a large range of metal-metal separations without a significant change in energy. The excited state surface, possessing a formal, metal-metal single bond, allows for only a small variation in metal-metal separation.

For  $\text{Ir}_2(\text{TMB})_4^{2+}$ , the  $^1(d\sigma^* \rightarrow p\sigma)$  transition is to lower energy of the transition for  $\text{Rh}_2(\text{TMB})_4^{2+}$ . However, the  $d_z^2 \rightarrow p_z$  transition for the Ir(I) monomer is also to lower energy of the Rh(I) monomer transition, a shift in reference point. In comparing the monomer-to-dimer shifts, the bathochromic shift observed for Ir(I) is larger than for Rh(I),  $7600\text{ cm}^{-1}$  compared to  $6700\text{ cm}^{-1}$ . The energy separation of the covalent and ionic states is less for Ir than for Rh, allowing for a greater mixing of the higher-energy states, yielding a greater stabilization of the  $^1(d\sigma^*p\sigma)$  excited state.

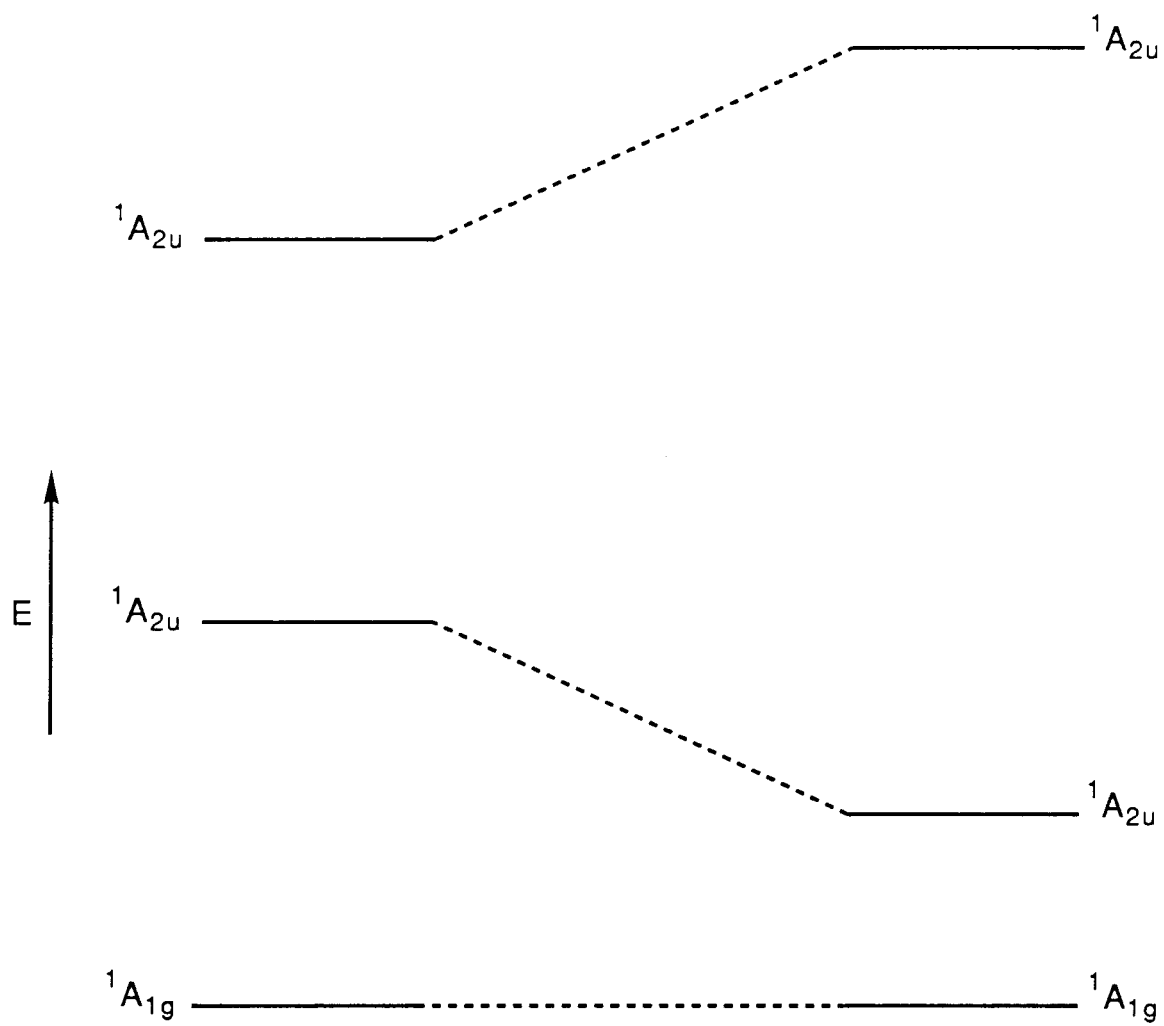
Within the series of complexes discussed,  $\text{Rh}_2\text{b}_4^{2+}$  appears to be an anomaly. The monomer-to-dimer shift of the  $d_z^2 \rightarrow p_z$  transition is much larger than that seen for analogous Rh(I) isocyanide complexes and is, in fact, larger than the shift observed for  $\text{Ir}_2(\text{TMB})_4^{2+}$ . This surprising result can be understood if one reasons that the bridge ligand destabilizes the  $d^8$ - $d^8$  ground state in favor of the contracted single-bonded Rh complex. The spectral band widths and Stokes shift observed for  $\text{Rh}_2\text{b}_4^{2+}$  are less than for  $\text{Rh}_2(\text{TMB})_4^{2+}$  and  $\text{Rh}_2(\text{CNC}_6\text{H}_5)_8^{2+}$ , indicating that the molecular distortion upon excitation is less.<sup>2</sup> This observation could be understood as a result of unfavorable ligand interactions at shorter Rh-Rh distances. The oxidation potential for  $\text{Rh}_2\text{b}_4^{2+}$  is estimated to be 500 mV less positive of  $\text{Rh}_2(\text{TMB})_4^{2+}$ , making the complex more easily oxidized than  $\text{Ir}_2(\text{TMB})_4^{2+}$ .<sup>24</sup> This result indicates that the  $\text{Rh}_2\text{b}_4^{2+}$  ground state is destabilized relative to the ground state of other binuclear  $d^8$  complexes. Therefore, the shift of the  $d_z^2 \rightarrow p_z$  transition for  $\text{Rh}_2\text{b}_4^{2+}$  is a combination of ground state and excited-state effects.

The important point of the above discussion is that the energy of the  $d\sigma^* \rightarrow p\sigma$  transition is governed by the ability of the excited state, rather than being controlled by ground-state effects, to undergo a contraction along the metal-metal coordinate. Another way of describing the phenomenon is that at the weak coupling limit, the excited state appears as an excited monomer coupled to a ground-state monomer. Increasing the metal-metal interaction serves to delocalize the excitation over the entire complex, lowering the energy of the state (Figure 2.7). This concept appears to be general to the large class of binuclear metal complexes. For the series of complexes  $Rh_2(CNR)_4(dppm)_2^{2+}$  (Table 2.9), the increased ability of the complex to accommodate a smaller Rh-Rh separation results in a systematic red shift of the  $d\sigma^* \rightarrow p\sigma$  transition energy (the monomer reference point does not vary with variation in the alkyl isocyanide). One final point that should be emphasized is that only the shift in the  $d_z^2 \rightarrow p_z$  transition energy between monomer and dimer is comparable within a series of complexes. Comparison of absolute transition energies for a series of complexes can lead to erroneous results.

## Photophysics

Excitation into the 625 nm band of  $Ir_2(TMB)_4^{2+}$  at room temperature in fluid solution results in two emissions (Figure 2.8). The higher-energy band centered at 735 nm is quite short-lived ( $< 2$  ns) and is assigned to the fluorescence  $^1A_{2u} \rightarrow ^1A_{1g}$ . The second band centered at 1080 nm is assigned to the phosphorescence  $^3A_{2u} \rightarrow ^1A_{1g}$  with a temperature-independent lifetime of  $210 \pm 20$  ns. No vibronic structure has been resolved in emission or absorption even at very low temperatures; therefore, the  $\lambda_{00}$  transition for the  $^1,^3A_{2u}$  states can be estimated only from either the overlap of the absorption and the emission spectra or by averaging the absorption and the emission maxima.<sup>32</sup> Either estimate places the  $\lambda_{00}$  transition for the  $^3A_{2u}$  state at 940 nm and at

**Figure 2.7.** Qualitative valence bond state diagram for weak d/p interaction of two  $d^8$  monomers. The energies of the  $^1A_{2u}$  states are dependent on the amount of stabilization, contraction along the metal-metal coordinate.



**Table 2.9.**  $^1(d_z^2 \rightarrow p_z)$  transition energies for  $d^8$  complexes.

Metal Complex	$d_z^2 \rightarrow p_z$ , nm <sup>a</sup>	$\Delta E$ , cm <sup>-1</sup> <sup>b</sup>	M-M, Å <sup>c</sup>	Ref.
Rh(CO)Cl(P(C <sub>6</sub> H <sub>5</sub> ) <sub>3</sub> )	367			30
Rh <sub>2</sub> (CO) <sub>2</sub> Cl <sub>2</sub> (dppm) <sub>2</sub>	442	4600	3.2386(5)	26, 27
Rh(CO)Cl(As(C <sub>6</sub> H <sub>5</sub> ) <sub>3</sub> )	356			25
Rh <sub>2</sub> (CO) <sub>2</sub> Cl <sub>2</sub> (dpam) <sub>2</sub>	466	6600	3.396(1)	26, 27
Rh(CNCH <sub>3</sub> ) <sub>2</sub> (P(C <sub>6</sub> H <sub>5</sub> ) <sub>3</sub> ) <sup>+</sup>	458			29
Rh <sub>2</sub> (CN-t-butyl) <sub>4</sub> (dppm) <sub>2</sub> <sup>2+</sup>	523	2700	d	29
Rh <sub>2</sub> (CNC <sub>6</sub> H <sub>11</sub> ) <sub>4</sub> (dppm) <sub>2</sub> <sup>2+</sup>	550	3600	d	29
Rh <sub>2</sub> (CN-n-C <sub>4</sub> H <sub>9</sub> ) <sub>4</sub> (dppm) <sub>2</sub> <sup>2+</sup>	560	4000	d	29
Rh <sub>2</sub> (CNCH <sub>3</sub> ) <sub>4</sub> (dppm) <sub>2</sub> <sup>2+</sup>	573	4400	d	29
Rh <sub>2</sub> (b) <sub>2</sub> (dppm) <sub>2</sub> <sup>2+</sup>	625	5800	d	30
Rh <sub>2</sub> (CNC <sub>6</sub> H <sub>5</sub> ) <sub>4</sub> (dppm) <sub>2</sub> <sup>2+</sup>	606	e	d	31

a. For mononuclear  $d^8$  complexes,  $\lambda_{\max}$  for the  $^1(d_z^2 \rightarrow p_z)$  transition. For binuclear  $d^8$  complexes,  $\lambda_{\max}$  for the  $^1(d\sigma^* \rightarrow p\sigma)$  transition.

b. Red shift of the  $^1(d_z^2 \rightarrow p_z)$  transition in going from monomer to dimer.

$$\Delta E = E(^1(d_z^2 \rightarrow p_z)) - E(^1(d\sigma^* \rightarrow p\sigma)).$$

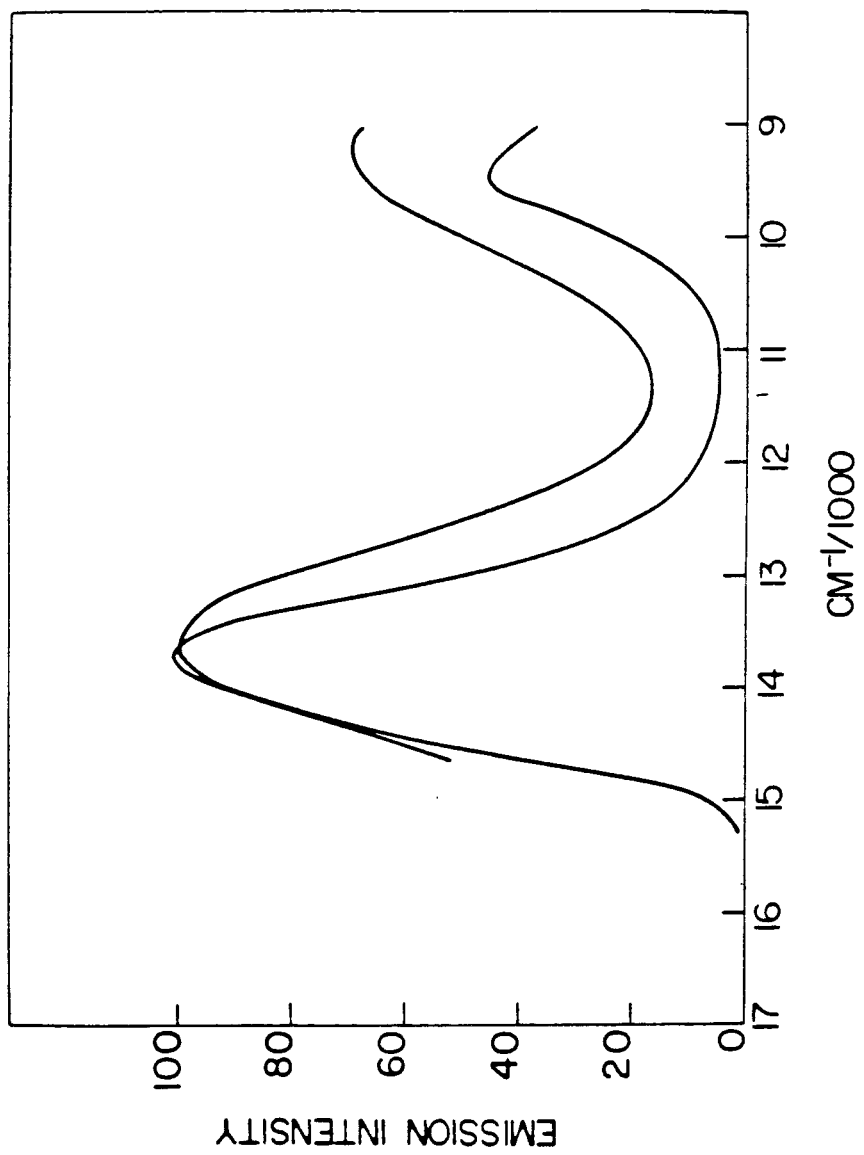
c. Crystallographically determined metal-metal separation.

d. Not known.

e. Not comparable to alkyl isocyanide complexes that are due to shift in reference point.

Spectral data for Rh(CNC<sub>6</sub>H<sub>5</sub>)<sub>2</sub>(P(C<sub>6</sub>H<sub>5</sub>)<sub>3</sub>)<sub>2</sub><sup>+</sup> not known.

**Figure 2.8.** Emission spectra of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in 2-MeTHF/Propionitrile at 298 K and 77 K. Taken from Reference 1.





680 nm for the  $^1A_{2u}$  state, 1.3 eV (30 kcal/mol) and 1.8 eV (42 kcal/mol) above the  $^1A_{1g}$  ground state, respectively.

Table 2.10 summarizes the photophysical data for  $\text{Ir}_2(\text{TMB})_4^{2+}$  and compares these data to the analogous  $\text{Rh}(\text{I})$  complexes. The Stokes shifts and emission band widths are comparable for the series of complexes, indicating a similar molecular distortion that accompanies the  $d\sigma^* \rightarrow p\sigma$  excitation. From the emission lifetime ( $\tau$ ) and quantum yield data ( $\Phi_{\text{em}}$ , emission quantum yield), the radiative rate ( $k_r$ ) and nonradiative rate ( $k_{nr}$ ) for the  $^1A_{2u}$  and  $^3A_{2u}$  states can be determined (Table 2.11). The nonradiative rate can be obtained, assuming that the nonradiative processes dominate the excited state deactivation.<sup>3,34,37,38</sup>

$$\tau = \frac{1}{(k_r + k_{nr})} \quad 2.1$$

$$\tau \approx \frac{1}{k_{nr}} \quad 2.2$$

From the measured quantum yields, the radiative rates can be estimated.

$$^1\Phi_{\text{em}} = \frac{^1k_r}{(^1k_r + ^1k_{nr})} \quad 2.3$$

$$^3\Phi_{\text{em}} = \Phi_{\text{isc}} \frac{^3k_r}{(^3k_r + ^3k_{nr})} \quad 2.4$$

For  $\text{Rh}_2\text{b}_4^{2+}$ , the  $\Phi_{\text{isc}}$  (intersystem crossing quantum yield) has been estimated to be 0.8.<sup>39</sup> For those systems that lack good quantum yield data, the  $^3A_{2u}$  state of  $\text{Rh}_2(\text{TMB})_4^{2+}$  and  $\text{Ir}_2(\text{TMB})_4^{2+}$ , the Strickler-Berg equation can be used to estimate  $k_r$ .<sup>37</sup>

$$k_r = 2.88 \times 10^{-9} (n^2)^{-3} \langle \nu^{-3} \rangle^{-1} (g_1/g_2) \int \epsilon d\nu \quad 2.5$$

$n$  - refractive index of the medium

Table 2.10. Photophysical data for  $\text{Ir}_2(\text{TMB})_4^{2+}$ ,  $\text{Rh}_2(\text{TMB})_4^{2+}$ , and  $\text{Rh}_2\text{b}_4^{2+}$ .

	$\lambda_{\text{max}}$	$\tau$		$\Phi_{\text{em}}$		Stokes shift	Ref.
	nm	298 K	77 K	298 K	77 K	$\text{cm}^{-1}$	
$\text{Ir}_2(\text{TMB})_4^{2+}$							1
$^1\text{A}_{2\text{u}}$	735	< 2 ns		$2.5 \times 10^{-3}$ a		2400	
$^3\text{A}_{2\text{u}}$	1080	$210 \pm 20$ ns	$200 \pm 10$ ns	$10^{-3}$ b		2900	
$\text{Rh}_2(\text{TMB})_4^{2+}$							2, 33, 34
$^1\text{A}_{2\text{u}}$	614	$820 \pm 20$ ps		$4.6 \times 10^{-2}$ c		3200	
$^3\text{A}_{2\text{u}}$	740	$25 \pm 5$ ns	20.5 $\mu\text{s}$	$6 \times 10^{-4}$ d	0.51 c		
$\text{Rh}_2\text{b}_4^{2+}$							2, 34, 35, 36
$^1\text{A}_{2\text{u}}$	656	1.3 ns		$5.6 \times 10^{-2}$ c	0.08 c	2800	
$^3\text{A}_{2\text{u}}$	810	$8.5 \pm 5$ $\mu\text{s}$	12.5 $\mu\text{s}$	0.32 c	0.64	2600	

a. Error in determination  $\pm 50\%$ .

b. Error in determination  $\pm 100\%$ .

c. Error in determination  $\pm 15\%$ .

d. Estimate,  $\Phi_{\text{em}}(298 \text{ K})/\Phi_{\text{em}}(77 \text{ K}) = \tau(298 \text{ K})/\tau(77 \text{ K})$ . Assumes  $\Phi_{\text{isc}}$  and  $k_{\text{r}}$  to be temperature independent.

**Table 2.11.** Radiative rates ( $k_r$ ) and nonradiative rates ( $k_{nr}$ ) for the  $^1A_{2u}$  and  $^3A_{2u}$  excited states of  $\text{Ir}_2(\text{TMB})_4^{2+}$ ,  $\text{Rh}_2(\text{TMB})_4^{2+}$ , and  $\text{Rh}_2\text{b}_4^{2+}$ .

	$^1A_{2u}$ excited state			
	$^1k_{nr}$		$^1k_r$	
	298 K	77 K	298 K	77 K
$\text{Ir}_2(\text{TMB})_4^{2+}$	$9.5 \times 10^9$		$2.3 \times 10^7$ a	
$\text{Rh}_2(\text{TMB})_4^{2+}$	$1.2 \times 10^9$	$3.7 \times 10^8$	$5.6 \times 10^7$	
			$4.8 \times 10^7$ a	
$\text{Rh}_2\text{b}_4^{2+}$	$7.7 \times 10^8$	$5.4 \times 10^8$	$4.3 \times 10^7$	
	$^3A_{2u}$ excited state			
	$^3k_{nr}$		$^3k_r$	
	298 K	77 K	298 K	77 K
$\text{Ir}_2(\text{TMB})_4^{2+}$	$4.8 \times 10^6$	$5.0 \times 10^6$	$4.8 \times 10^3$ b	
			$9 \times 10^4$ c	
$\text{Rh}_2(\text{TMB})_4^{2+}$	$4.0 \times 10^7$	$4.9 \times 10^4$	$3.0 \times 10^4$	$3.1 \times 10^4$
$\text{Rh}_2\text{b}_4^{2+}$	$1.2 \times 10^5$	$8.0 \times 10^4$	$4.7 \times 10^4$	$5.1 \times 10^4$

a. Calculated with Equation 2.6.

b. From the measure  $^3\Phi_{em}$ .

c. Calculated with Equation 2.6, using 77 K absorption data. This assumes that  $k_r$  is temperature-independent.

$\langle \nu^{-3} \rangle$  - mean of the cube of the reciprocal of the emission frequency

$g_1, g_2$  - the degeneracies of the lower and upper states

$\nu$  - the emission maximum

$$k_r \approx 5 \times 10^{-9} \nu^2 \epsilon_{\max} \Delta \nu \quad 2.6$$

$\epsilon_{\max}$  - extinction coefficient of the absorption band

$\Delta \nu$  - absorption band width (full width at half the maximum height)

The radiative rate for the series of complexes is found to be quite constant and temperature-independent. Similar radiative rates are expected for  $\text{Rh}_2\text{b}_4^{2+}$  and  $\text{Rh}_2(\text{TMB})_4^{2+}$  but is somewhat surprising for  $\text{Ir}_2(\text{TMB})_4^{2+}$ . One might have expected that the increased spin-orbit coupling for the third-row transition metal would have increased the radiative rate. This increase may be offset by the decrease in emission energy, which also accompanies the change in metal.

The nonradiative rates, unlike the radiative rates, vary through the series. The  $^1k_{nr}$  changes because of the varying flexibility of the ligand system and the decreased energy gap through the series. For the  $^3k_{nr}$ , the expected linear relationship between  $\ln k_{nr}(77 \text{ K})$  and the emission energy ( $E_{em}$ ) is observed.<sup>40</sup> At room temperature, the  $k_{nr}$  for  $\text{Rh}_2(\text{TMB})_4^{2+}$  increases dramatically, while for  $\text{Ir}_2(\text{TMB})_4^{2+}$  only a slight variation is observed. This marked temperature dependence of the nonradiative rate for  $\text{Rh}_2(\text{TMB})_4^{2+}$  has been studied previously.<sup>34,41</sup>

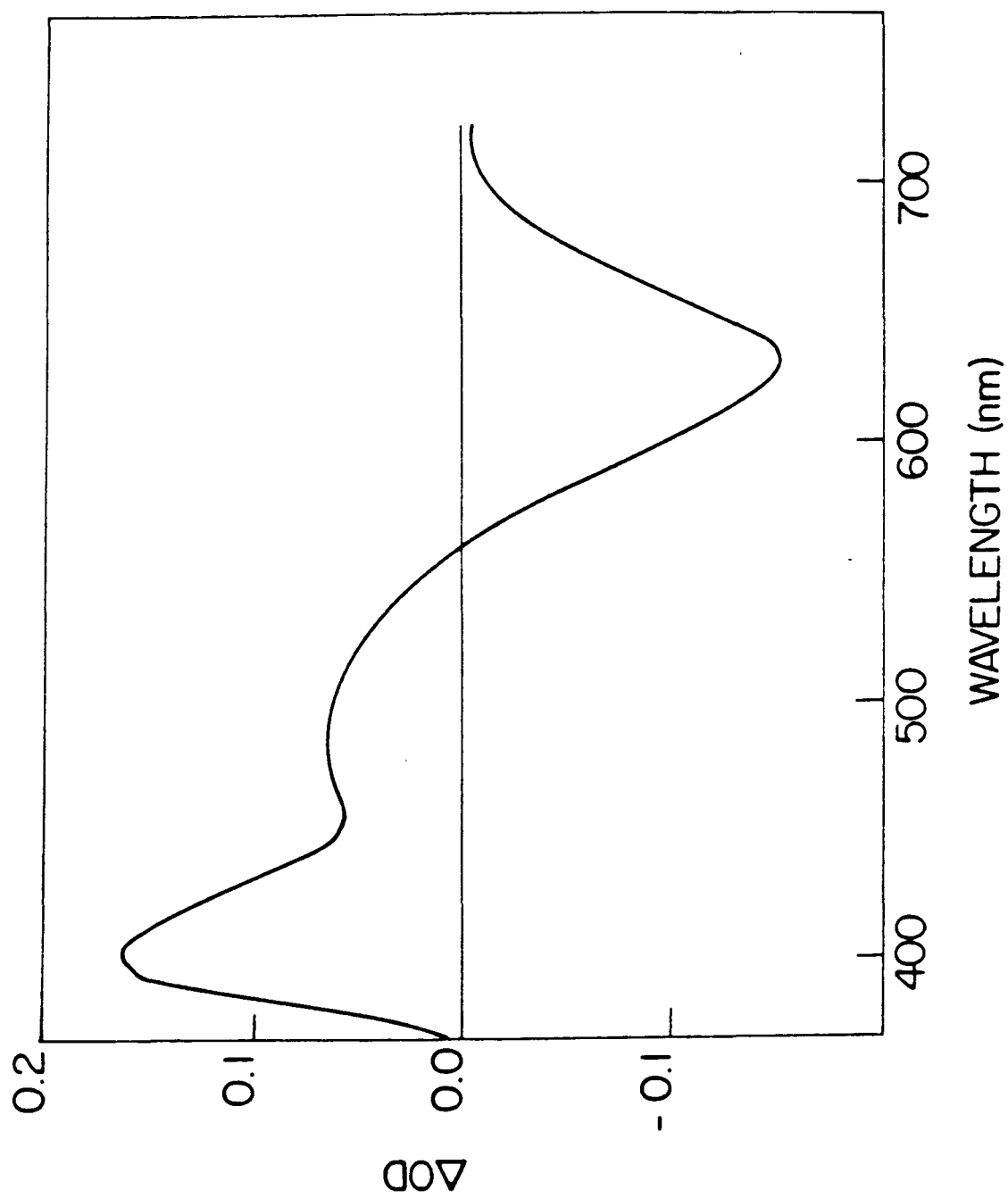
The earlier study of the temperature dependence of the  $^3A_{2u}$  state discussed the nonradiative rate constants for the binuclear  $d^8$  complexes within the framework of the strong-coupling model as described by Engleman and Jortner.<sup>34</sup> It was thought that the  $^3(d\sigma^*p\sigma)$  excited state is thermally deactivated to a  $dd$  excited state, whose equilibrium geometry is  $D_{2d}$  distorted about the metal center, and that such a state would undergo a rapid nonradiative decay to the ground state. Such  $D_{2d}$  distortions are observed for

certain dd excited states of  $d^8$  square planar complexes. The  $^3B_2$  and  $^3B_1$  ligand field states having the configuration  $(d_z^2)^1(d_{x^2-y^2})^1$  are states that could couple to the  $^3A_{2u}$  state. Ballhausen *et al.* have suggested that a similar configuration in  $d^8$  monomers is unstable with respect to a  $D_{2d}$  distortion, which would result in intersection with the ground-state surface and account for the effective intersystem crossing.<sup>42</sup> That the  $^3A_{2u}$  state of  $Ir_2(TMB)_4^{2+}$  does not show a temperature-dependent lifetime is support of this analysis, since the ligand field states are thermally inaccessible. A recent NMR structural study of Rh and Ir isocyanide complexes proposed that rotation about the metal-metal axis contributes substantially to the deactivation of the  $^3A_{2u}$  state, the rotation coupling the  $^3A_{2u}$  state to intermediate dd states.<sup>22</sup> Again, for  $Ir_2(TMB)_4^{2+}$ , such states would be higher in energy; therefore, rotation about the metal-metal axis would not deactivate the  $^3A_{2u}$  state.

There has been no direct measurement of the  $^1A_{2u}$  excited-state lifetime for  $Ir_2(TMB)_4^{2+}$ . Two estimates of the  $^1A_{2u}$  lifetime of  $Ir_2(TMB)_4^{2+}$  can be made. If one assumes the radiative rates for  $Rh_2(TMB)_4^{2+}$  and  $Ir_2(TMB)_4^{2+}$  are the same, from the ratio of the  $^1\Phi_{em}$ , an estimate of 40 ps ( $\tau = 1/k_{nr}$ ) is obtained. Using the calculated radiative rate,  $^1\tau$  is estimated to be 100 ps. A  $^1A_{2u}$  lifetime of  $70 \pm 30$  ps is in line with other binuclear  $Ir(I)$  complexes.<sup>3</sup>

The transient absorption spectrum of  $Ir_2(TMB)_4^{2+}$  has been measured previously (Figure 2.9).<sup>1</sup> The transient has a lifetime of  $185 \pm 15$  ns, in good agreement with the lifetime obtained by measuring the decay of the phosphorescence. The decrease in absorbance at 625 nm is due to bleaching of the ground state, with the bands at 400 and 510 nm attributed to absorptions by the  $^3A_2$  state. For  $Rh_2b_4^{2+}$  and  $Rh_2(TMB)_4^{2+}$ , these absorptions have been assigned to  $d\pi \rightarrow p\sigma$  and  $d\sigma \rightarrow d\sigma^*$ , respectively.<sup>36</sup>

**Figure 2.9.** Transient difference spectrum for  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in  $\text{CH}_3\text{CN}$  at 25 °C. Taken from Reference 1.



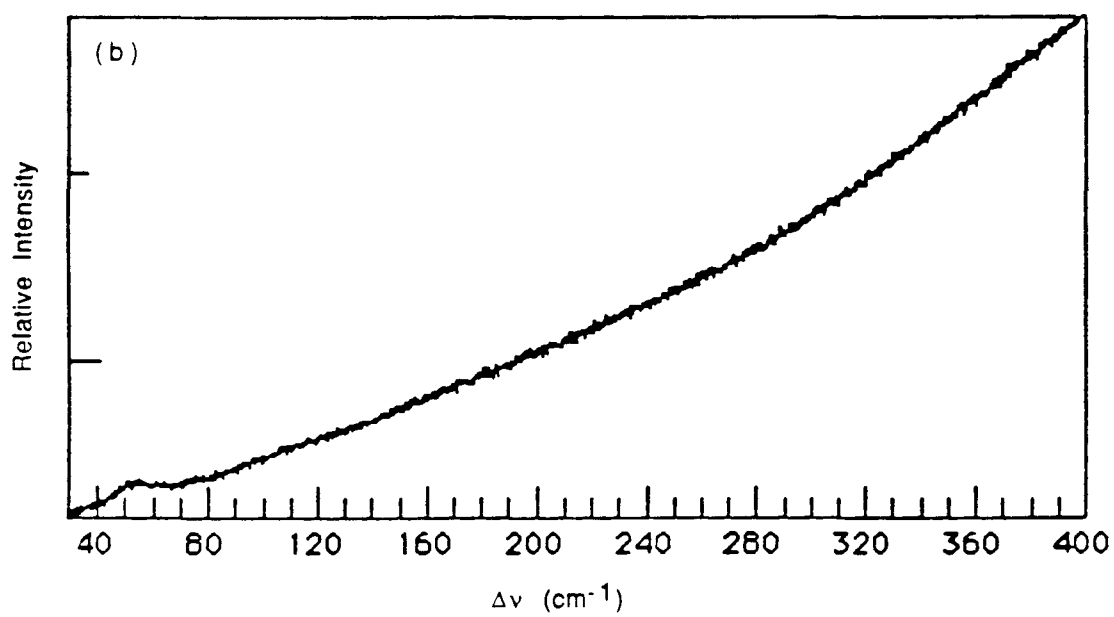
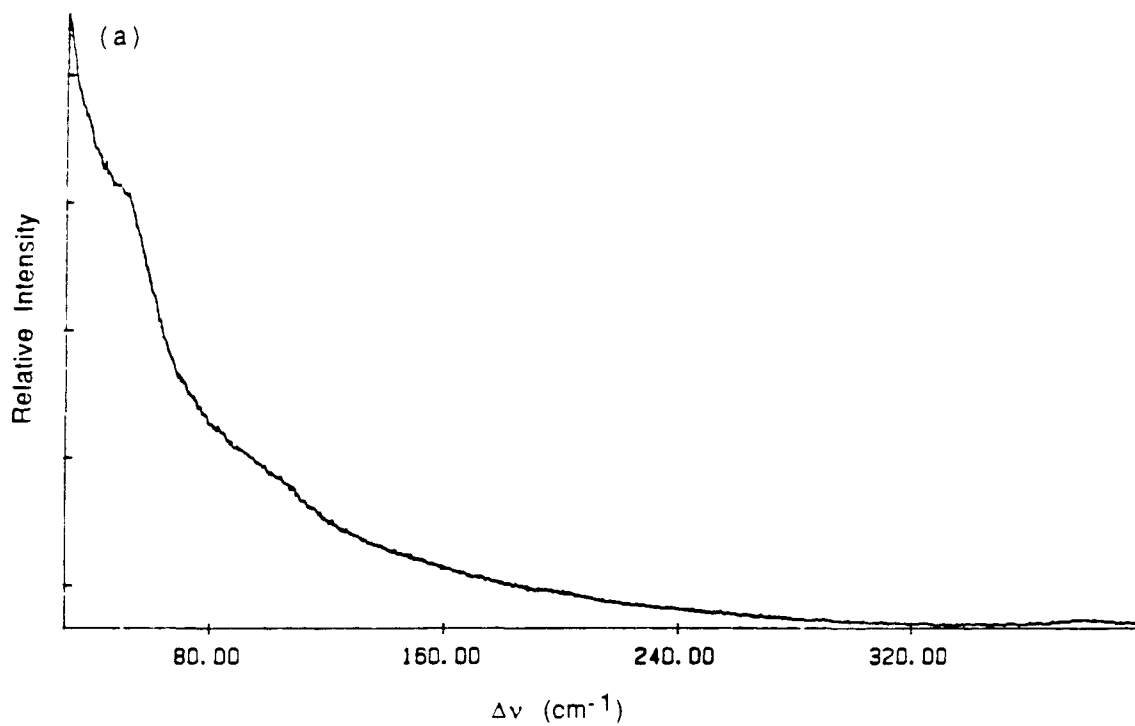
## Vibrational Spectroscopy

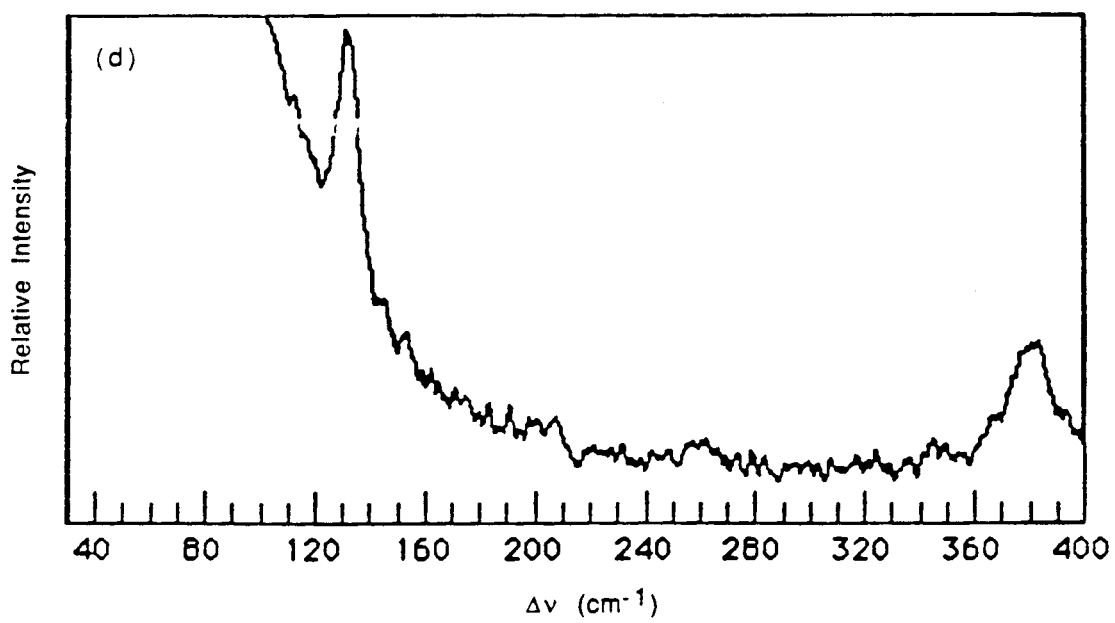
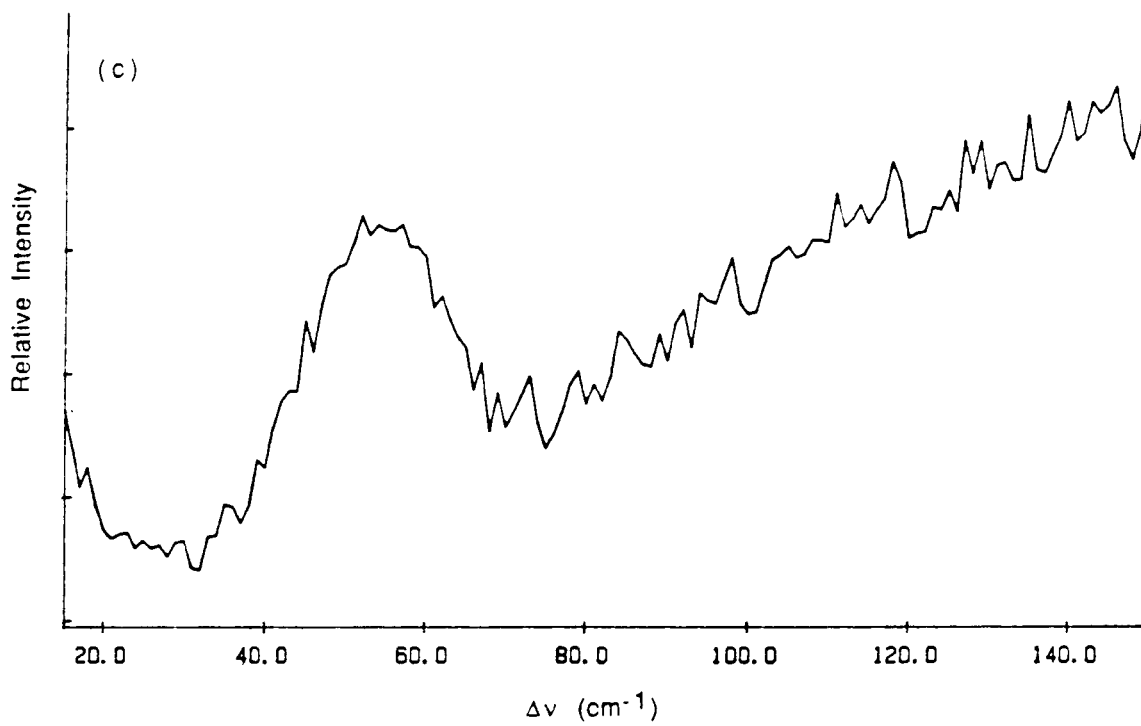
The ground state of the binuclear  $d^8$  complexes possesses only a weak metal-metal bonding interaction, with the  $d\sigma^*p\sigma$  excited states possessing a formal metal-metal single bond. The metal-metal interactions in the ground and excited states of  $\text{Ir}_2(\text{TMB})_4^{2+}$  have been probed using resonance Raman spectroscopy. The spectra are shown in Figure 2.10. The low-frequency features for some of the spectra are poorly resolved because of Rayleigh scattering of the laser line. Other spectra show a substantial base-line slope because of emission from the fluorescent excited state. The ground state spectrum shows an intense peak at  $53\text{ cm}^{-1}$  with a feature at  $106\text{ cm}^{-1}$ ,  $2\nu(\text{Ir-Ir})$ , which may be attributed to the first overtone of the  $53\text{ cm}^{-1}$  vibration. The  $^3A_2$  excited-state spectrum exhibits an intense peak at  $132\text{ cm}^{-1}$ , with possibly a weak overtone at  $264\text{ cm}^{-1}$ . These bands are assigned to the Ir-Ir stretching modes. For the excited-state spectrum, the assignment of the  $132\text{ cm}^{-1}$  band to the excited state Ir-Ir vibration is supported by the dependence of the spectra on the per-pulse laser power. As the incident-per-pulse laser energy is varied from 2 to 5 mJ/pulse, the  $132\text{ cm}^{-1}$  band increases in intensity as the  $53\text{ cm}^{-1}$  intensity decreases (Figures 2.10e and 2.10f). Similar per-pulse laser-power dependence has been observed for  $\text{Rh}_2\text{b}_4^{2+}$ .<sup>43</sup> Upon excitation to the  $^3A_{2u}$  excited state, there is a sixfold increase in the metal-metal restoring force (Table 2.12). This is in line with what is observed for other binuclear  $d^8$  complexes.

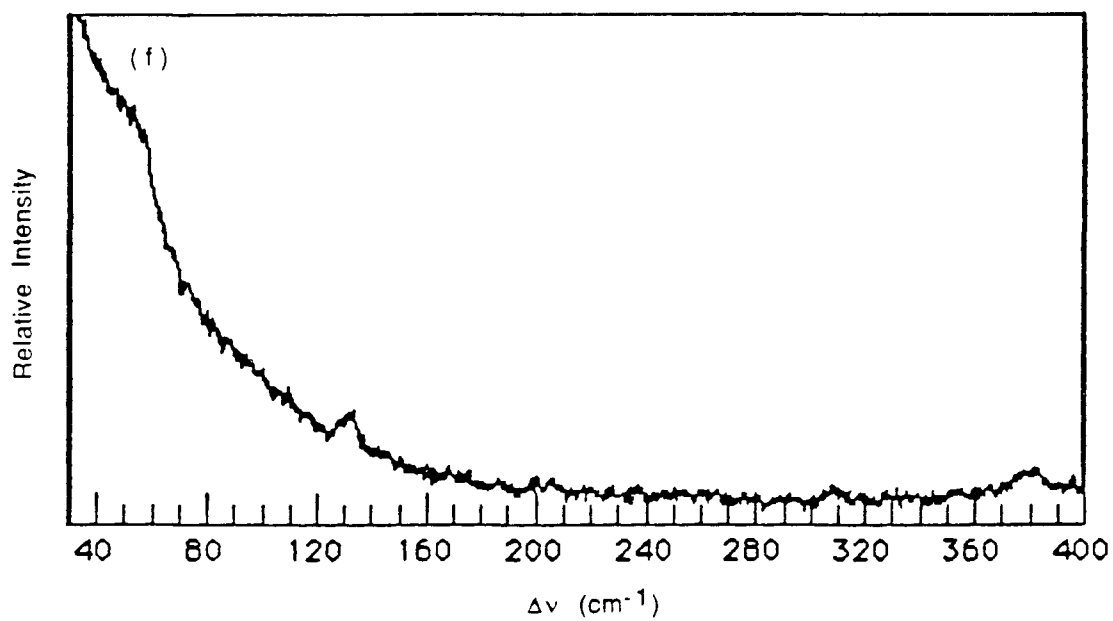
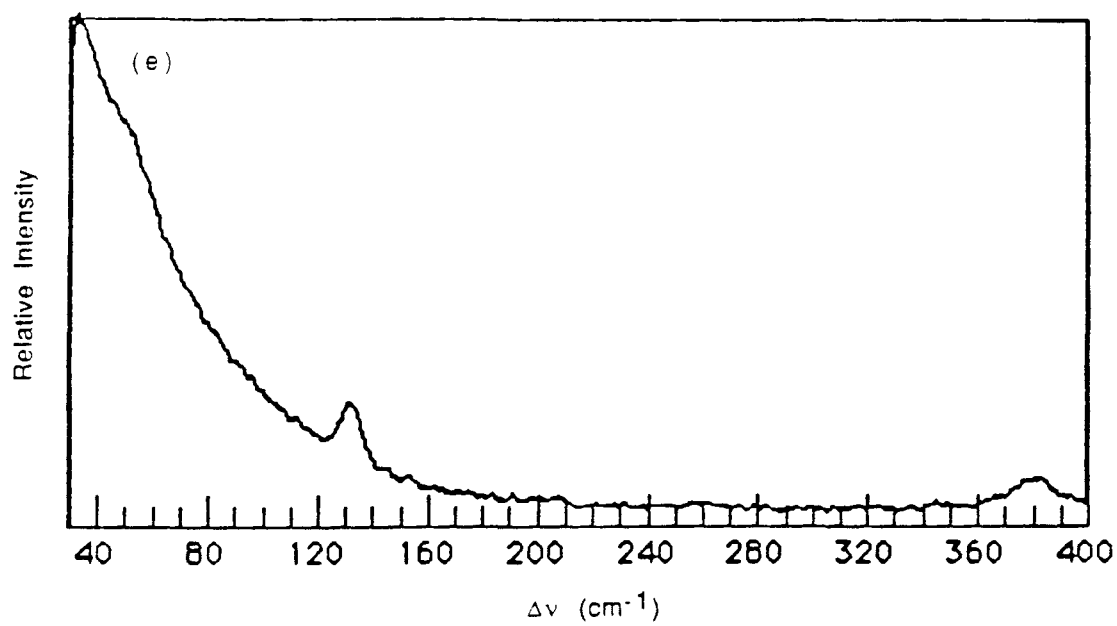
It was stated earlier that the metal-metal interaction in  $\text{Ir}_2(\text{TMB})_4^{2+}$  is greater (stronger metal-metal bonding) than in the analogous  $\text{Rh(I)}$  complex. This is manifest in the vibrational data. The metal-metal restoring force for the  $^1A_{1g}$  state increases by a factor of 1.9 in going from  $\text{Rh}_2(\text{TMB})_4^{2+}$  to  $\text{Ir}_2(\text{TMB})_4^{2+}$ . The metal-metal bond strength in the  $^3A_{2u}$  state of  $\text{Ir}_2(\text{TMB})_4^{2+}$  is also greater than in the  $\text{Rh(I)}$  dimers. The Rh-Rh bond strength in the  $^3A_{2u}$  state of  $\text{Rh}_2\text{b}_4^{2+}$  is estimated to be  $42\text{ kcal/mol}$ .<sup>14</sup> The shorter metal-metal bond length calculated for  $\text{Ir}_2(\text{TMB})_4^{2+}$  and larger metal-metal



**Figure 2.10.** Resonance Raman spectra of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$ : (a)  $\text{CH}_3\text{CN}$  solution,  $\lambda_{\text{ex}} = 568.2 \text{ nm}$ ; (b)  $\text{CH}_3\text{CN}$  solution,  $\lambda_{\text{ex}} = 647.1 \text{ nm}$ ; (c)  $\text{K}_2\text{SO}_4$  pellet,  $\lambda_{\text{ex}} = 647.1 \text{ nm}$ ; (d) Time resolved,  $\text{CH}_3\text{CN}$  solution,  $\lambda_{\text{ex}} = 565 \text{ nm}$ , 10 scans; (e) Time resolved,  $\text{CH}_3\text{CN}$  solution,  $\lambda_{\text{ex}} = 565 \text{ nm}$ , 10 scans, 2 mJ/P; (f) Time resolved,  $\lambda_{\text{ex}} = 565 \text{ nm}$ , 10 scans, 5 mJ/P.







**Table 2.12.** Observed metal-metal vibrational frequencies and calculated force constants. Observed metal-metal distances and calculated metal-metal distances using Woodruff's relationship.

	<sup>1</sup> A <sub>1g</sub>		<sup>3</sup> A <sub>2u</sub>		Ref.
	$\nu(\text{M-M})$	F	$\nu(\text{M-M})$	F	
	cm <sup>-1</sup>	mdyne/Å	cm <sup>-1</sup>	mdyne/Å	
Rh <sub>2</sub> b <sub>4</sub> <sup>2+</sup>	79	0.19	144	0.63	43
Rh <sub>2</sub> (TMB) <sub>4</sub> <sup>2+</sup>	55	0.09			44
Ir <sub>2</sub> (TMB) <sub>4</sub> <sup>2+</sup>	53	0.16	132	0.99	

	Metal-Metal Separation			
	<sup>1</sup> A <sub>1g</sub>		<sup>3</sup> A <sub>2u</sub>	
	X-ray, Å	W.R., Å <sup>a</sup>	W.R., Å <sup>a</sup>	$\Delta$ , Å <sup>b</sup>
Rh <sub>2</sub> b <sub>4</sub> <sup>2+</sup>	3.242(1)	3.17 <sup>c</sup>	2.96	0.28 <sup>d</sup>
Rh <sub>2</sub> (TMB) <sub>4</sub> <sup>2+</sup>	3.262(1)	3.23		0.4
Ir <sub>2</sub> (TMB) <sub>4</sub> <sup>2+</sup>	3.199(1)	3.23	2.87	0.33

a. Woodruff's relationship: Rh,  $D = 1.83 + 1.45\exp(-F/2.53)$ ; Ir,  $D = 2.01 + 1.31\exp(-F/2.36)$ .<sup>45</sup>

b. Decrease in metal-metal separation upon excitation,  $\Delta = \text{W.R.}(\text{}^1\text{A}_{1g}) - \text{W.R.}(\text{}^3\text{A}_{2u})$ .

c. The large discrepancy between the calculated and observed metal-metal distances is a consequence of ligand contribution to the observed vibrational frequency; using the crystallographically determined Rh-Rh distance  $F_{\text{calc}} = 0.07$ ,  $\nu_{\text{calc}} = 47 \text{ cm}^{-1}$

d. Frank-Condon analysis yields  $\Delta Q = 0.31 \text{ Å}$ .<sup>14</sup>

restoring force in the  $^3A_{2u}$  excited state imply a stronger metal-metal interaction.

While it is not necessarily the case that the increased force constant is an indication of a stronger bonding interaction, there is precedent in the literature for such a correlation within a series of analogous metal-metal bonded compounds.<sup>46,47</sup> Larger force constants for Ir in comparison to Rh have been observed for complexes that possess formal metal-metal single bonds.<sup>44</sup> The major distortion upon excitation is along the Ir-Ir coordinate; a contraction of 0.3 Å is expected.

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### **Chapter 3**

#### **Detailed Study of the Electronic Spectrum of $\text{Ir}_2(\text{TMB})_4^{2+}$**

## Introduction

The electronic spectra of face-to-face, metal-metal bonded  $d^8$ - $d^8$  complexes have been extensively discussed on the basis of a molecular orbital (MO) model.<sup>1-10</sup> In this picture, the metal-based  $d_{z^2}$  (filled) and  $p_z$  (empty) orbitals each interact strongly to produce pairs of bonding and antibonding orbitals. Configuration interaction between the empty and filled sets (which might be visualized alternatively as dative d-p bonding) provides the weak ground-state bond, while the lowest-energy singlet and triplet excited states, derived from the  $d\sigma^* \rightarrow p\sigma$  one-electron transition, are strongly stabilized relative to their respective monomer  $d_{z^2} \rightarrow p_z$  states.

While this picture accounts nicely for several important aspects of the metal-metal interaction, it has also been noted that several monomer  $d \rightarrow p$  excitations are only slightly perturbed in the  $d^8$ - $d^8$  complexes.<sup>1</sup> This observation and the rather long metal-metal bonds (3.1 - 3.3 Å)<sup>11,12</sup> suggest that  $d^8$ - $d^8$  species should be examined from a valence bond (VB) "weak-coupling" viewpoint.

In addition to beginning such an examination, and to complement earlier investigations of rhodium complexes, ( $\text{Rh}_2(\text{TMB})_4^{2+}$ ;  $\text{Rh}_2\text{b}_4^{2+}$ , b = 1,3-diisocyanopropane), we undertook a detailed study of the electronic spectrum of  $\text{Ir}_2(\text{TMB})_4^{2+}$  (TMB = 2,5-diisocyano-2,5-dimethylhexane). A change in the metal from rhodium to iridium should perturb the electronic structure, thereby altering the characteristic  $d^8$ - $d^8$  electronic spectrum and allowing us to test key aspects of the MO and VB models. Electronic absorption and magnetic circular dichroism (MCD) spectra of  $\text{Rh}_2(\text{TMB})_4^{2+}$  and  $\text{Ir}_2(\text{TMB})_4^{2+}$  are reported along with polarized single-crystal absorption spectra of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2 \cdot \text{CH}_3\text{C}_6\text{H}_5$ .

## Experimental

### Synthesis

All synthetic procedures were carried out with standard Schlenk techniques unless otherwise specified. Tetrahydrofuran (THF) was distilled from  $\text{CaH}_2$  prior to use. All other solvents were from freshly opened bottles with no further purification, except as noted. All solvents were Schlenk-degassed prior to use. Standard procedures were used to prepare 2,5-diisocyano-2,5-dimethylhexane (TMB),<sup>11</sup>  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$ ,<sup>13</sup>  $[\text{Ir}(\text{COD})\text{Cl}]_2$  (COD = 1,5-cyclooctadiene),<sup>14</sup>  $[\text{Ir}(\text{COD})_2]\text{BF}_4$ ,<sup>15</sup> and  $[\text{Rh}_2(\text{TMB})_4](\text{PF}_6)_2$ .<sup>11</sup> The literature preparation of the compound  $\text{Ir}(\text{COD})(\text{acac})$ ,<sup>16</sup> a precursor to  $[\text{Ir}(\text{COD})_2]\text{BF}_4$ , was modified as noted below. All other chemicals were of reagent grade or comparable quality and were used as received. The  $^1\text{H}$  NMR spectra were obtained on a 400-MHz JNM-GX400 FT NMR spectrometer. The IR spectra were taken on a Beckman IR 4240.

#### $[\text{Ir}_2(\text{TMB})_4](\text{PF}_6)_2$

A filtered methanol solution of  $\text{NaPF}_6$  was added to a stirred methanol solution of  $[\text{Ir}_2(\text{TMB})_4]\text{Cl}_2$  (prepared by the method of Reference 13). Immediate precipitation of a blue powder occurred. The solution was filtered and the blue material recrystallized by slow evaporation of a toluene/acetonitrile solution. Electronic absorption spectra of this compound were identical with that of the  $\text{B}(\text{C}_6\text{H}_5)_4^-$  salt except for the absence of absorption that was due to the  $\text{B}(\text{C}_6\text{H}_5)_4^-$  anion in the UV region.

#### $\text{Ir}(\text{COD})(\text{acac})$

A THF solution of  $\text{Ti}(\text{acac})_3$  (acac = acetylacetone) was added with stirring to a THF solution of  $[\text{Ir}(\text{COD})\text{Cl}]_2$ . Formation of a white precipitate,  $\text{TiCl}_3$ , occurred immediately. The solution was stirred for approximately 1 h to ensure complete reaction. THF was removed under vacuum, and pentane was added. The solution was

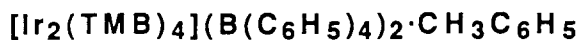
filtered through a Celite plug to yield a clear, bright-yellow filtrate. The pentane was removed under vacuum to yield a bright-yellow crystalline material. The material was further purified by sublimation at 70-80 °C onto a water-cooled cold finger condenser (0 °C) at  $10^{-5}$  to  $10^{-6}$  torr. The final material was a bright-yellow powder. Calculated for Ir(COD)(acac): C, 39.09; H, 4.79. Found: C, 38.3; H, 4.6. NMR:  $\delta$ (CDCl<sub>3</sub>, 20 °C); 1.60 (quartet, CH<sub>2</sub>, 4H), 1.98 (singlet, CH<sub>3</sub>, 6H), 2.23 (multiplet, CH<sub>2</sub>, 4H), 3.93 (multiplet, CH=CH, 4H), 5.49 (singlet, CH, 1H). This preparative method appeared to yield a product that was less contaminated with chloride impurities than the literature preparation.

#### *[Ir(CN-*t*-butyl)<sub>4</sub>]BF<sub>4</sub>*

To a clear yellow CH<sub>3</sub>CN solution of [Ir(COD)<sub>2</sub>]BF<sub>4</sub> was added, via syringe, a large excess of *t*-butyl isocyanide (CN-*t*-butyl). A rapid color change from yellow to orange was observed. The solution was stirred for approximately 1 h. The solvent and volatiles were then removed under high vacuum, yielding an extremely air-sensitive bright-orange solid.

Acetonitrile in this preparation was distilled twice from CaH<sub>2</sub>, degassed with a minimum of 5 freeze-pump-thaw cycles, and stored under vacuum over activated alumina. The *t*-butyl isocyanide was thoroughly degassed by 5 freeze-pump-thaw cycles. The reaction was done in a Vacuum Atmospheres inert atmosphere box to ensure oxygen exclusion. Earlier attempts to prepare the desired compound using standard Schlenk techniques yielded materials that appeared to be partially oxidized according to UV-Vis spectroscopic criteria. Attempts to prepare Ir(CN-*t*-butyl)<sub>4</sub><sup>+</sup> directly from [Ir(COD)Cl]<sub>2</sub> yielded material that did not give reproducible spectra. We infer some partial oxidation of the material enhanced by the presence of chloride ion. NMR:  $\delta$ (CD<sub>3</sub>CN, 20 °C); 1.49 (singlet, CH<sub>3</sub>). IR: (Nujol mull, NaCl plates) 2190 cm<sup>-1</sup> ( $\nu$ (NC), vs).

## X-ray Data Collection and Reduction for



Crystal data are given in Table 3.1. The thin-flat plate crystal was of a form suited for spectroscopic measurements, but ill-suited for x-ray diffraction. It was glued on a glass fiber with the *c* axis approximately parallel to the fiber. The *b* axis is perpendicular to the plate. The crystal was much larger than the collimator; it was centered on a Nonius CAD-4 diffractometer equipped with graphite-monochromated MoK $\alpha$  radiation. Unit cell dimensions were obtained from the setting angles of 14 reflections with  $8.4^\circ < 2\theta < 10.1^\circ$ ; the long *b* axis caused difficulties in indexing the reflections. A data set of 1654 reflections in  $\pm h, k, l$  out to  $9^\circ$  in  $2\theta$  was collected in a  $\theta$ - $2\theta$  scan at  $8^\circ$  per min, and then a second set of 148 data with *h* and *k* = 0 or 1 was collected with an  $\omega$  scan. The  $\theta$ - $2\theta$  data near the *b* axis were deleted and the two data sets merged. The data were corrected for absorption, and Lorentz and polarization factors were applied. The four independent iridium atoms were located from the Patterson map and their positions refined with 6 cycles of least squares (Table 3.2). A subsequent Fourier map did not have sufficient resolution to locate the atoms of the TMB ligands, so refinement was terminated.

Calculations were done with programs of the CRYM Crystallographic Computing System and ORTEP. Scattering factors and corrections for anomalous scattering were taken from a standard reference.<sup>17</sup> The function minimized in least squares was  $\sum w(F_o^2 - F_c^2)^2$ , where  $w = 1/\sigma^2(F_o^2)$ . Variances of the individual reflections were assigned based on counting statistics plus an additional term,  $(0.14)I^2$ . Variances of the merged reflections were determined by standard propagation of error plus an additional term,  $0.14 < I^2$ . The absorption correction was done by Gaussian integration over an 8x8x8 grid. Transmission factors varied from 0.151 to 0.945.

**Table 3.1.** Crystal and Intensity Collection Data.

Formula: $\text{Ir}_2\text{N}_8\text{C}_{94}\text{B}_2\text{H}_{112}$	Formula Weight: 1772.03
Crystal Color: Green	Habit: flat plate
$a = 17.678(8) \text{ \AA}$	$\alpha = 90^\circ$
$b = 55.03(4) \text{ \AA}$	$\beta = 110.54(8)^\circ$
$c = 18.01(3) \text{ \AA}$	$\gamma = 90^\circ$
$v = 16408(8) \text{ \AA}^3$	$z = 8$
$d_{\text{calc}} = 1.435(1) \text{ g cm}^{-3}$	
$\lambda = 0.71073 \text{ \AA}$	T: room
Graphite monochromator	
Space group: $P2_1/c$	Absences: $h0l, l = 2n+1; 0k0, k = 2n+1$
Crystal size: $1.37 \times 0.60 \times 0.016 \text{ mm}$	$\mu = 34.88 \text{ cm}^{-1}$
CAD-4 Diffractometer	$\theta$ - $2\theta$ scan, $\omega$ scan
$2\theta$ range: $0$ - $18^\circ$	Octants collected: $\pm h, k, l$
Number of reflections measured: 1802	
Number of independent reflections: 1505	
Number with $F_0^2 > 0$ : 1381	
Number with $F_0^2 > 3\sigma(F_0^2)$ : 1052	
Goodness-of-fit for merging data: 1.51	



**Table 3.2.** Coordinates of the Iridium Atoms.

Atom	x	y	z
Ir1	0.4155(16)	0.1059(5)	0.0742(16)
Ir2	0.4828(15)	0.1168(5)	0.2580(16)
Ir3	0.9180(15)	0.1293(5)	0.5440(15)
Ir4	0.9659(15)	0.1447(5)	0.7277(16)

## Spectroscopic Measurements

All solvents were dried and degassed by standard methods.<sup>18,19</sup> Absorption spectra presented in Figures 3.2 and 3.6-3.9 were recorded with a Cary 17 spectrophotometer. Absorption and MCD spectra presented in Figures 3.4 and 3.5 were determined simultaneously and synchronously along the same light path by means of a computer-controlled spectrometer described previously.<sup>20</sup> A field strength of 7.0 T was provided by a superconducting magnet system (Oxford instruments SM2-7, fitted with a bore tube held at room temperature). The  $\Delta\epsilon_M$  values are presented in units of  $M^{-1} \text{ cm}^{-1} \text{ T}^{-1}$ , where T is the magnetic field strength in Tesla units. All spectra were corrected for solvent blank. Single-crystal polarized absorption measurements employed dual Glan-Thompson, air-spaced calcite polarizers. Low-temperature absorption experiments were performed with two temperature-control systems: a quartz optical dewar of local design (77 K measurements); or a CTI-Cryogenics model 70 Cryodyne Cryocooler.

Propionitrile/2-methyl-THF (1:2 volume ratio) solutions for 77 K glassy matrix experiments were prepared on a high-vacuum line in an evacuable 1 cm pathlength quartz cuvette attached to a high-vacuum Teflon stopcock.

Absorption spectra at ~20 K were recorded of single crystals mounted on quartz flats. Crystal orientation was established with a polarizing microscope; the crystals transmitted light yellow-green and dark-blue light for polarization parallel to their two extinction directions determined between crossed polarizers. The former extinction was found to correspond to  $\parallel a$  of the (010) crystal face. The extinctions were very sharp with white light, and are therefore inferred to be wavelength-independent in the visible region. The crystals were carefully masked with heat-conducting grease (fine copper powder suspended in vacuum grease) in order to provide good thermal contact with the cold station of the cryostat. The thickness of the  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  crystal whose spectra are given in Figures 3.7 and 3.8 was determined with a micrometer to be

21  $\mu\text{m}$ . The thickness of the crystal whose spectra are shown in Figure 3.9 was not determined.

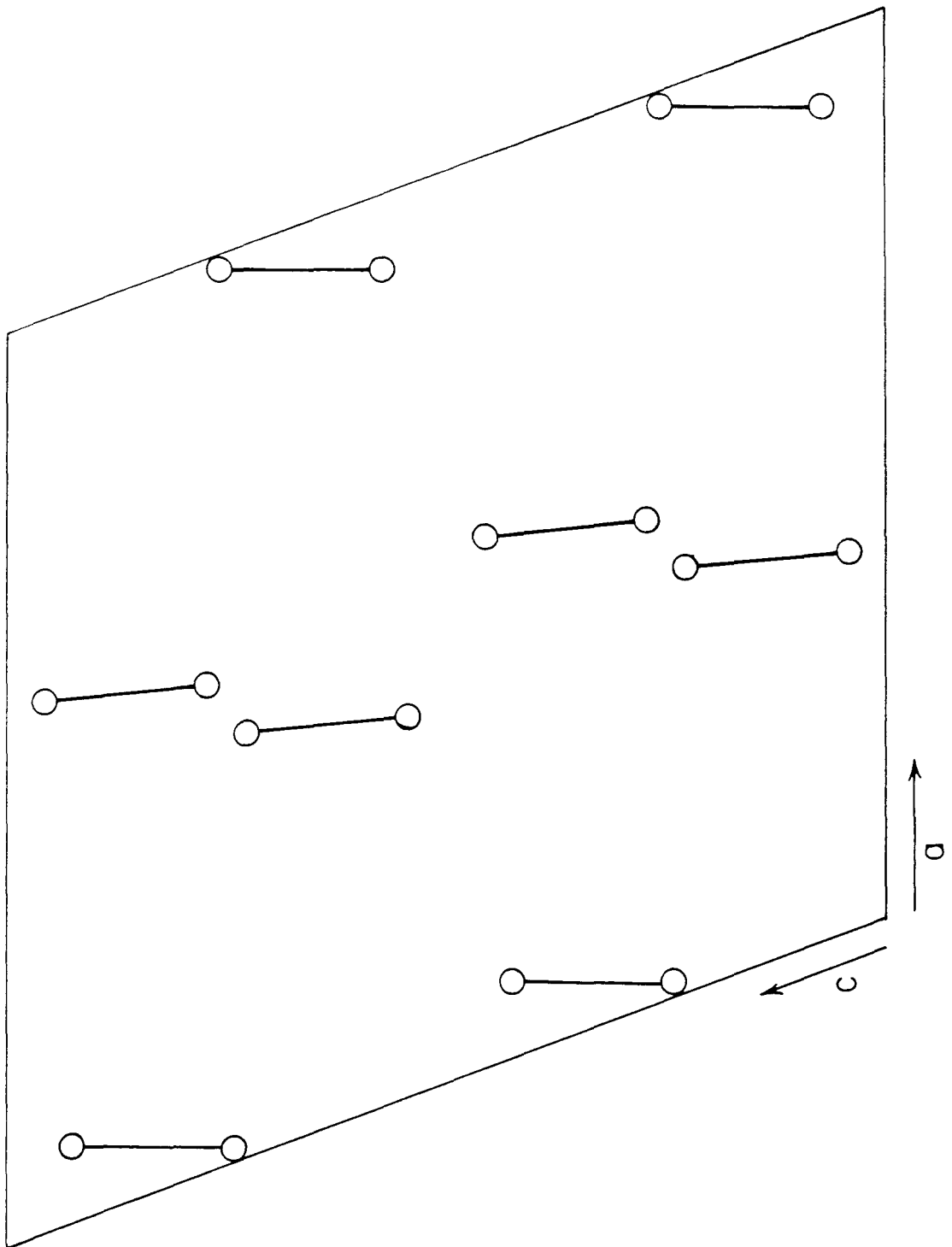
## Results and Discussion

### Crystal Structure of $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_2)_2 \cdot \text{CH}_3\text{C}_6\text{H}_5$

In the course of our studies, we required thin single crystals for spectroscopic purposes. The crystal structure of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2 \cdot 2(\text{CH}_3\text{CN})$ , which had been determined previously,<sup>12</sup> revealed a cation orientation amenable to polarized spectroscopy, but crystals obtained from  $\text{CH}_3\text{CN}$  solution were chunky blue blocks, far too thick for spectroscopy. Attempts to grow thin crystals eventually yielded clusters of thin blue-green plates<sup>21</sup> by slow evaporation of toluene-acetonitrile solutions, from which high-quality spectroscopic data were obtained. X-ray examination of these crystals showed them not to be isomorphous with those of the structurally characterized material. NMR spectra showed that these crystals contained toluene ( $n \approx 1$ ) instead of the acetonitrile of crystallization that was found in the previously characterized phase.<sup>12,22</sup>

The structure of this phase was partially determined, as described in the experimental section. The well-developed crystal face was (010); in Figure 3.1 we show a (010) view of the Ir-Ir units of the structure. The Ir<sub>2</sub> bond distances of the two independent Ir<sub>2</sub> units in this structure (3.16(4) and 3.23(4) Å) are, within their large esd's, equivalent to those found in the earlier study, 3.199(1) Å.<sup>12</sup> The Ir<sub>2</sub> units are all aligned nearly perpendicular to the *a* axis, which corresponds to an optical extinction, and lie very nearly in the (010) plane. Thus, our single-crystal polarized spectra should approximate molecular *z* and *x,y* spectra for the  $\perp a$  and  $\parallel a$  polarizations, respectively. It should be noted that for a monoclinic crystal face that does not contain the *b* axis, extinctions are not required to be parallel to a crystallographic axis, nor are they required to be wavelength-independent. However, both properties were found (within experimental error).

**Figure 3.1.** Projection of Ir-Ir units onto the (010) plane. Ir position indicated by o. The two metal atoms of a dimer unit are connected by a solid line.



### Electronic Spectrum of $\text{Ir}(\text{CN-}t\text{-butyl})_4^+$

For comparison to binuclear  $\text{Ir}(\text{I})$  complexes, it is important to have a reliable monomer electronic absorption spectrum. This apparently simple task proved to be difficult. We eventually obtained reproducible spectra for  $\text{Ir}(\text{CN-}t\text{-butyl})_4^+$ , when a synthetic route (see the experimental section) was employed that rigorously excluded both halide and oxygen. The isolated  $\text{BF}_4^-$  salt is remarkably air-sensitive, instantaneously changing color from orange to blue-green upon exposure to trace contamination of air.

The probable explanation of this color change is the formation of partially oxidized oligomers, analogous to well-characterized rhodium species such as  $\text{Rh}_4\text{b}_8^{6+}$ , which have similar colors.<sup>23,26</sup>

Well-characterized  $\text{Rh}(\text{I})$  isocyanide oligomeric systems, which are relatively air-insensitive, hence easier to handle, are thermally labile (because of very weak ground-state  $\text{Rh-Rh}$  bonding) but not photosensitive, since the excited states are strongly bound.<sup>1,30,31</sup> The similar spectroscopic properties of binuclear  $\text{Rh}(\text{I})$  and  $\text{Ir}(\text{I})$  isocyanide complexes that are documented here and elsewhere suggest that their thermal and photochemical properties would not be very different. In contrast, partially oxidized  $\text{Rh}(\text{I})$  isocyanide oligomers are thermally much less labile (the  $\text{Rh-Rh}$  ground state is strongly bound) but undergo dissociative photochemistry.<sup>23</sup> Moreover, net photoreduction has been noted for  $\text{Rh}$  isocyanide systems.<sup>13</sup> With regard to authentic  $\text{Ir}(\text{I})$  oligomers, we note that samples of  $[\text{Ir}(\text{CN-}t\text{-butyl})_4]\text{BF}_4$  during dissolution in  $\text{CH}_3\text{CN}$  show flashes of blue-green color near the dissolving crystals, the color disappearing with diffusion into the bulk solvent. This may represent  $\text{Ir}(\text{I})$  oligomerization in very concentrated solutions,<sup>27</sup> as previously reported for  $\text{Rh}(\text{I})$  isocyanides, including the *t*-butyl isocyanide complex.<sup>31</sup> We have not attempted to quantify this observation, however, because of the extreme difficulties involved in working with this compound.

The electronic absorption spectrum found for  $[\text{Ir}(\text{CN-}t\text{-butyl})_4]\text{BF}_4$  in dilute solution is shown in Figure 3.2; data are summarized in Table 3.3. The observed bands are assigned to  $d_z^2 \rightarrow p_z$ ,  $d_{xz,yz} \rightarrow p_z$ , and  $d_{xy} \rightarrow p_z$  excitations, as indicated by the extensive work of Mason and coworkers.<sup>28,32,33</sup> The features at 291 and 343 nm could each be assigned either to pure electronic transitions (Table 3.3) or to vibronic sidebands involving one quantum of  $\nu(\text{NC})$  built upon lower energy transitions. For the 291 nm feature, the vibronic assignment is excluded because the intensity is much too high relative to that seen for authentic  $\nu(\text{NC})$  sidebands in the present work (*vide infra*) as well as in a previous investigation.<sup>1</sup> The interpretation of the 343 nm feature is, however, ambiguous.

Our results are consistent with the conventional MO scheme for square-planar  $d^8$  complexes containing  $\pi$ -backbonding ligands.<sup>26,32-34</sup> The filled d level ordering is  $d_z^2 > d_{xz,yz} > d_{xy}$ , and no electronic transitions involving the empty  $d_{x^2-y^2}$  orbital have been located experimentally. As usual for such complexes, the last point only suggests the probable energetic ordering  $d_{x^2-y^2} > p_z$  (i.e., that the LUMO is  $p_z$ ), since d-d transitions to  $d_{x^2-y^2}$  are dipole-forbidden, and could be obscured by the intense (allowed)  $d \rightarrow p$  absorptions.

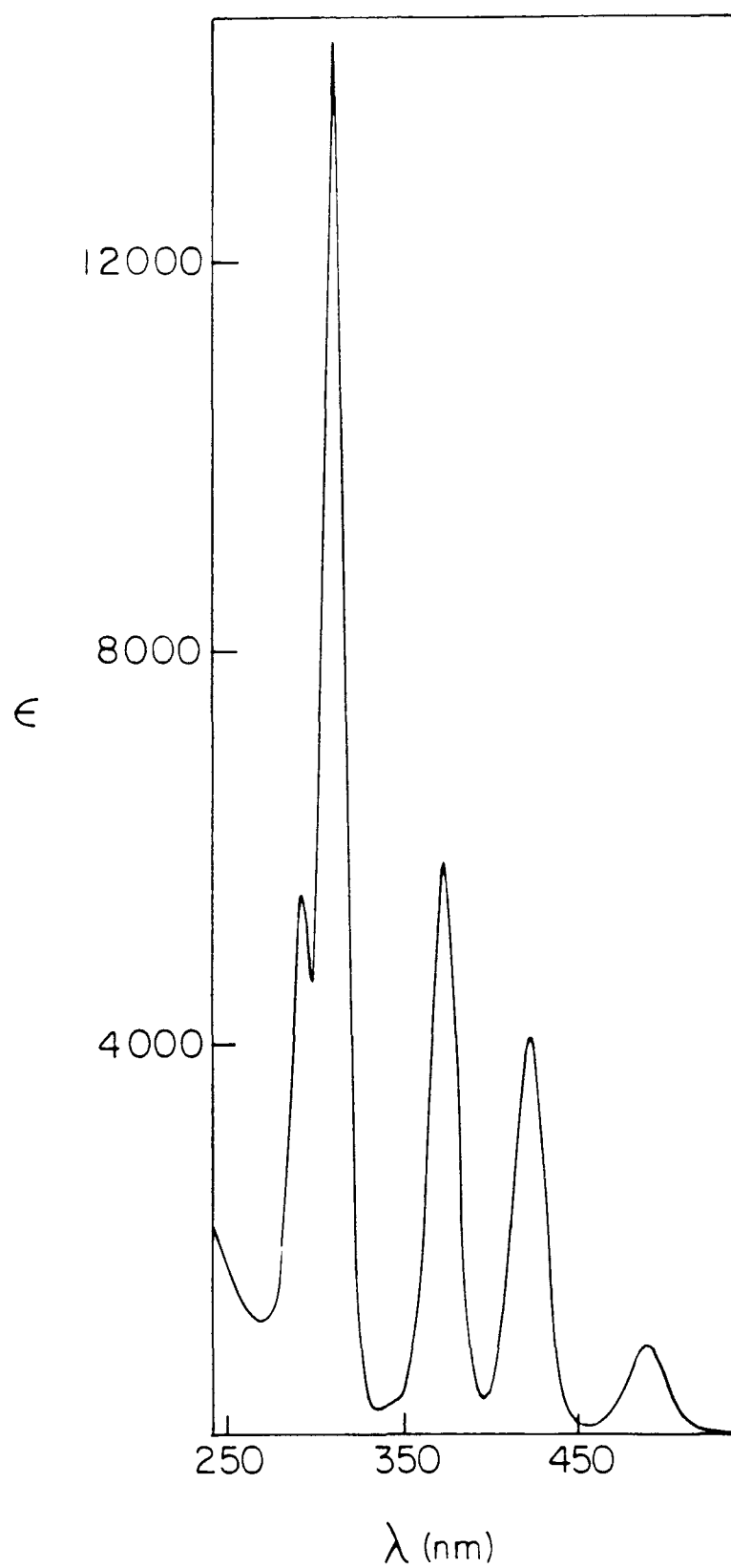
### Valence-Bond Model

Consider a dimer with no one-electron coupling of the two centers. For clarity, we consider only singlet states of a  $D_{4h}$  symmetry dimer. Low-lying molecular states involving  $d_z^2 \rightarrow p_z$  and  $d_{xz,yz} \rightarrow p_z$  excitations are shown in Figure 3.3.

The single-center excitations are Davydov-coupled<sup>35</sup> in symmetric and antisymmetric fashion to yield, for the  $^1A_{2u}$  ( $d_z^2 \rightarrow p_z$ ) monomer excited state, for example,  $^1A_{1g}$  and  $^1A_{2u}$  dimer states. The MO picture gives a total of four excited states correlating to  $d_z^2 \rightarrow p_z$  excitations. In the VB picture, the remaining two excited states are a  $^1A_{1g}$ ,  $^1A_{2u}$  pair corresponding to excitation of a  $d_z^2$  electron on one center into a  $p_z$



**Figure 3.2.** Electronic absorption spectrum of  $[\text{Ir}(\text{CN-t-butyl})_4]\text{BF}_4$  in  $\text{CH}_3\text{CN}$  at 25 °C.



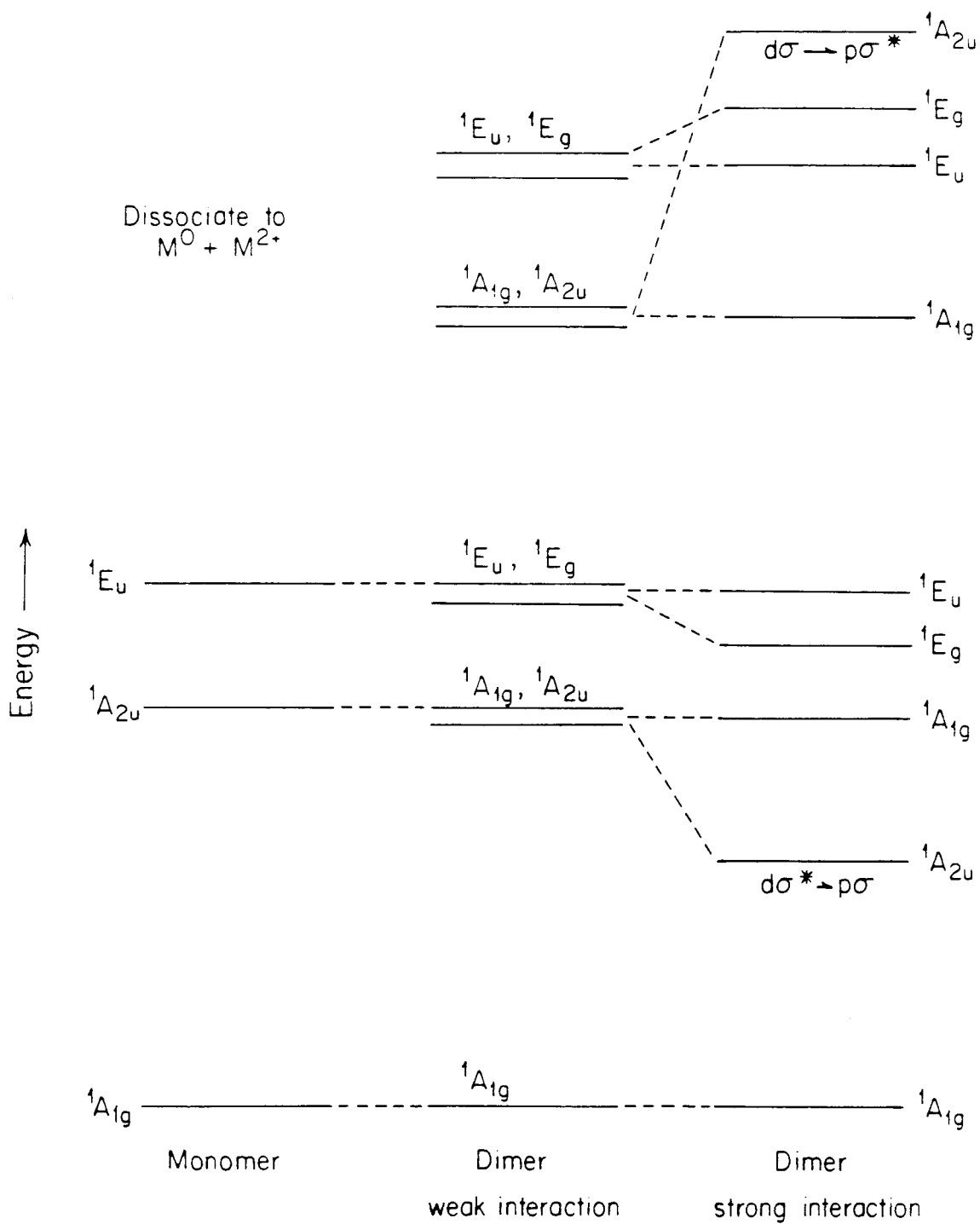
**Table 3.3.** Absorption spectral data for  $[\text{Ir}(\text{CN-t-butyl})_4]\text{BF}_4$  measured in  $\text{CH}_3\text{CN}$  at 25 °C.

$^1\text{A}_{1g} \rightarrow$	$\lambda_{\text{max}}$ , nm ( $\epsilon$ , $\text{M}^{-1}\text{cm}^{-1}$ )	fwhm, $\text{cm}^{-1}$ a
$\text{E}_g(^3\text{A}_{2g}), d_z^2 \rightarrow p_z$	489 (870)	1130
$\text{A}_{2g}(^1\text{A}_{2g}), d_z^2 \rightarrow p_z$	423 (4100)	1120
$\text{E}_g(^3\text{E}_g), d_{xz,yz} \rightarrow p_z$	373 (5900)	1300
$\text{E}_g(^1\text{E}_g), d_{xz,yz} \rightarrow p_z$	308 (14500)	1360
$\text{E}_g(^3\text{B}_{1g}), d_{xy} \rightarrow p_z$	291 (5500)	b

a. fwhm - full width at half the maximum height.

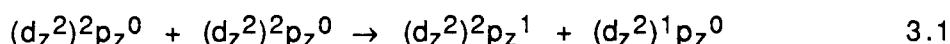
b. Not measured.

**Figure 3.3.** Qualitative VB state diagram for weak d/p interaction of two Rh(I) or Ir(I) monomers. See text. Note that for the strong interaction case, state energies are implicitly for the equilibrium bond distance of each individual state.



orbital on the other center; thus, they are "ionic" states that must dissociate to  $M^{2+}$  plus  $M^0$ . As such, they must lie at much higher energy, *even for no metal-metal bonding*.

The energy of the metal-to-metal charge transfer states are dependent on the ease of ionization of the metal center and the ability of a single metal center to accommodate an additional electron. Such a transition should scale roughly as  $IP_d - EA_p$ . An estimate of the molecular  $IP_d - EA_p$  can be obtained by considering the atomic  $IP_d - EA_p$ . The atomic IP and EA calculated for Rh and Ir are listed in Table 3.4.<sup>36</sup>



The transition energies, as defined above, are listed in Table 3.4, corrected for the  $1/R$  repulsion terms. The transition energies, as they stand, are not comparable to the molecular system, in that the molecular stabilization of the atomic charge transfer state (via delocalization of charge onto the ligands, in a simplistic viewpoint) has not been accounted for. An estimate of the molecular stabilization accompanying the metal-to-metal charge transfer can be obtained from earlier work on  $Re_2Cl_8^{2-}$ , a quadruply bonded ( $\sigma^2\pi^4\delta^2$ ) metal dimer.<sup>37</sup> Similar to the binuclear  $d^8$  complexes, the  $\delta$  interaction is best viewed in the VB "weak-coupling" limit. Thus, the  $^1(\delta \rightarrow \delta^*)$  transition is best described as a metal-to-metal charge-transfer (MMCT) transition. It can be shown that the energy separation of the  $^3(\delta\delta^*)$  state, the lowest energy covalent excited state, and the  $^1(\delta\delta^*)$  state, the lowest energy charge-transfer excited state, is  $2K$ :<sup>38</sup>

$$K = 1/2[(aa|aa) - (aa|bb)], \quad 3.3$$

**Table 3.4.** Metal-to-metal charge transfer transition energies.

	Calculated Atomic State Energies (hartrees)		
	$s^0d^8$	$p^1d^8$	$s^0d^7$
Ir	-103.51082	-103.54144	-102.83667
Rh	-108.36242	-108.47397	-107.73884

---

	IP (eV)	-EA (eV)
Ir	15.47	-3.19
Rh	16.97	-3.03

---

$M^+ + M^+ \rightarrow M^{2+} + M^0$			
	IP-EA (eV)	(IP-EA) - 1/R (eV)	(IP-EA) - 1/R + MS (eV) <sup>a</sup>
Ir	12.8	8.3	5.7 (45000 cm <sup>-1</sup> )
Rh	13.9	9.4	6.8 (55000 cm <sup>-1</sup> )

a. M.S. - Molecular stabilization = 2.6 eV.

where  $(aa|aa)$  is the energy required to transfer a  $\delta$  electron from metal center A to metal center B at infinite A-B distance. The  $(aa|bb)$  term is the electrostatic correction to finite R that we have already applied, as noted above. A value for the molecular stabilization of an MMCT state can be obtained by comparing the energy separation of the  $^3(\delta\delta^*)$  and  $^1(\delta\delta^*)$  excited states to the estimate, using atomic IP and EA values. For  $\text{Re}_2\text{Cl}_8^{2-}$ , the molecular stabilization of an MMCT state is 2.65 eV. This estimate is obtained using the experimentally observed  $^1(\delta\delta^*)$  transition energy<sup>39</sup> (which is greater than, but very close to, the singlet-triplet energy separation) and the calculated energy separation using *ab initio* atomic energies.<sup>40</sup>

Using 2.6 eV as a rough estimate of the molecular stabilization of the charge-transfer state, an estimate of the MMCT state energy for the Ir(I) and Rh(I) complexes is obtained (Table 3.4). The transition energies are indicated to be in the UV with those of Ir lower than those of Rh. These estimates are upper limits to the transition energies. The calculated IP and EA for Ir and Rh may be in error because of correlation effects not accounted for in the calculations. However, the excitation in question does not require an increased pairing of electrons; therefore, energy differences (in particular, the predicted strong decrease in energy for  $\text{Ir}_2$  versus  $\text{Rh}_2$ ) may not have a large error associated with them. These are very crude estimates, and the crude agreement with experiment that will be evident later should be taken only to indicate that our assignments are not absurd.

The VB wave functions are equal mixtures of the two MO wavefunctions of the same symmetry. Thus, the two  $^1A_{2u}$  wave functions are given by  $d\sigma^*p\sigma \pm d\sigma p\sigma^*$ , the plus combination corresponding to the lowest energy state. If we increase the metal-metal one-electron interaction (shorten the metal-metal distance), the two  $^1A_{2u}$  states will mix in such a way as to resolve into the two MO wave functions, as indicated on the right-hand side of Figure 3.3.<sup>41</sup> The same thing should happen for the other  $d \rightarrow p$  excited states if the metal-metal distance is decreased. However, there is much less



reason for them to do so. Both the  $^1E_u$  and  $^1A_{1g}$  states are, at long distance, equal mixtures of bonding-bonding and antibonding-antibonding MO states; the  $^1A_{1g}$  states, for example, are  $(d\sigma^*p\sigma^* \pm d\sigma p\sigma)$ . Thus, bonding contributions tend to cancel out. It is difficult to predict the relative bonding-antibonding contributions, particularly since the various interactions ( $d\sigma$ - $d\sigma$ ,  $d\sigma$ - $p\sigma$ ,  $p\sigma$ - $p\sigma$ ,  $d\pi$ - $d\pi$ , etc.) undoubtedly maximize stabilization at greatly different metal-metal distances. However, it is reasonable that much less stabilization would arise than for  $^1A_{2u}$  ( $d\sigma^* \rightarrow p\sigma$ ), and our experimental data (*vide infra*) support the idea that the  $^1E_u$  and  $^1A_{1g}$  excitations are nearly vertical. There is, moreover, an interesting comparison available in diatomic mercury,  $Hg_2$ , which in its ground state ( $s\sigma^2s\sigma^{*2}$ ) is an extremely weakly bound van der Waals molecule.<sup>42-44</sup> However, corresponding to the atomic  $^1S \rightarrow ^3P$  transition at 254 nm,  $Hg_2$  shows two allowed transitions. One,  $s\sigma^* \rightarrow p\sigma$ , is to a strongly bound excited state, and is richly structured in  $v(Hg_2)$  and red-shifted from the atomic line. The second,  $(s\sigma/s\sigma^*) \rightarrow (p\pi/p\pi^*)$ , is vertical, and is shifted from the atomic line by only 57  $cm^{-1}$ . These two  $Hg_2$  states are clearly analogous to the low-lying  $^1A_{2u}$  and  $^1E_u$  excited states of  $d^8$ - $d^8$   $Rh_2$  and  $Ir_2$  species.

Finally, we consider the  $^1E_g$  excited states. The VB wave functions of these states  $(d\pi^*p\sigma \pm d\pi p\sigma^*)$  indicate that the lowest  $^1E_g$  should be stabilized by resolving into the antibonding-to-bonding  $d\pi^*p\sigma$  wave function. However, it is unclear how much stabilization can actually be expected, since the  $d\pi$  overlap will be very small except at very short metal-metal distances, which would be precluded by repulsive metal-metal and ligand-ligand interactions. We compromise in Figure 3.3 by suggesting a smaller stabilization than for  $^1A_{2u}$ .

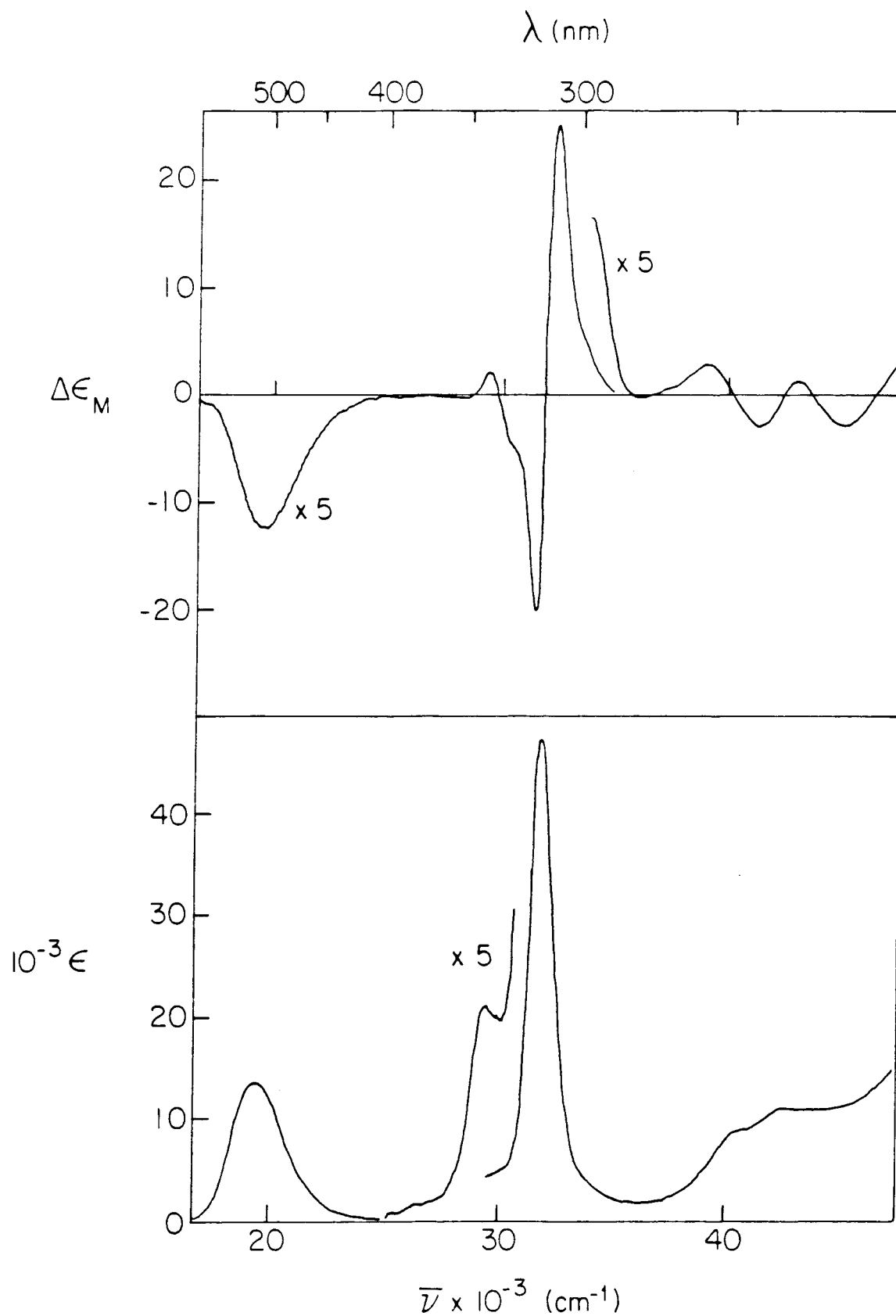
The important points that emerge from this discussion are that, in contrast to a simple MO picture, a VB model readily accommodates both weakly and strongly perturbed monomer transitions, and that it identifies fully half of the dimer electronic excited states as ionic (MMCT); these ionic states should lie at considerably higher

energy than the monomer excited states. Both of these points prove to be crucial to understanding the electronic spectra of  $d^8$ - $d^8$  complexes.

### Solution Electronic Absorption and MCD Spectra

The electronic absorption and MCD spectra of  $\text{Rh}_2(\text{TMB})_4^{2+}$  are shown in Figure 3.4. The assignments have been discussed extensively elsewhere, with supporting single-crystal data,<sup>1,11</sup> and our present (entirely consistent) data are offered largely for comparisons with the Ir complex. The theory necessary for interpretation of the MCD spectrum has been given elsewhere.<sup>3</sup> In the following discussion, we will use  $D_{4h}$  symmetry labels. The symmetry of the TMB complexes is only  $D_4$ , because of significant torsion of the  $\text{MC}_4$  units with respect to each other.<sup>12</sup> The effect of the lower symmetry is that  $D_{4h}$  g/u labels should be dropped. However, the only effect of the lower symmetry upon the states under consideration is that  $D_{4h}$  states of  $E_g$  symmetry become formally x,y-dipole-allowed, and these transitions are expected to be weak, since it depends upon low-symmetry-induced mixing of the  $D_{4h}$   $E_u$  and  $E_g$  states. The low-energy feature ( $\lambda_{\text{max}}$  517 nm) is assigned to the  $d\sigma^* \rightarrow p\sigma$  singlet ( $A_{1g}(^1A_{1g}) \rightarrow A_{2u}(^1A_{2u})$ ) transition. This band shows the expected -B term in the MCD spectrum.<sup>45</sup> The strong band centered at 314 nm is attributed to the  $A_{1g}(^1A_{1g}) \rightarrow E_u(^1E_u)$  ( $d\pi \rightarrow p\sigma$ ) excitation. This transition is to a degenerate excited state and the MCD band appears to be dominated by a +A term. A weak band ( $\lambda_{\text{max}}$  340 nm) is observed to lower energy of the  $E_u(^1E_u)$  band. From low-temperature, polarized single-crystal spectra, this band is deconvoluted into two features,<sup>1</sup> an x,y-polarized component assigned to the  $E_u$  component of the  $^3E_u$  transition ( $\lambda_{\text{max}}$  342 nm) and a z-polarized  $A_{2u}$  component ( $\lambda_{\text{max}}$  335 nm). The  $E_u$  component of the  $^3E_u$  state should have +A and +B terms associated with it and the  $A_{2u}$  state, a -B term. The two  $^3E_u$  states appear to be resolved in the MCD spectrum. A slight  $-\Delta\epsilon$  is observed on the low-

**Figure 3.4.** Electronic absorption (lower curves) and MCD (upper curves) spectra of  $[\text{Rh}_2(\text{TMB})_4](\text{PF}_6)_2$  in  $\text{CH}_3\text{CN}$  at 25 °C. Note that the MCD scales in the regions 210 - 295 nm and 400 - 600 nm have been expanded fivefold; similarly, the absorptivity scale has been expanded fivefold in the 325 - 400 nm region.  $\Delta\epsilon$  has units of  $(\text{M cm T})^{-1}$ .



energy side of the MCD band at 330 nm ( $\sim 3.00 \mu\text{m}^{-1}$ ), indicating a weak +A term; this transition is dominated by a +B term. Similarly, weak +A terms have been observed for analogous states of monomeric Rh(I) isocyanides.<sup>28</sup> To higher energy, a transition with a -B term is observed, consistent with the presence of an  $A_{2u}$  state. Several weak features are observed to higher energy of the  $E_u(1E_u)$  transition. The shoulder at 290 nm appears to have a +B term, consistent with an  $E_u(3B_{1u}) d_{xy} \rightarrow p\sigma$  assignment. However, this would also be consistent with the feature being a vibronic  $E_u(1E_u) + \nu(\text{NC})$  band, and the energy and low intensity of the band are reasonable for either assignment. The two bands above 250 nm ( $4.00 \mu\text{m}^{-1}$ ) appear to possess A terms, indicating degenerate excited states, although the interpretation of the data for the higher-energy band is ambiguous. Emission polarization ratios of these bands measured for 77 K glassy solutions are also consistent with x,y-polarization, again most clearly for the lower-energy band.<sup>46</sup> These bands have not been reported in previous work (largely because they are obscured by anion absorption for  $\text{B}(\text{C}_6\text{H}_5)_4^-$  salts), but we note here that similar features are seen in the spectra of all other binuclear Rh(I) isocyanide complexes that we have examined, whereas monomers such as  $\text{Rh}(\text{CN-}t\text{-butyl})_4^+$  show no resolved absorption in this region.<sup>46</sup> Assignments listed in Table 3.5 will be justified later.

The electronic absorption and MCD spectra of  $[\text{Ir}_2(\text{TMB})_4](\text{PF}_6)_2$  are shown in Figure 3.5. In analogy to  $\text{Rh}_2(\text{TMB})_4^{2+}$ , the intense absorption centered at 625 nm ( $\epsilon = 11,200 \text{ M}^{-1}\text{cm}^{-1}$ ) is assigned to the fully allowed  $A_{1g}(1A_{1g}) \rightarrow A_{2u}(1A_{2u}) (d\sigma^* \rightarrow p\sigma)$  excitation. This band shows the expected -B term in the MCD. The intense features below 300 nm, having maxima at 318 ( $\epsilon = 22,000 \text{ M}^{-1}\text{cm}^{-1}$ ) and 372 nm ( $\epsilon = 9750 \text{ M}^{-1}\text{cm}^{-1}$ ) are attributed to  $d\pi \rightarrow p\sigma$  transitions to  $E_u(1E_u)$  and  $E_u(3E_u)$  states, respectively. The 318 nm absorption, while complicated by the overlap of the weaker bands to low and high energy, displays the expected +A term in the MCD. The  $A_{1g}(1A_{1g}) \rightarrow E_u(3E_u)$  transition appears in the MCD as a skewed band because of the contribution of

**Table 3.5.** Absorption maxima for  $\text{Rh}_2(\text{TMB})_4^{2+}$  and  $\text{Ir}_2(\text{TMB})_4^{2+}$ . The convention we adopt in this and the following table is that the previously adopted MO  $d\sigma^* \rightarrow p\sigma$  assignments are retained for such excitations, but other transitions are labeled either as single-center excitations (e.g.,  $d_{xz,yz} \rightarrow p_z$ ) or MMCT.

$^1A_{1g} \rightarrow$	[Rh <sub>2</sub> (TMB) <sub>4</sub> ](PF <sub>6</sub> ) <sub>2</sub> in CH <sub>3</sub> CN (25 °C)		MCD	
	$\lambda_{\text{max}}$ , nm ( $\epsilon$ , M <sup>-1</sup> cm <sup>-1</sup> )	fwhm, cm <sup>-1</sup> a	$\nu$ , cm <sup>-1</sup> × 10 <sup>-3</sup>	$\Delta\epsilon_M$ , M <sup>-1</sup> cm <sup>-1</sup> T <sup>-1</sup>
	$\nu$ , cm <sup>-1</sup> × 10 <sup>-3</sup>			
$E_u(^3A_{2u}), d\sigma^* \rightarrow p\sigma$	b			
$^1A_{2u}, d\sigma^* \rightarrow p\sigma$	517 (13600)	2800	1.95	-2.45
	1.93			
$E_u(^3E_u), d_{xz,yz} \rightarrow p_z$	340 (4200, sh)		{ 2.78 2.87	{ -0.076 0.0
	2.94		2.95	+2.38
$A_{2u}(^3E_u), d_{xz,yz} \rightarrow p_z$	c		3.04 (sh)	-4.67
$E_u(^1E_u), d_{xz,yz} \rightarrow p_z$	314 (47500)	1400	{ 3.15 3.19	{ -20.1 0.0
	3.18		3.24	+25.2
$E_u(^3B_{1u}), d_{xy} \rightarrow p_z$	290 (3500, sh)		~3.3 (sh)	~+5
and/or $E_u(^1E_u) + \nu(\text{CN})$	3.45			
			{ 3.91 4.02	{ +0.059 0.0
MMCT(?)	245 (8600, sh)		4.14	-0.581
	4.08			
MMCT(?)	234 (10700, sh)		4.24	0.0
	4.27		4.30	+0.291
			4.51	-0.531

[Ir <sub>2</sub> (TMB) <sub>4</sub> ](B(C <sub>6</sub> H <sub>5</sub> ) <sub>4</sub> ) <sub>2</sub> in CH <sub>3</sub> CN (25 °C)					
<sup>1</sup> A <sub>1g</sub> →	λ <sub>max</sub> , nm (ε, M <sup>-1</sup> cm <sup>-1</sup> )		fwhm, cm <sup>-1</sup> <sup>a</sup>	MCD <sup>d</sup>	
	ν, cm <sup>-1</sup> × 10 <sup>-3</sup>			ν, cm <sup>-1</sup> × 10 <sup>-3</sup>	Δε <sub>M</sub> , M <sup>-1</sup> cm <sup>-1</sup> T <sup>-1</sup>
E <sub>u</sub> ( <sup>3</sup> A <sub>2u</sub> ), dσ* → pσ					
<sup>1</sup> A <sub>2u</sub> , dσ* → pσ	625	(11200)	2200	1.61	-1.26
	1.60				
E <sub>u</sub> ( <sup>3</sup> E <sub>u</sub> ), d <sub>xz,yz</sub> → p <sub>z</sub>	372	(9750)	1000	{ 2.63 2.66 2.71	{ -0.118 0.0 +6.25
	2.69				
+ ν(CN)	b				
A <sub>2u</sub> ( <sup>3</sup> E <sub>u</sub> ), d <sub>xz,yz</sub> → p <sub>z</sub>	333	(4000, sh)		3.03	-12.0
	3.00				
E <sub>u</sub> ( <sup>1</sup> E <sub>u</sub> ), d <sub>xz,yz</sub> → p <sub>z</sub>	318	(22,000)	1300	{ 3.09 (sh) 3.13 3.23	{ -10.1 0.0 +8.98
	3.14				
E <sub>u</sub> ( <sup>3</sup> B <sub>1u</sub> ), d <sub>xy</sub> → p <sub>z</sub>	305	(9500, sh)			
	3.28				
MMCT				{ 3.46 3.51 3.56	{ -1.05 0.0 +0.23
E <sub>u</sub> ( <sup>3</sup> E <sub>u</sub> ), d <sub>xz,yz</sub> → p <sub>z</sub>	-280	(8500, sh)			
	~3.6				
E <sub>u</sub> ( <sup>1</sup> E <sub>u</sub> ), d <sub>xz,yz</sub> → p <sub>z</sub>	268	(17000)	1500	{ 3.65 3.69 3.75	{ -1.34 0.0 +3.49
	3.73				
				3.88	-1.94
"dσ → pσ"	-210	(25000)		4.31	-1.20
	~4.7				

$[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in a glassy matrix of  $\text{CH}_3\text{CH}_2\text{CN}/2\text{-MeTHF}$  (77 K)

$^1\text{A}_{1g} \rightarrow$	$\lambda_{\text{max}}, \text{nm} (\epsilon, \text{M}^{-1}\text{cm}^{-1})$	fwhm, $\text{cm}^{-1}$ <sup>a</sup>
$\text{E}_u(^3\text{A}_{2u}), d\sigma^* \rightarrow p\sigma$	520 (150)	1500
$^1\text{A}_{2u}, d\sigma^* \rightarrow p\sigma$	635 (25500)	1200
$\text{E}_u(^3\text{E}_u), d_{xz,yz} \rightarrow p_z$	373 (17000)	650
+ $\nu(\text{CN})$	344 (1500, sh)	
$\text{A}_{2u}(^3\text{E}_u), d_{xz,yz} \rightarrow p_z$	331 (5000)	700
$\text{E}_u(^1\text{E}_u), d_{xz,yz} \rightarrow p_z$	320 (38000)	690
$\text{E}_u(^3\text{B}_{1u}), d_{xy} \rightarrow p_z$	304 (6000)	800
MMCT		
$\text{E}_u(^3\text{E}_u), d_{xz,yz} \rightarrow p_z$	280 (11000, sh)	
$\text{E}_u(^1\text{E}_u), d_{xz,yz} \rightarrow p_z$	268 (32000)	1000
" $d\sigma \rightarrow p\sigma^*$ "	e	

a. fwhm = full width at half the maximum height.

b. Not observed.

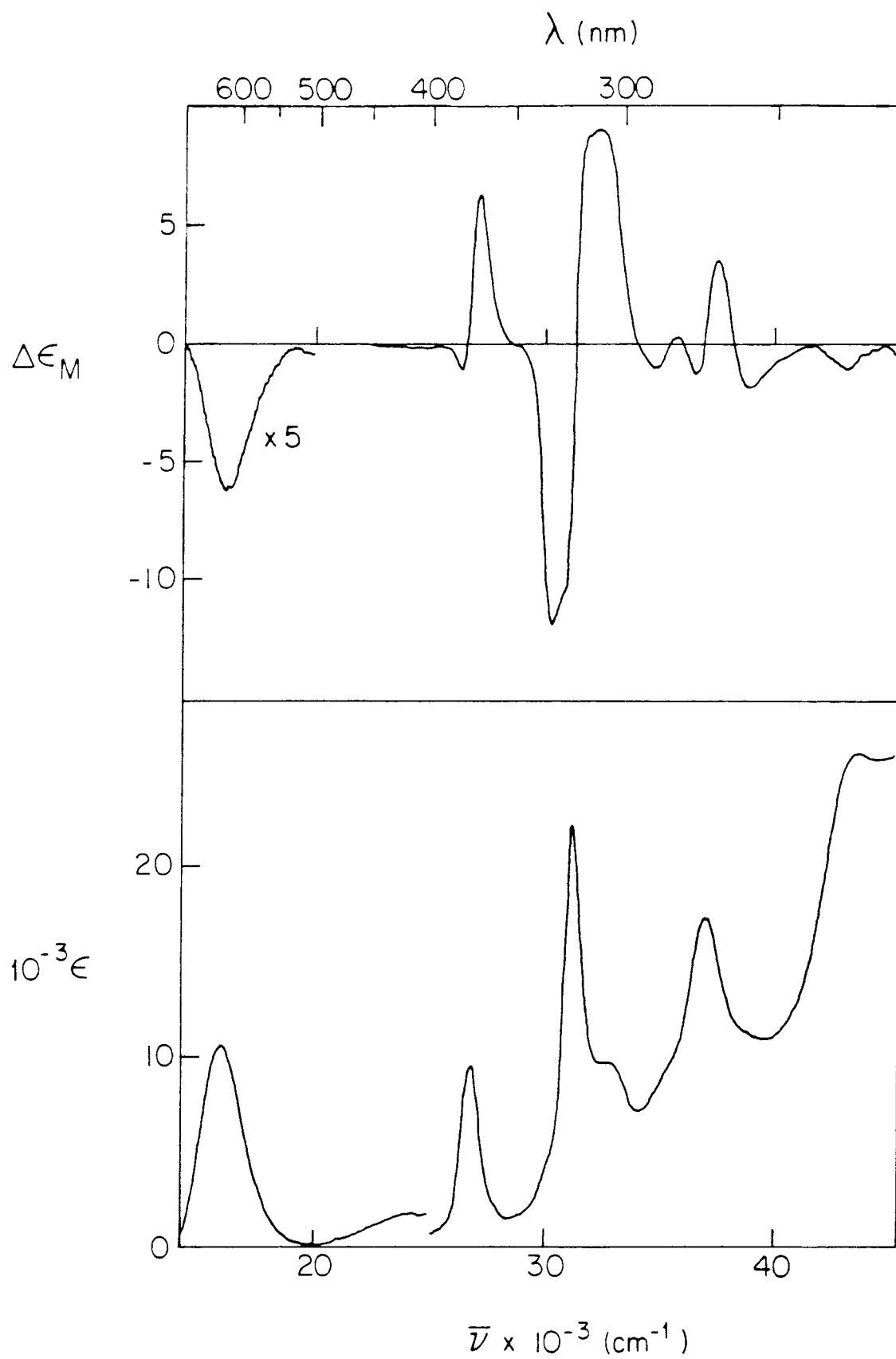
c. Polarized single-crystal spectra resolve this band into two features:  $\text{E}_u(^3\text{E}_u), \lambda_{\text{max}} 342 \text{ nm}$ ;  
 $\text{A}_{2u}(^3\text{E}_u), \lambda_{\text{max}} 335 \text{ nm}$ .

d. For  $\text{PF}_6^-$  salt.

e. Not determined because of anion absorption.



**Figure 3.5.** Electronic absorption (lower curve) and MCD (upper curves) spectra of  $[\text{Ir}_2(\text{TMB})_4](\text{PF}_6)_2$  in  $\text{CH}_3\text{CN}$  at 25 °C. Note that the MCD scale has been expanded fivefold in the 700 - 500 nm region.  $\Delta\epsilon$  has units of  $(\text{M cm T})^{-1}$ .



the +A and +B terms. In comparison to  $\text{Rh}_2(\text{TMB})_4^{2+}$ , there is a much larger +A term contributing to this transition. A feature centered at 820 nm is observed to lower energy of the 625 nm absorbance in both single crystals and concentrated solutions. The MCD of this band was not measured because of spectrometer limitations. This band is assigned to the  $^3(d\sigma^* \rightarrow p\sigma)$  transition. To lower energy of the 372 nm band, a weak shoulder is observed. A similar weak feature is observed for  $\text{Rh}_2(\text{TMB})_4^{2+}$ . Several features not well resolved for  $\text{Rh}_2(\text{TMB})_4^{2+}$  are observed for  $\text{Ir}_2(\text{TMB})_4^{2+}$ . To both the high- and low-energy side of the 318 nm band, shoulders are observed (305 and 333 nm). The low-energy shoulder appears as a -B term in the MCD spectrum, consistent with an assignment to the  $A_{2u}(^3E_u)$  transition. The shoulder to higher energy, although overlapped with the  $^1E_u$  band, appears to be dominated by a +B term, perhaps with a small -A component, in accord with the  $E_u(^3B_{1u}) d_{xy} \rightarrow p\sigma$  transition. In addition, bands to higher energy of 300 nm are resolved for both the  $\text{B}(\text{C}_6\text{H}_5)_4^-$  and  $\text{PF}_6^-$  salts. Neither of these features has an analogue in monomer spectra. The two lower-energy bands exhibit +A terms in the MCD, indicating transitions to degenerate excited states. The highest-energy shoulder possesses a weak -B term and no A term, suggestive of a nondegenerate excited state. Assignments are listed in Table 3.5.

The energies and widths of the bands derived from the lowest-energy  $d_z^2 \rightarrow p_z$  excitation for  $\text{Ir}_2(\text{TMB})_4^{2+}$  are in stark contrast to the  $d_z^2 \rightarrow p_z$  absorption in the spectrum of  $\text{Ir}(\text{CN-}t\text{-butyl})_4^+$ . The large red shift has been attributed to the smaller d-p gap in the dimer.<sup>10</sup> The greater widths of the  $A_{2u}(^1A_{2u})$  and  $E_u(^3A_{2u})$  bands are a consequence of the enhanced metal-metal bonding interaction in the  $^1,^3(d\sigma^*p\sigma)$  excited states. Promotion of an electron from  $d\sigma^*$  to  $p\sigma$  produces a large molecular distortion; this distortion is mainly a contraction along the metal-metal coordinate.<sup>1</sup>

In contrast to the  $d\sigma^* \rightarrow p\sigma$  transition, features assigned to  $d\pi \rightarrow p\sigma$  excitation have energies and widths very similar to those observed for the monomer, implying very little if any molecular distortion along the metal-metal coordinate. If one adheres

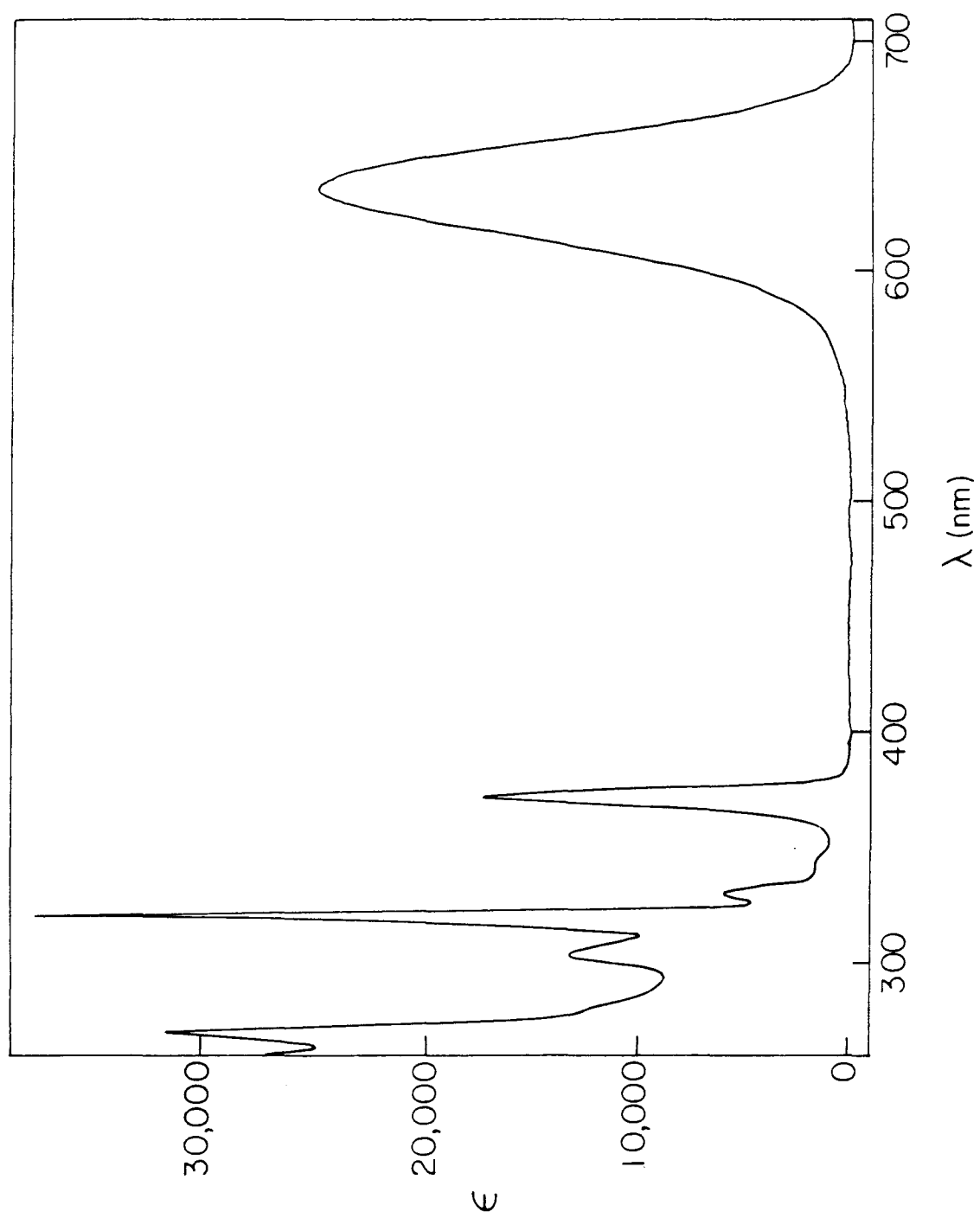
strictly to the MO model for the binuclear complex, a shift in the  $d\sigma^* \rightarrow p\sigma$  transitions implies a shift in the  $d\pi \rightarrow p\sigma$  excitations, unless  $p\sigma/d\sigma$  and  $p\sigma/p\sigma$  bonding interactions are negligible,<sup>1</sup> which would contradict the ground-state bonding picture. Since a large shift is observed for the  $d\sigma^* \rightarrow p\sigma$  transitions and very little if any shift is seen for  $d\pi \rightarrow p\sigma$ , the VB model is preferred; according to VB, the low-energy  $d\sigma^* \rightarrow p\sigma$  transition is red-shifted from the monomer  $d_z^2 \rightarrow p_z$  because of a favorable metal-metal interaction at shorter distance, with the  $^1,^3E$  states only slightly shifted from their reference monomer energies.

The weak features observed to lower energy of the  $A_{1g}(^1A_{1g}) \rightarrow E_u(^3E_u)$  transition have been the subject of much speculation. Arguments against assignment to transitions to  $d_x^2-y^2$ -derived levels (e.g., "ligand-field" excited states) have been presented, and other assignments have been suggested:  $d_{xy} \rightarrow p_z$ ,  $d\sigma^* \rightarrow p\sigma^*$ , and  $d\sigma \rightarrow p\sigma$ .<sup>1</sup>

Figure 3.6 shows the low-temperature (77 K) absorption spectrum of  $[Ir_2(TMB)_4](B(C_6H_5)_4)_2$ . Upon cooling, the expected sharpening and red shift of  $A_{1g}(^1A_{1g}) \rightarrow A_{2u}(^1A_{2u})$  ( $d\sigma^* \rightarrow p\sigma$ ) are observed. The red shift of the maximum indicates that the effective vibrational frequency that is coupled to the electronic transition is higher in the excited state than in the ground state.<sup>47</sup>

The high-energy portion of the spectrum resolves impressively upon cooling. The bands assigned to  $d\pi \rightarrow p\sigma$  excitations sharpen but do not shift. The shoulders on each side of the 320 nm band are resolved into distinct features with widths similar to those of the intense  $d\pi \rightarrow p\sigma$  absorptions. A weak shoulder at 344 nm attributable to one quantum of  $\nu(NC)$  ( $2260\text{ cm}^{-1}$ ) is resolved. The similarity of the high-energy portion of the low-temperature absorption spectrum of  $Ir_2(TMB)_4^{2+}$  to that of  $Ir(CN-t\text{-butyl})_4^+$  is remarkable.

**Figure 3.6.** Electronic absorption spectrum of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in a glassy matrix of  $\text{CH}_3\text{CH}_2\text{CN}/2\text{-MeTHF}$  (1:2 volume ratio) at 77 K.



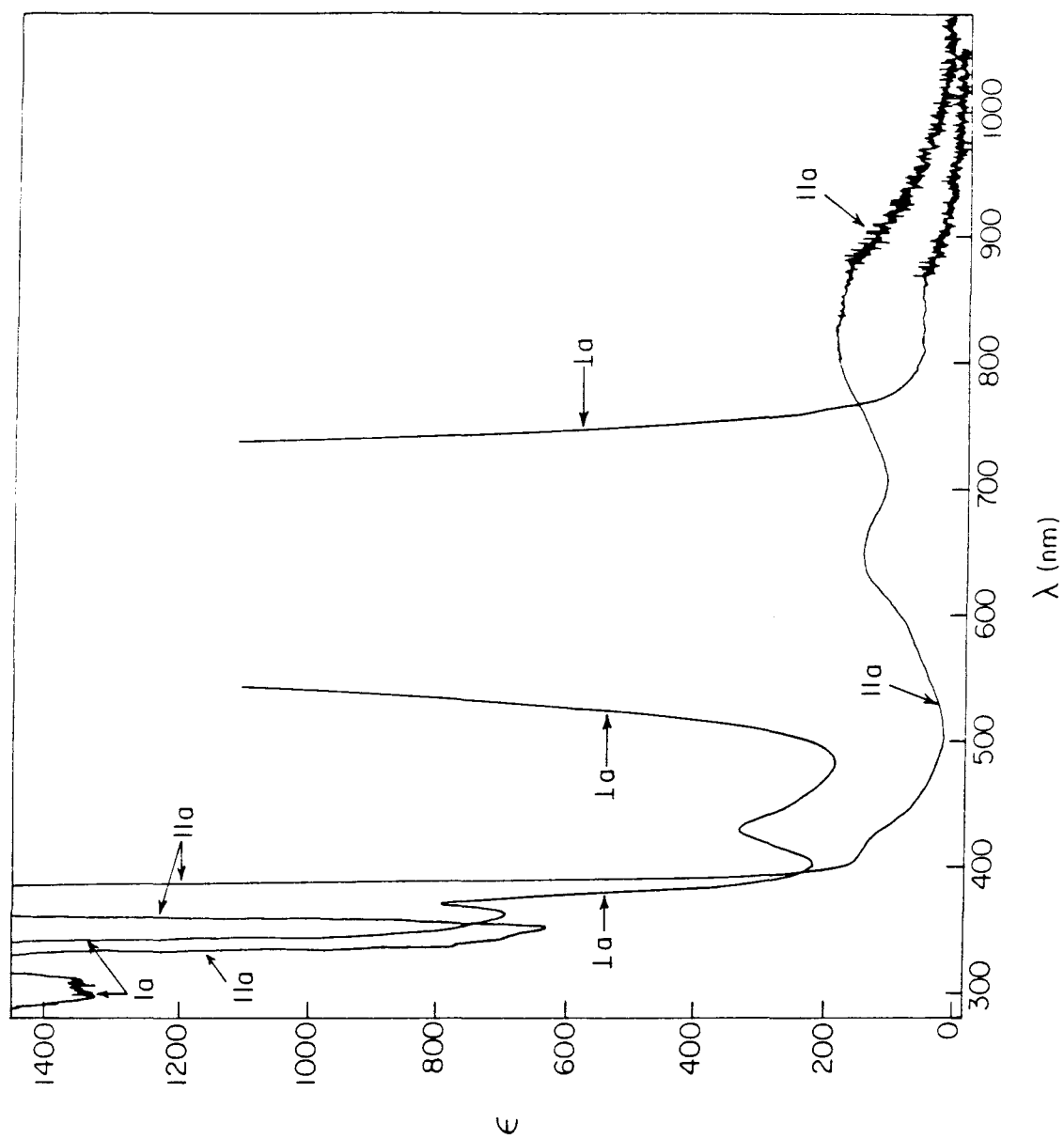
### Single-Crystal Polarized Spectra of $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2 \cdot \text{CH}_3\text{C}_6\text{H}_5$

To verify the assignments of the solution and 77 K glass spectra, and to probe the nature of the weak transitions, single-crystal polarized spectra were measured. The room-temperature polarized single-crystal absorption spectra of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2 \cdot \text{CH}_3\text{C}_6\text{H}_5$  are shown in Figure 3.7. Consistent with the earlier assignment,  $A_{1g}(^1A_{1g}) \rightarrow A_{2u}(^1A_{2u})$  ( $\lambda_{\text{max}}$  645 nm) shows the expected z-polarization. At lower energy is the x,y-polarized  $A_{1g}(^1A_{1g}) \rightarrow E_u(^3A_{2u})$  band ( $\lambda_{\text{max}}$  830 nm). These two features are the  $d\sigma^* \rightarrow p\sigma$  transitions.

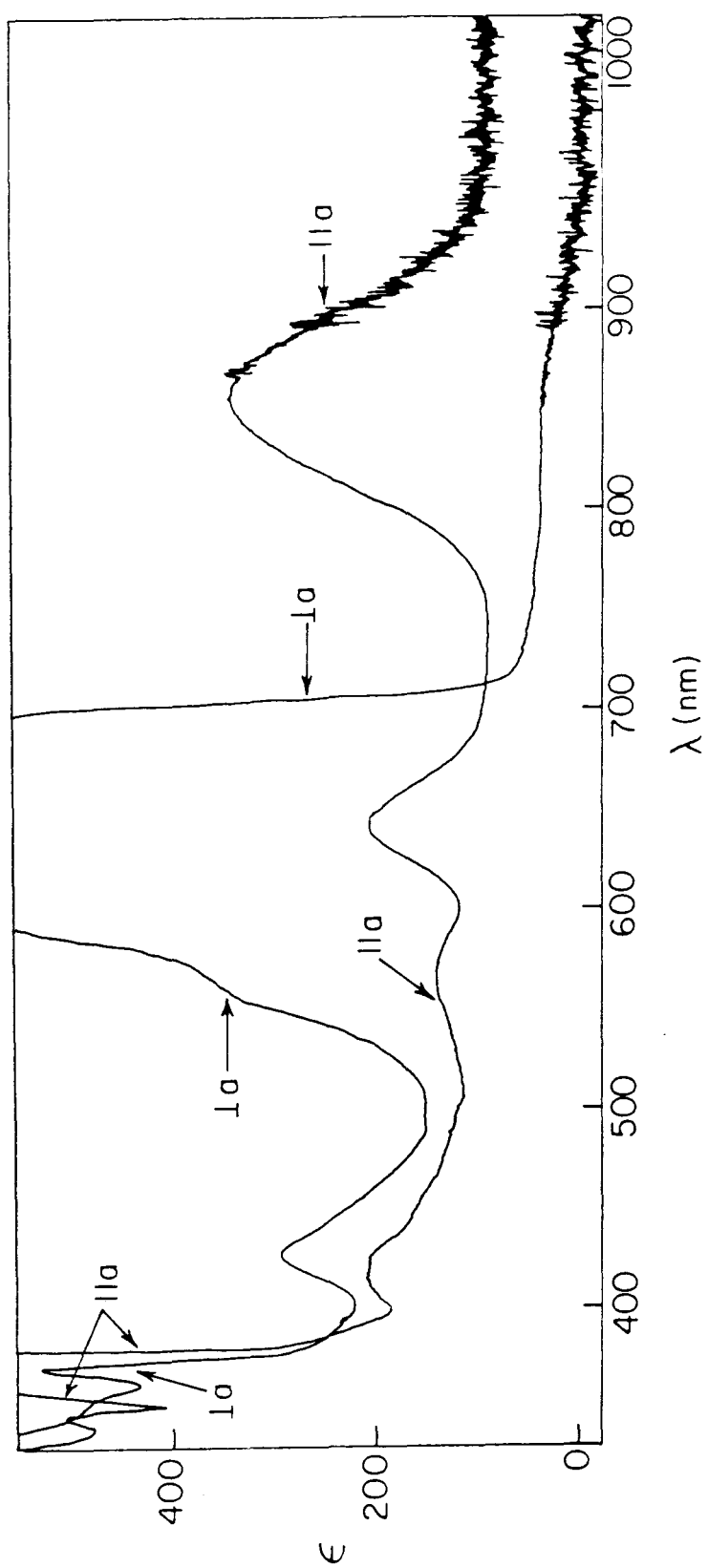
Upon cooling to 24 K, the spectrum sharpens and reveals additional features not resolved at room temperature (Figure 3.8). To the high-energy side of the  $A_{2u}(^1A_{2u})$  band, a feature assignable to one quantum of  $\nu(\text{NC})$  ( $2100 \text{ cm}^{-1}$ ) is observed in both polarizations ( $\lambda_{\text{max}}$  565 nm). (Modes of  $e_g$  and  $a_{1g}$  symmetries presumably lead to  $\perp z$  and  $\parallel z$  intensities, respectively.) This feature appears as a broad shoulder in the x,y-polarized room-temperature spectrum. The weak feature at 640 nm in the  $\parallel a$  spectra at both room temperature and 24 K is interpreted in terms of a vibronically induced  $\perp z$  intensity component of the transition to  $A_{2u}(^1A_{2u})$ , presumably involving low-frequency modes of  $e_g$  symmetry.<sup>48</sup> The observation that the intensity of this feature decreases significantly as the temperature is reduced to 24 K is consistent with the vibronic interpretation. The band assigned to the transition to  $E_u(^3A_{2u})$  sharpens and red-shifts, as expected because of the higher-frequency metal-metal vibration in the excited state. Similar to the  $A_{1g}(^1A_{1g}) \rightarrow E_u(^3A_{2u})$  band in the spectrum of  $\text{Rh}_2(\text{TMB})_4^{2+}$ , but in contrast to  $\text{Rh}_2\text{b}_4^{2+}$ , no vibronic structure is resolved at 24 K.<sup>1</sup> The metal-metal vibration (which is known from a Raman study to have values for the  $\text{Ir}_2$  compound of 55 and  $132 \text{ cm}^{-1}$  in the ground and  $^3A_{2u}$  excited states, respectively)<sup>12</sup> is expected to be the primary promoting mode for this electronic transition. The absence of a resolved vibronic progression built on the electronic

**Figure 3.7.** Polarized single-crystal absorption spectra of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  at 25 °C. The  $\parallel a$  base line is vertically offset by 20  $\epsilon$  units.





**Figure 3.8.** Polarized single-crystal absorption spectra of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  at 24 K. The  $\parallel a$  base line is vertically offset by 100  $\epsilon$  units.



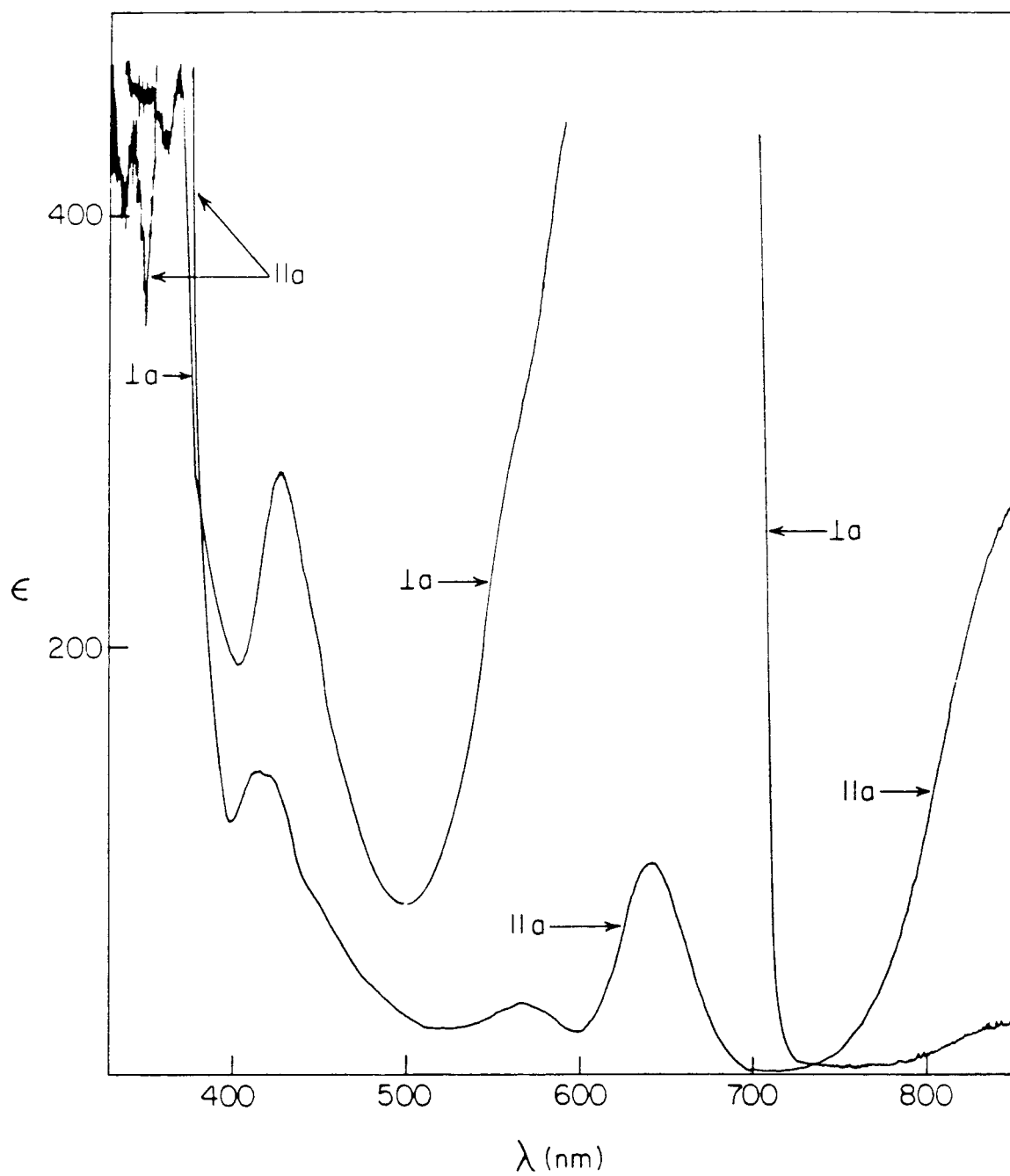
absorption band may be due to broadening attributable to vibronic coupling involving a low-energy ligand mode.

Several well-resolved bands are observed at higher energy than the  $d\sigma^* \rightarrow p\sigma$  absorption. The intense feature at 371 nm displays x,y-polarization, as expected for the  $A_{1g}(^1A_{1g}) \rightarrow E_u(^3E_u)$  system. The weak band at 344 nm, resolved in both room-temperature and 24 K spectra, is x,y-polarized; it represents one quantum of  $\nu(\text{NC})$  ( $1970\text{ cm}^{-1}$ ) built on the transition to  $E_u(^3E_u)$ . The sharp band at 369 nm in the z-polarized spectrum is assigned to  $A_{1g}(^1A_{1g}) \rightarrow E_u(^3E_u)$ ; it appears in this polarization because of the slight misalignment of the chromophore in the crystal. The shoulder to higher energy is assigned to one quantum of  $\nu(\text{NC})$ , presumably an  $e_g$  mode in this polarization.

Unambiguous polarizations for the higher-energy transitions could not be obtained because available crystals were too thick to allow maxima to be measured in either polarization. Indirect polarizations can be obtained from the room-temperature spectra (Figure 3.7), which could be extended to higher optical density than the low-temperature spectra. The band at 328 nm appears to be z-polarized, since intensity in the  $\perp a$  polarization goes off scale near 340 nm well before it does in the  $\parallel a$  polarization, but then comes back on scale at  $\sim 318$  nm. This polarization is consistent with the  $A_{1g}(^1A_{1g}) \rightarrow A_{2u}(^3E_u)$  assignment. The intense band(s) to higher energy are similarly indicated to be x,y-polarized, since the absorption  $\perp a$  remains on scale down to nearly 290 nm. These observations are consistent with our assignments of the intense solution bands at 318 and 305 nm as the x,y-polarized transitions  $A_{1g}(^1A_{1g}) \rightarrow E_u(^1E_u)$  and  $A_{1g}(^1A_{1g}) \rightarrow E_u(^3B_{1u})$ .

Several weak features are resolved to lower energy of the  $A_{1g}(^1A_{1g}) \rightarrow E_u(^3E_u)$  band. Low-temperature single-crystal spectra of a thicker crystal of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2 \cdot \text{CH}_3\text{C}_6\text{H}_5$  are shown in Figure 3.9. The weak features in the 400- 500 nm region are better resolved and show distinct polarizations. Similar

**Figure 3.9.** Polarized single-crystal absorption spectra of a thicker crystal of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  at 22 K.



features have been reported for the analogous Rh complex.<sup>1</sup> Because of drastic thermal narrowing of the intense bands to higher and lower energy (Figures 3.7 and 3.8), the temperature dependence of the intensity of these bands is obscured. The most intense band, which peaks at 428 nm, is predominantly z-polarized. In x,y-polarization, the shape of the absorption is reproducibly very different, and we infer additional weak shoulders at 415 and 450 nm. The absorption maxima in the polarized single-crystal spectra of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  are set out in Table 3.6.

### Assignments

Many possible assignments have been considered for the weak features observed at lower energy than the  $A_{1g}(^1A_{1g}) \rightarrow E_u(^3E_u)$  band. One initially attractive assignment, suggested for the  $\text{Rh}_2$  complexes, was excitations into the  $d\sigma^*$  (Rh-C) orbitals derived from  $\text{Rh}(d_{x^2-y^2})$ ,<sup>49</sup> as excited states of this type have been postulated to be thermally accessible from  $^3A_{2u}(d\sigma^* \rightarrow p\sigma)$ .<sup>50</sup> However, the halfwidths of these bands are far too narrow for transitions to  $d_{x^2-y^2}$ -derived orbitals, because the excited state in question would be strongly distorted along  $\nu(\text{M-C})$  coordinates.

The  $d_{xy} \rightarrow p_z$  transitions ( $d\delta, \delta^* \rightarrow p\sigma$ ), which are dipole-forbidden for both mononuclear and binuclear  $d^8$  complexes, were viewed as another attractive possibility for the weak features.<sup>1</sup> The assignment of the 291 nm band in the spectrum of  $[\text{Ir}(\text{CN-}t\text{-butyl})_4]\text{BF}_4$  to a transition to  $E_u(^3B_{1g})$  (derived from  $d_{xy} \rightarrow p_z$ ) indicates that it is not likely that the weak feature at 430 nm represents an analogous excitation. The  $d_{xy} \rightarrow p_z$  transitions evidently are only weakly perturbed, as are the  $d_{xz,yz} \rightarrow p_z$  excitations.

Another possibility has been suggested for the weak features. For  $\text{Rh}_2\text{b}_4^{2+}$  and  $\text{Rh}_2(\text{TMB})_4^{2+}$ ,  $d\sigma \rightarrow p\sigma$  and  $d\sigma^* \rightarrow p\sigma^*$  excitations could be responsible for these bands.<sup>1</sup> (The relevant excited state,  $d\sigma^*p\sigma^* + d\sigma p\sigma$ , may not be strongly bound (*vide infra*) and therefore is best described as an excited monomer weakly coupled to a ground-state

**Table 3.6.** Absorption maxima for polarized single-crystal spectra of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$ .

$^1\text{A}_{1g} \rightarrow$	pol.	at 25 °C		pol.	at 24 K	
		$\lambda_{\text{max}}$ , nm ( $\epsilon$ , $\text{M}^{-1}\text{cm}^{-1}$ )	fwhm, $\text{cm}^{-1}$		$\lambda_{\text{max}}$ , nm ( $\epsilon$ , $\text{M}^{-1}\text{cm}^{-1}$ )	fwhm, $\text{cm}^{-1}$
$\text{E}_\text{u}(\text{}^3\text{A}_{2\text{u}})$ , $\text{d}\sigma^* \rightarrow \text{p}\sigma$	x,y	830 (180, $\parallel a$ ; 60, $\perp a$ )	2200	x,y	855 (250, $\parallel a$ ; 40, $\perp a$ )	1300
$^1\text{A}_{2\text{u}}$ , $\text{d}\sigma^* \rightarrow \text{p}\sigma$	z	645 (140, $\parallel a$ ; $\perp a$ <sup>a</sup> )		z	642 (110, $\parallel a$ ; $\perp a$ <sup>a</sup> )	1100
+ $\nu(\text{CN})$	z; x,y	565		z; x,y	568 (50, $\parallel a$ )	
a	x,y	~455 (40 sh, $\parallel a$ )		x,y	~450 (65 sh, $\parallel a$ )	
$^1\text{A}_{1g}$ , $\text{d}\sigma/\sigma^* \rightarrow \text{p}\sigma/\sigma^*$	z > x,y	430 (~300, $\perp a$ )	1400	z > x,y	428 (300, $\perp a$ )	1500
a				x,y	415	
$\text{E}_\text{u}(\text{}^3\text{E}_\text{u})$ , $\text{d}_{\text{xz,yz}} \rightarrow \text{p}_\text{z}$	x,y	371		x,y	369	
+ $\nu(\text{CN})$	z; x,y	345 (sh)		z; x,y	344	
$\text{A}_{2\text{u}}(\text{}^3\text{E}_\text{u})$ , $\text{d}_{\text{xz,yz}} \rightarrow \text{p}_\text{z}$	z	~328				
$\text{E}_\text{u}(\text{}^1\text{E}_\text{u})$ , $\text{d}_{\text{xz,yz}} \rightarrow \text{p}_\text{z}$	x,y	>320				

a. Off-scale.

b. Not assigned, see text.



monomer.) Such an assignment is quite appealing for  $\text{Ir}_2(\text{TMB})_4^{2+}$ . The spectral bandwidth ( $1500\text{ cm}^{-1}$ ) is too great to correspond to a  $d\pi \rightarrow p\sigma$  transition, but is similar to the bandwidths associated with  $d_z^2 \rightarrow p_z$  excitations. Such transitions are dipole-forbidden; hence, the absorptions should be weak, with intensities induced by vibronic or crystal site-symmetry couplings. An interesting consequence of such an assignment is the correspondence of the 430 nm band in the spectrum of  $\text{Ir}_2(\text{TMB})_4^{2+}$  to the  $A_{1g}(^1A_{1g}) \rightarrow A_{2u}(^1A_{2u})$  ( $d_z^2 \rightarrow p_z$ ) absorption at 423 nm for  $\text{Ir}(\text{CN-}t\text{-butyl})_4^+$ . A similar correspondence between dimer and monomer band positions has been noted previously for the  $d\pi \rightarrow p_z$  and  $d_{xy} \rightarrow p_z$  excitations.

Note that only one  $^1A_{1g}(d_z^2 \rightarrow p_z)$  transition is expected near 430 nm, since the other one (MO model) is to an ionic transition in the VB model (Figure 3.3).

Possible assignments for the additional weak features near 430 nm, as suggested by the VB model, are transitions to  $^1,^3E_g$  bound states arising from  $d\pi/\pi^* \rightarrow p\sigma/\sigma^*$  excitations. Such features might be somewhat red-shifted from the monomer transitions (Figure 3.1) and have widths comparable to the bound  $d\sigma^* \rightarrow p\sigma$  excitation. The low intensity would suggest spin-forbidden transitions, possibly  $^3(d\pi/\pi^* \rightarrow p\sigma/\sigma^*)$ .

If the  $d\sigma/\sigma^* \rightarrow p\sigma/\sigma^*$  assignment is correct, a corresponding band should be observed for the analogous Rh complexes. Indeed, for  $\text{Rh}_2(\text{TMB})_4^{2+}$  and  $\text{Rh}_2\text{b}_4^{2+}$ , bands at 375 and 385 nm, respectively, are observed, each with the appropriate width and polarization.<sup>46</sup> (The  $A_{2u}(^1A_{2u})$  transition for  $\text{Rh}(\text{CNet})_4^+$  is observed at 380 nm.)<sup>31</sup>

In addition to accommodating both weakly and strongly perturbed monomer transitions, the VB model identifies fully half of the dimer electronic excited states as ionic (MMCT). Earlier, we placed the MMCT transitions for  $\text{Ir}(\text{I})$  and  $\text{Rh}(\text{I})$  at approximately 240 and 200 nm, respectively. The correspondence between the predicted MMCT transition energies and the broad features at 270 and 240 nm for  $\text{Ir}_2(\text{TMB})_4^{2+}$  and  $\text{Rh}_2(\text{TMB})_4^{2+}$ , respectively, makes assignment of the observed bands

to such transitions very attractive. The A-term behavior in the MCD spectrum suggests transitions to  $^1,^3E_u$  ( $d_{xz,yz} \rightarrow p_z$ ) states. The peak separation observed for  $\text{Ir}_2(\text{TMB})_4^{2+}$  ( $2000\text{ cm}^{-1}$ ) is consistent with the expected singlet-triplet state splitting. For the covalent  $d_{xz,yz} \rightarrow p_z$  transition, peak separations of  $1400\text{ cm}^{-1}$  ( $A_{2u}(^3E_u)$ ) and  $4500\text{ cm}^{-1}$  ( $E_u(^3E_u)$ ) are observed. The large bandwidths are attributed to ligand distortions caused by the change in electron density around the metal. In analogy to  $\text{Ir}_2(\text{TMB})_4^{2+}$ , the high-energy features for  $\text{Rh}_2(\text{TMB})_4^{2+}$  are tentatively assigned to MMCT ( $d_{xz,yz} \rightarrow p_z$ ) transitions. While this is but one of the possible assignments (other possibilities are transitions to  $d_{x^2-y^2}$  states, which would have to be from  $d\pi/\pi^*$  levels in order to account for the A terms, or MLCT states), the large blue shift of the bands in going from Ir to Rh is consistent with the ionic nature of the transition, and is not obviously consistent with any other interpretation. In line with this discussion, the highest energy band for  $\text{Ir}_2(\text{TMB})_4^{2+}$ , which shows a weak -B term, may be attributable to an ionic  $^1A_{2u}$  state derived from a  $d\sigma/\sigma^* \rightarrow p\sigma/\sigma^*$  transition (e.g., the  $d\sigma \rightarrow p\sigma^*$  state of the MO limit).

We have not made any assignments to states derived from excitations to the  $d_{x^2-y^2}$  orbital. Such transitions, which are formally allowed for the dimer and have been assigned<sup>2</sup> for  $d^8$ - $d^8$   $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$ , are not expected to give rise to intense bands.

## Conclusions

The MO model accounts for some of the important aspects of the metal-metal interaction in binuclear  $d^8$  complexes, but it is unable to accommodate all of the observed spectral features. The VB "weak-coupling" model allows for dimer transitions that are highly perturbed from the monomer bands as well as monomer-like dimer excitations. The VB model also identifies half of the dimer electronic excited states as ionic, involving metal-to-metal charge transfer; these MMCT states do not correlate to any monomer states. With the VB model, we have been able to rationalize the energies

and widths of the previously observed bands; in addition, we have arrived at plausible assignments for bands that were either not interpreted satisfactorily or not observed in earlier work. There is a striking similarity of our picture of the  $d^8$ - $d^8$  states to that derived from the spectroscopy of van der Waals molecules. The VB model does not change any of the MO-based interpretations of the thermal chemistry, photochemistry, or photophysics of these systems. What it does do is emphasize that weakly bound ground-state systems will show just a few strongly stabilized excited states that correlate to monomer states, and these are the only ones that MO theory describes correctly. The weakly perturbed monomer-like states, and, particularly, the ionic states, are better described by the VB model.

The VB model is applicable to the general class of dimeric systems in which the metals are nonbonded, in a formal sense, and can be viewed as weakly coupled (e.g.,  $[\text{Ir}(\mu\text{-pz})(\text{COD})]_2$ ,<sup>51</sup> binuclear  $d^8$  A-frame;  $\text{Au}_2(\text{dmpm})_2^{2+}$ ,<sup>52</sup> binuclear  $d^{10}$  complex;  $\text{Pt}_2(\text{dppm})_3$ ,<sup>53</sup> binuclear  $d^{10}$  complex). In all of these cases, minimal perturbations of monomer states are expected except when collapse to a strongly bonded state can occur.

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22. Determined from the ratio of the integrated intensities for toluene- $\text{CH}_3$  resonance and aromatic proton resonances, to the integrated intensities of the TMB resonances. A signal corresponding to the methyl resonance of the  $\text{CH}_3\text{CN}$  was not seen.
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41. The higher energy  $^1A_{2u}$  state (" $d\sigma \rightarrow p\sigma^*$ ") obviously does not demand a shorter metal-metal bond. However, its extreme antibonding character mandates a higher energy, as supported by calculations on analogous states of Hg<sub>2</sub>.
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## Chapter 4

**Thermal Reactivity of  $\text{Ir}_2(\text{TMB})_4^{2+}$**

## Introduction

Binuclear  $d^8$  complexes are known to undergo two-electron oxidative addition reactions similar to mononuclear  $d^8$  complexes.<sup>1-9</sup> Unlike the mononuclear systems in which the metal is oxidized by two electrons,<sup>10</sup> the dimers undergo addition reactions in which the metals are oxidized by one electron with formation of a metal-metal single bond. Interest in the addition reactions of binuclear complexes arises from two sources. In an effort to characterize the excited-state chemistry of a system, it is necessary to understand the thermal ground-state reactivity. Second, it is thought that binuclear complexes are more attractive for activating organic molecules than mononuclear complexes because of the possibility of cooperative involvement of metal sites.<sup>11-15</sup>

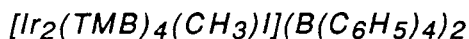
While much effort has been expended toward understanding the various oxidative addition pathways for mononuclear systems, little is known for bimetallic complexes. Some general statements can be made by considering the electronic structure and the three-dimensional structure of these systems. Similar arguments for the orbital involvement in addition reactions can be made for the dimers as have been made for the mononuclear complexes.<sup>10</sup> This is not necessarily apparent within the MO formalism; however, within the VB model, where the ground state is best described as two weakly coupled monomers, this is more apparent. For some dimer complexes, the metals have been observed to react independently, mimicking the monomer reactivity.<sup>16</sup> One of the major considerations in the reactivity of bimetallic complexes is the influence of the three-dimensional structure on the reactivity. For a cage-like complex, concerted addition to one or both metal centers is unlikely. For an open A-frame complex, concerted additions to one or both metals are possible and have been observed.<sup>16-20</sup> One added feature that bimetallic systems possess, which may enhance their reactivity in comparison to monomers, is the formation of a metal-metal bond in the final product. It has been suggested that the increased enthalpy of reaction for oxidative addition may be attributed to formation of a metal-metal bond.<sup>7</sup>

Earlier work found that  $\text{Rh}_2\text{b}_4^{2+}$ ,  $\text{Rh}_2(\text{TMB})_4^{2+}$ ,  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$ , and  $[\text{Ir}(\mu\text{-pz})\text{COD}]_2$  react with dihalides to yield the corresponding  $d^7\text{-}d^7$  dimer,  $\text{X-M-M-X}$ .<sup>1,2,9</sup>  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$  and  $[\text{Ir}(\mu\text{-pz})\text{COD}]_2$  were also found to react with  $\text{CH}_3\text{I}$  to yield  $\text{CH}_3\text{-M-M-I}$ .<sup>2,9</sup> It was found that  $\text{Ir}_2(\text{TMB})_4^{2+}$  reacts with dihalides ( $\text{Cl}_2$ ,  $\text{Br}_2$ ,  $\text{I}_2$ ),  $\text{CH}_3\text{I}$ ,  $\text{HCl}$ , and  $\text{CH}_2(\text{CN})_2$  to yield the corresponding  $d^7\text{-}d^7$  complex,  $\text{X-M-M-Y}$  ( $\text{X} = \text{Y}$ ,  $\text{Cl}_2$ ,  $\text{Br}_2$ ,  $\text{I}_2$ ;  $\text{X} = \text{I}$ ,  $\text{Y} = \text{CH}_3$ ;  $\text{X} = \text{Cl}$ ,  $\text{Y} = \text{H}$ ;  $\text{X} = \text{CH}(\text{CN})_2$ ,  $\text{Y} = \text{H}$ ).<sup>3</sup> Because of the greater reactivity expected for Ir and the interest in the photochemical reactivity of the excited states of  $\text{Ir}_2(\text{TMB})_4^{2+}$ , extension of the thermal chemistry was pursued.

## Experimental

### Synthesis

All synthetic procedures were carried out with standard Schlenk techniques. All solvents were from freshly opened bottles with no further purification. All solvents were degassed prior to use. Standard procedures were used to prepare  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$ .<sup>21</sup> Methyl iodide was distilled prior to use. All other chemicals were of reagent grade or of comparable quality and were used as received. The  $^1\text{H}$  NMR spectra were obtained on a 400 MHz JNM-GX400 FT NMR spectrometer.



To a  $\text{CH}_3\text{CN}$  solution of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$ , a slight excess of freshly distilled  $\text{CH}_3\text{I}$  was added via syringe. The color of the solution changed immediately upon addition, from clear blue to clear yellow. The solution was stirred for 10 minutes. Solvent and excess  $\text{CH}_3\text{I}$  were removed under vacuum to yield an orange-yellow powder that was recrystallized from acetone to yield yellow crystalline material. A red-orange emission was observed from the recrystallized material at 77 K in the solid state. Calculated for  $[\text{Ir}_2(\text{TMB})_4(\text{CH}_3)\text{I}](\text{B}(\text{C}_6\text{H}_5)_4)_2$ : C, 58.7; H, 5.9; N, 6.1. Calculated for  $0.85[\text{Ir}_2(\text{TMB})_4(\text{CH}_3)\text{I}](\text{B}(\text{C}_6\text{H}_5)_4)_2$  and  $0.15[\text{Ir}_2(\text{TMB})_4(\text{I})_2](\text{B}(\text{C}_6\text{H}_5)_4)_2$ : C, 58.1; H, 5.8; N, 6.1. Calculated for  $[\text{Ir}_2(\text{TMB})_4(\text{CH}_3)\text{I}](\text{B}(\text{C}_6\text{H}_5)_4)_2 \cdot 2\text{CH}_3\text{COCH}_3$ : C, 58.9; H, 6.2; N, 5.8. Found: C, 58.8; H, 6.1; N, 5.7. Sample dried in vacuum: C, 58.3, H, 5.9; N, 6.1.  $^1\text{H}$  NMR:  $\delta(\text{CD}_3\text{CN}, 20^\circ\text{C})$ ; 0.95 (singlet,  $\text{CH}_3$ , 3H), 1.53 (broad singlet,  $\text{CH}_2$ , 16H), 2.09 (singlet,  $\text{CH}_3$ , 48H), 6.83 (triplet-of-triplets, 8H), 6.98 (triplet, 16H), 7.25 (multiplet, 16H).

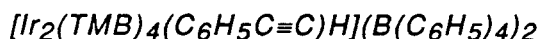


To a clear, blue acetonitrile solution of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$ , 1.4 equivalents of acetylchloride was added via syringe. An immediate color change was

observed, blue to yellow. The solution was stirred for 10 minutes. The final solution was a cloudy yellow. Solvent and excess acetylchloride were removed under vacuum to yield a light yellow powder. Attempts to recrystallize this material from acetone led to a rapid color change of the solution, yellow to blue-green.

### NMR Tube Reactions

Acetonitrile- $d_3$  was freeze-pump-thaw degassed, stored under vacuum over activated alumina, and vacuum-transferred into the NMR tube. Cyclohexadienes were distilled from  $\text{NaBH}_4$  under argon, freeze-pump-thaw degassed, stored under vacuum, and protected from light. 9,10-dihydroanthracene was recrystallized three times from absolute ethanol. A  $\text{CD}_3\text{CN}$  solution of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  was prepared in a grease-free/vacuum-adapted NMR tube. To this tube was added the desired substrates. Gases ( $\text{HCl}$  and  $\text{HI}$ ) were freeze-pump-thaw degassed or scrubbed ( $\text{H}_2$ ) with the appropriate catalyst prior to addition. Liquids were either added via syringe in an inert atmosphere ( $\text{R}_3\text{SnH}$  and  $\text{R}_3\text{SiH}$ ) or vacuum-transferred from an appropriate storage flask (1,4-cyclohexadiene and 1,3-cyclohexadiene). Solids ( $\text{CH}_2(\text{CN})_2$  and 9,10-dihydroanthracene) were added to crystalline  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  and the mixture dissolved in  $\text{CD}_3\text{CN}$ .



To the  $\text{CD}_3\text{CN}$  solution of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  (0.005 g) was added via syringe an excess of phenylacetylene. The solution was agitated at room temperature overnight.  $^1\text{H}$  NMR:  $\delta(\text{CD}_3\text{CN}, 20^\circ\text{C})$ ; -10.79 (singlet, Ir-H), 1.52 (broad singlet,  $\text{CH}_3$ ), 1.91 (broad singlet,  $\text{CH}_2$ ), 6.83 (triplet), 6.98 (triplet), 7.26 (multiplet).

*$[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  + 9,10-dihydroanthracene*

$[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  and freshly recrystallized 9,10-dihydroanthracene were combined in a grease-free/vacuum-adapted NMR tube. Acetonitrile- $d_3$  was vacuum-transferred onto the solids. The NMR tube was wrapped in foil and placed in the dark at room temperature overnight to allow the solid material to dissolve. The NMR spectrum of the sample was measured after 12 hrs. The sample, wrapped in foil, was placed in an 80 °C oil bath for 20 hrs.  $^1\text{H}$  NMR:  $\delta(\text{CD}_3\text{CN}, 20\text{ }^\circ\text{C})$ ; -10.62 (singlet, Ir-H); anthracene, 7.50 (multiplet, 4H), 8.05 (quartet, 4H), 8.50 (singlet, 2H).

### Spectroscopic Measurements

All solvents were dried and degassed by standard methods.<sup>22,23</sup> Absorption spectra were recorded with a Cary 17 spectrophotometer, a Hewlett-Packard 8450A spectrophotometer, or a Shimadzu UV-260 spectrophotometer. Spectra were obtained of solutions prepared on a high-vacuum line in a cell consisting of a 10 ml Pyrex bulb, a 1 cm pathlength quartz cell, and a Teflon vacuum valve. Several spectra of the complexes that are air-stable in solution were measured, of samples prepared in the air in 1 cm pathlength quartz cuvettes. Substrates were introduced to  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  solutions as described for the NMR tube reactions.

### Electrochemical Measurements

Cyclic voltametric measurements were conducted with a Princeton Applied Research (PAR) model 173 potentiostat/galvanostat, a model 175 universal programmer, and a model 179 digital coulometer. Cyclic voltammograms (CVs) were plotted on a Houston Instruments Omnigraphic 2000 x,y-recorder. Fast-sweep CVs were recorded on a Tektronix 5223 digitizing oscilloscope. The cell geometry used to measure the CVs is described elsewhere.<sup>24</sup> CVs were measured at a BAS Pt button working electrode, with a Pt wire auxiliary electrode, and a sodium saturated calomel

electrode (SSCE) as the reference. Dichloromethane was distilled from  $\text{CaH}_2$  under an atmosphere of argon, freeze-pump-thaw degassed, and stored under argon. Tetra-*n*-butyl ammonium hexafluorophosphate ( $\text{TBAPF}_6$ ) was recrystallized 3 times from 95% ethanol. CVs of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  were measured for dichloromethane solutions under argon that contained 0.1 M  $\text{TBAPF}_6$  as the supporting electrolyte and  $10^{-3}$  M compound.

### **Additional Measurements**

Elemental analyses were obtained at the CalTech Analytical Laboratory.



## Results and Discussion

### Thermal Reactivity

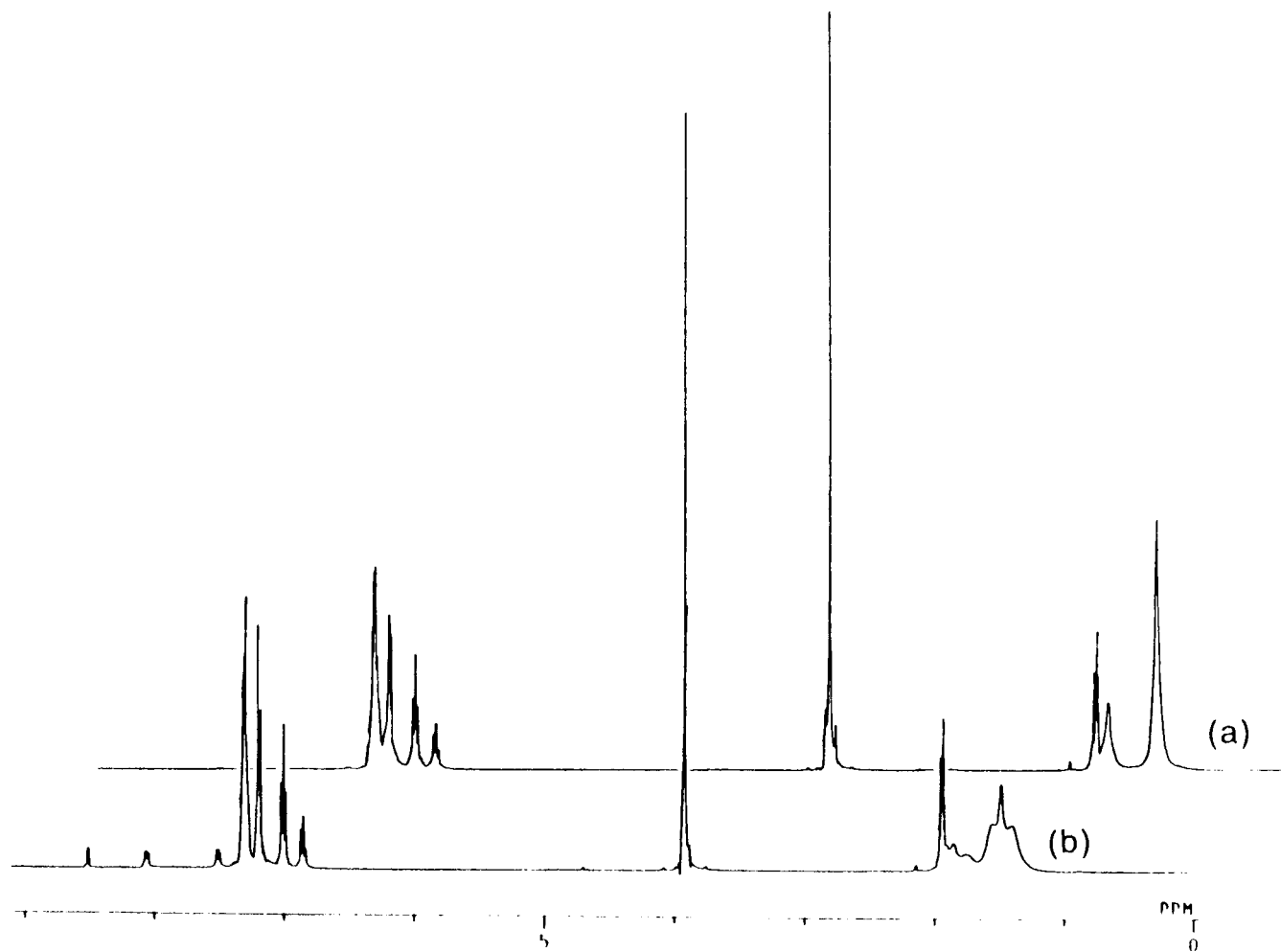
#### *Nonpolar and Low Polarity Reagents ( $H_2$ , C-H, Si-H, Sn-H)*

As was noted by T.P. Smith, no reaction with  $H_2$  is observed.<sup>3</sup> This is not surprising in that no reaction with  $H_2$  has been observed for the mononuclear alkyl isocyanide complexes. This has been understood to be a result of the decrease in electron density at the metal because of the  $\pi$  acidity of the isocyanide ligands. If phosphine is exchanged for isocyanide, as in the complex  $Rh_2(dppm)_2(CN-n\text{-hexyl})_4^{2+}$ , slow thermal addition of  $H_2$  is observed. In analogy to  $[Ir(CO)Cl(dppm)]_2$ ,<sup>16</sup>  $H_2$  addition is thought to proceed in a concerted fashion to one metal with rapid transfer of one hydride to the other metal center.

There are no known examples of concerted C-H addition to a binuclear  $d^8$  complex. For  $Ir_2(TMB)_4^{2+}$ , preliminary results indicate that thermolysis with 9,10-dihydroanthracene in  $CD_3CN$  at 80 °C yields products that are derived from C-H bond cleavage. Figure 4.1 shows the growth with time of anthracene and a metal product independently characterized as  $Ir_2(TMB)_4H_2^{2+}$ . The most likely pathway to the observed products is a free-radical process. Over a 12 hr period a 78% yield of anthracene, based on starting metal concentration, is obtained. This is the first example of a thermal C-H bond cleavage reaction by a binuclear  $d^8$  complex. There is no observable reaction with 1,4-cyclohexadiene or 1,3-cyclohexadiene at 25 or 80 °C.

For substrates described as carbon acids ( $CH_2(CN)_2$ ,  $pK_a \sim 11$ ;  $C_6H_5C\equiv CH$ ,  $pK_a \sim 18.5$ ),<sup>25</sup> reactions with  $Ir_2(TMB)_4^{2+}$  are observed and yield products best represented as  $H-Ir-Ir-R$ . Earlier work by T.P. Smith suggested that the reaction of malononitrile was not an acid/base reaction in that acetyl acetone ( $pK_a \sim 9$ ) and benzoic acid ( $pK_a \sim 10.7$ ) did not show similar reactivity.<sup>3</sup> Rigorous exclusion of light did not alter the reaction with malononitrile. It was thought that the prior coordination of the CN groups to the iridium center may be necessary for reaction, although no evidence to

**Figure 4.1.** NMR spectra of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  and 9,10-dihydroanthracene (a) prior to thermolysis, and (b) after 12 hrs thermolysis in 80 °C oil bath.



support this suggestion could be obtained. The lack of reactivity of acetyl acetone may be due to bad steric interactions with the methyl groups of the TMB ligand in the final product. Dimerization of benzoic acid or the instability of the final product are possible explanations of the lack of reactivity of this substrate.<sup>26</sup>

For  $\text{C}_6\text{H}_5\text{C}\equiv\text{CH}$ , a smooth reaction with  $\text{Ir}_2(\text{TMB})_4^{2+}$  is observed. Although the final product has not been isolated, the chemical shift of the Ir-H resonance and the nature of the aliphatic resonances lead to the suggestion that the product is  $\text{H-Ir-Ir-C}\equiv\text{CC}_6\text{H}_5$ . Similar to the reaction with malononitrile, the addition of phenylacetylene is thought to be an acid/base process. Homolytic cleavage of the C-H bond of phenylacetylene seems unlikely ( $D(\text{C}_6\text{H}_5\text{C}\equiv\text{C-H}) \sim 130 \text{ kcal/mol}$ ).<sup>27</sup>

Addition of acetylenes to Rh(I) and Ir(I) phosphine complexes has been recently studied.<sup>28</sup> It is suggested that these reactions rather than acid/base processes are concerted additions to the metal center. For  $\text{Ir}_2(\text{TMB})_4^{2+}$ , such a concerted addition is not likely, and an acid/base pathway is preferred. It may, in fact, be the case that for the mononuclear Rh(I) and Ir(I) phosphine complexes, the true reaction pathway is acid/base.

For  $\text{R}_3\text{SiH}$  and  $\text{R}_3\text{SnH}$ , rapid thermal reactions with  $\text{Ir}_2(\text{TMB})_4^{2+}$  are observed. In the optical absorption spectrum, a rapid bleach of the binuclear  $d^8$  spectrum is observed with growth of a spectrum that indicates oxidation of the metal complex, but not necessarily formation of a  $d^7$ - $d^7$  dimer complex. For  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$ , no thermal reaction with  $\text{R}_3\text{SiH}$  or  $\text{R}_3\text{SnH}$  is observed.<sup>29</sup> The NMR spectra of the reaction mixtures do reveal formation of metal hydrides.

#### *Electrophilic Reagents ( $\text{X}_2$ , $\text{RX}$ , $\text{HX}$ , $\text{CH}_3\text{COCl}$ )*

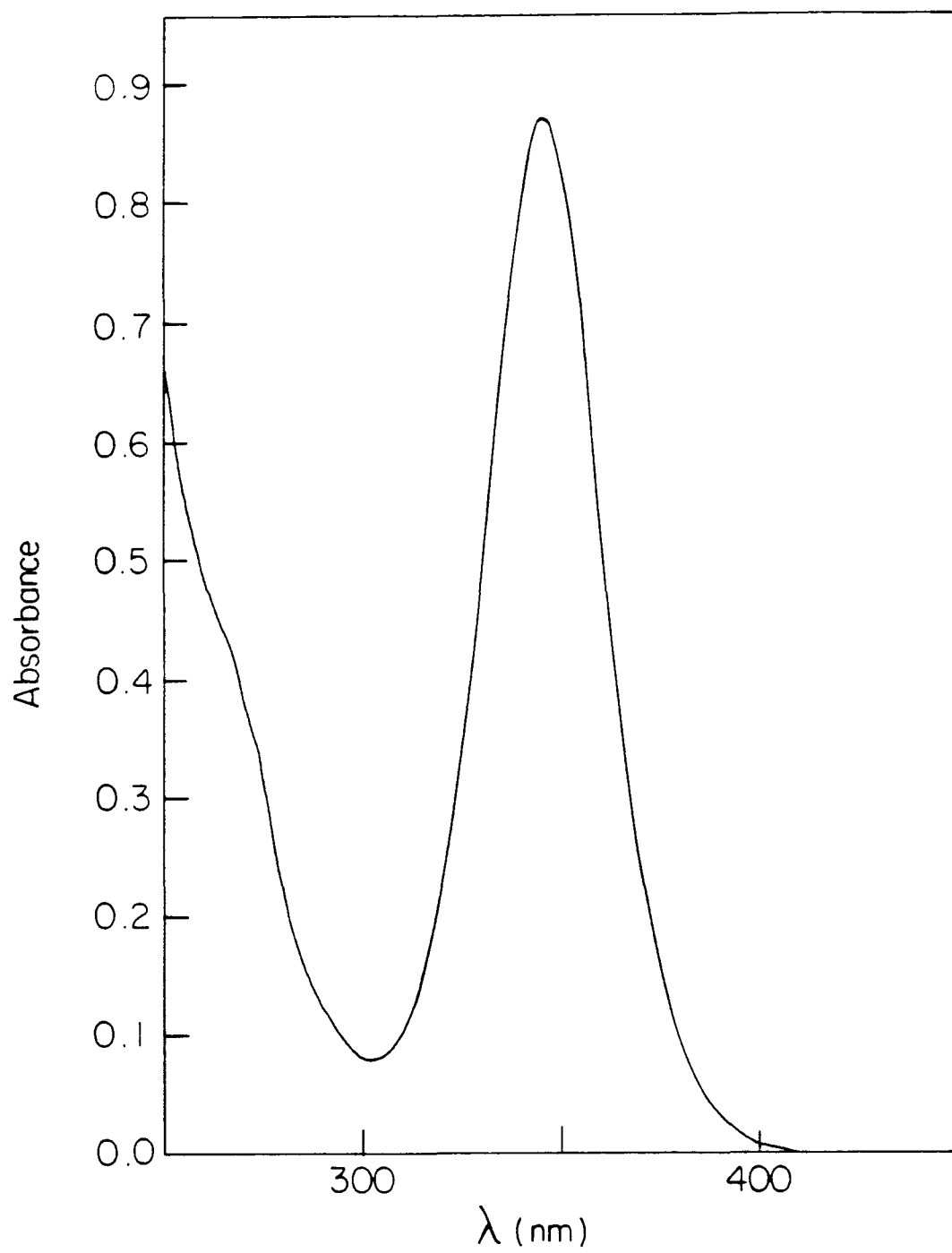
Previous work has shown that dihalides and mineral acids will add to  $\text{Ir}_2(\text{TMB})_4^{2+}$  oxidatively to yield the corresponding  $d^7$ - $d^7$  complex.<sup>3</sup> The only alkyl halide found to add thermally to  $\text{Ir}_2(\text{TMB})_4^{2+}$  is  $\text{CH}_3\text{I}$ . No thermal reaction is observed

with  $\text{CH}_2\text{Cl}_2$ ,  $\text{CHCl}_3$ ,  $\text{CCl}_4$ , or  $\text{ClCH}_2\text{CH}_2\text{Cl}$ . Previous discussion of the  $\text{CH}_3\text{I}$  reaction had concluded that the reaction follows a free-radical pathway. For analogous mononuclear  $d^8$  complexes, detailed studies have shown that the addition of  $\text{CH}_3\text{I}$ , unlike other alkyl halides, proceeds by an  $\text{S}_{\text{N}}2$  pathway,<sup>10</sup> contrary to the result for  $\text{Ir}_2(\text{TMB})_4^{2+}$ . Emphasized in the work on mononuclear complexes is the fine balance that exists between the concerted  $\text{S}_{\text{N}}2$  pathway and the electron-transfer, free-radical pathway. For  $\text{Ir}_2(\text{TMB})_4^{2+}$ , the balance may be swayed in favor of electron transfer.

Careful preparation and isolation of the addition products from the reaction of  $\text{CH}_3\text{I}$  and  $\text{Ir}_2(\text{TMB})_4^{2+}$  indicate that, unlike the earlier report, there is very little if any  $\text{Ir}_2(\text{TMB})_4\text{I}_2^{2+}$  ( $\text{Ir}_2\text{I}_2$ ) contamination. The  $\text{Ir}_2\text{I}_2$  by-product was used as evidence for a free-radical process. The UV-Vis spectrum of the isolated material (Figure 4.2) shows no evidence of  $\text{Ir}_2\text{I}_2$  contamination. The NMR spectrum of the isolated product appears quite clean. Neglecting the large uncertainty in the integrated intensities of the NMR resonances, an estimate of ~15%  $\text{Ir}_2\text{I}_2$  contamination is arrived at by comparing the ratio of  $\text{Ir-CH}_3$  to  $\text{TMB-CH}_2$ . However, from the combustion data a 15%  $\text{Ir}_2\text{I}_2$  impurity seems unlikely. Rather, the complex appears to be an acetone solvate. Thus, addition of  $\text{CH}_3\text{I}$  to  $\text{Ir}_2(\text{TMB})_4^{2+}$  appears to yield only  $\text{Ir}_2(\text{TMB})_4(\text{CH}_3)(\text{I})^{2+}$ .

These results do not distinguish between the possible pathways for addition of  $\text{CH}_3\text{I}$ . Alkyl halide addition to binuclear gold complexes is thought to proceed via an  $\text{S}_{\text{N}}2$  pathway.<sup>30,31</sup> Another possibility is an inner-sphere electron-transfer pathway in which a solvent-caged radical pair is formed. This pathway would yield products that are the same as that from an  $\text{S}_{\text{N}}2$  reaction. Evidence exists that suggests the addition of  $\text{I}_2$  to  $\text{Rh}_2\text{b}_4^{2+}$  follows an electron-transfer pathway.<sup>32</sup> No conclusion concerning the pathway of alkyl halide addition to  $\text{Ir}_2(\text{TMB})_4^{2+}$  is possible until more detailed experiments designed to test the various mechanism are done. Such experiments would be similar to those discussed for the mononuclear complexes.

**Figure 4.2.** Electronic absorption spectrum of  $[\text{Ir}_2(\text{TMB})_4(\text{CH}_3)\text{I}](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in  $\text{CH}_3\text{CN}$  at 25 °C.



Addition of acetylchloride to a solution of  $\text{Ir}_2(\text{TMB})_4^{2+}$  results in a rapid reaction with formation of a yellow material similar to that obtained from the reaction of  $\text{CH}_3\text{I}$ . The optical absorption spectrum of this material (Figure 4.3) indicates formation of a  $d^7$ - $d^7$  complex different from the previously prepared dichloride adduct and red-shifted from the spectrum expected for  $\text{Ir}_2(\text{TMB})_4(\text{CH}_3)(\text{Cl})^{2+}$ . The proposed product is  $\text{Ir}_2(\text{TMB})_4(\text{CH}_3\text{CO})(\text{Cl})^{2+}$ . Attempts to recrystallize the material lead to rapid decomposition. For analogous mononuclear  $\text{Ir}(\text{III})$  acyl complexes, decarbonylation occurs rapidly upon heating.<sup>10</sup> The yellow material is quite susceptible to decomposition, yielding blue-green materials upon heating, exposure to solvents such as acetone, and prolonged exposure to light.

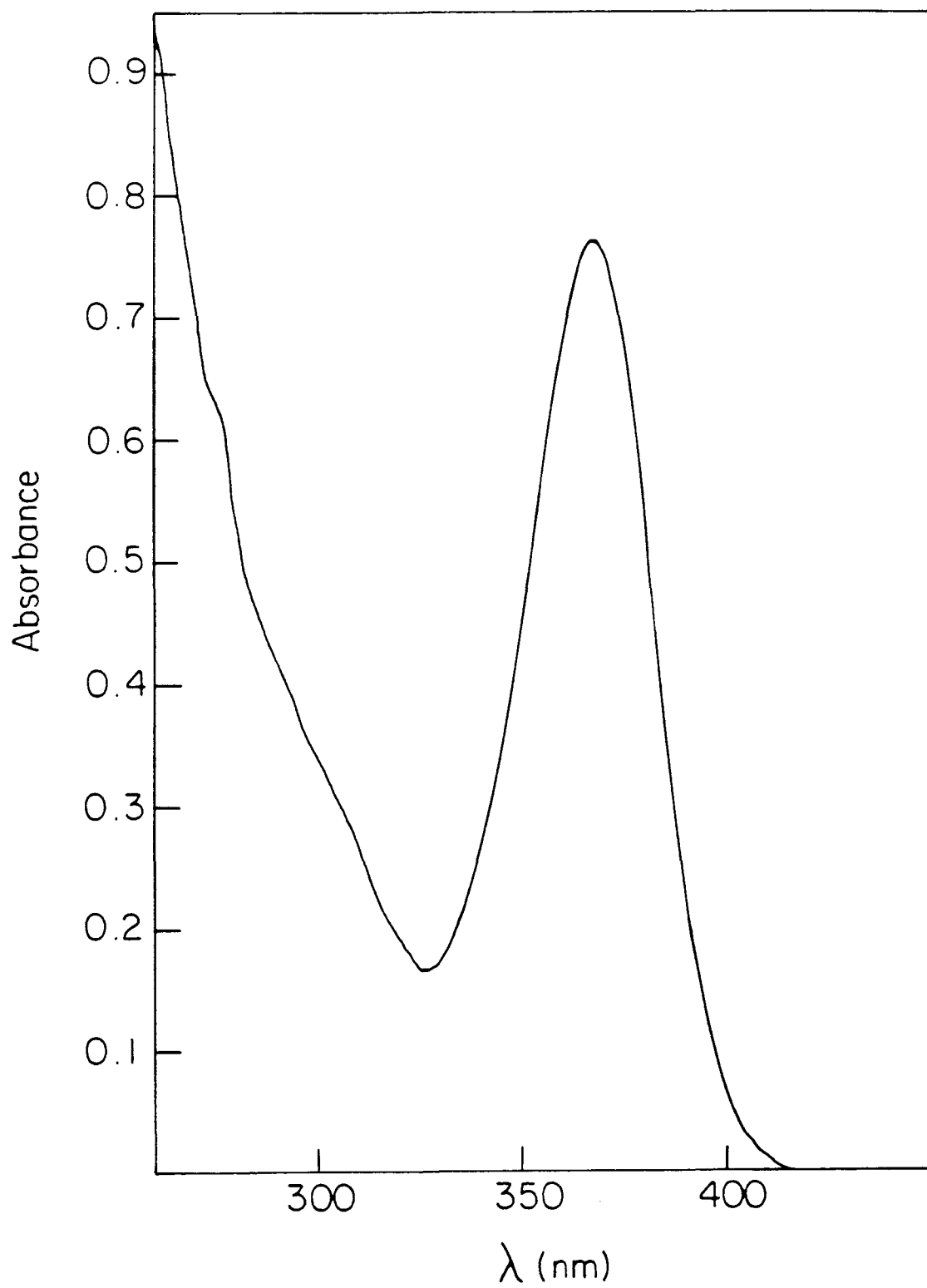
### Summary

The observed thermal reactivity for  $\text{Ir}_2(\text{TMB})_4^{2+}$  (Figure 4.4) conforms to the earlier expectations of net two-electron chemistry. While the details of the specific pathways for the various reactions are not known, the reactivity parallels that observed for the analogous mononuclear complexes. The reactivity of the monomers is controlled by the  $(d_z^2)^2$  configuration, with the reactivity of the dimer being dictated by the  $(d\sigma^*)^2$  configuration, or more accurately, two weakly coupled  $(d_z^2)^2$  metal centers. In contrast to the monomers, the metal centers of the dimer are oxidized by only one electron with formation of a metal-metal single bond. No complications that were due to reactions involving the isocyanide ligand were observed.<sup>33</sup>

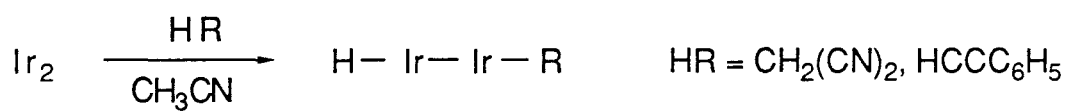
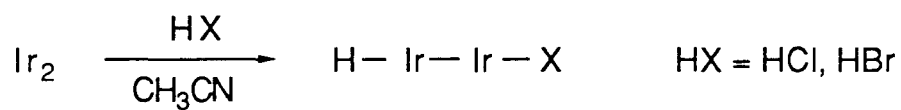
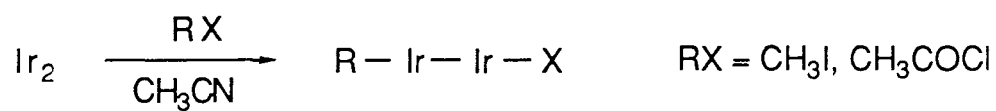
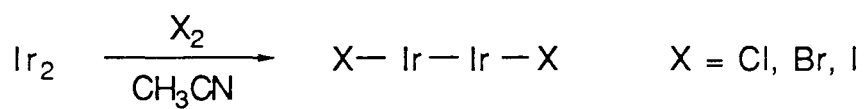
It will be shown in subsequent sections that the excited-state chemistry contrasts drastically to the ground-state chemistry. As stated earlier, the  $1,3(d\sigma^*p\sigma)$  excited states may be viewed as a diradical species, implying that the chemistry for such excited states will be one electron in nature, i.e., electron transfer and atom transfer, in contrast to the two-electron chemistry of the ground state. Free-radical reactions may occur for the ground state. These reactions involve oxidation of the metal complex,



**Figure 4.3.** Electronic absorption spectrum of  $[\text{Ir}_2(\text{TMB})_4(\text{CH}_3\text{CO})\text{Cl}](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in  $\text{CH}_3\text{CN}$  at 25 °C.



**Figure 4.4.** Observed thermal reactions of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$ .



generating a species whose electronic structure is similar to the triplet excited state, an oxidizing hole in the  $d\sigma^*$  orbital. It will be shown that electrochemical oxidation of  $\text{Rh}_2(\text{TMB})_4^{2+}$  to generate  $\text{Rh}_2(\text{TMB})_4^{3+}$  in the presence of alkyl halides and H-atom donors results in reactions that parallel those observed for the  $^3(d\sigma^*p\sigma)$  excited state.

### Ir-H Chemical Shifts

In Table 4.1, the Ir-H chemical shifts for a series of  $d^7$ - $d^7$  complexes of  $\text{Ir}_2(\text{TMB})_4^{2+}$  are listed. The unique feature of these complexes is the extreme sensitivity of the chemical shift of the Ir-H resonance to the nature of the ancillary group. A possible explanation for this behavior comes from the theory of Buckingham and Stephen in which the mixing of low-lying, electronic excited states into the ground-state wave function determines the hydride chemical shift.<sup>34,35</sup> For complexes of the type  $\text{YX}_4\text{MH}$ , the paramagnetic shielding,  $\sigma^P$ , is predicted to be very sensitive to the anisotropy of the complex. This is borne out for the limited data in Table 4.1.

In the Buckingham model, the relative chemical shifts for complexes of the type  $\text{YX}_4\text{MH}$  are governed by the ligand-field strength of the ligand Y; as the ligand-field strength of Y increases,  $\sigma^P$  should decrease. Within the limited data set for H-Ir-Ir-Y, the expected shift of the hydride resonance with increasing ligand-field strength of Y is observed. Electronic coupling of groups trans to one another through a metal-metal bond is known;<sup>36</sup> however, the trend in hydride chemical shift for the  $d^7$ - $d^7$  complexes of  $\text{Ir}_2(\text{TMB})_4^{2+}$  is surprising, considering the complexity of the electronic coupling of the metal centers. Further studies are necessary to determine how well the Buckingham model describes these systems.

### Electronic Spectra of $d^7$ - $d^7$ Complexes of $\text{Ir}_2(\text{TMB})_4^{2+}$

In Table 4.2, the optical absorption spectra for a series of  $d^7$ - $d^7$  complexes of  $\text{Ir}_2(\text{TMB})_4^{2+}$  are summarized. Figure 4.5 presents the spectra for several of these

**Table 4.1.** Ir-H Chemical Shifts.

H-Ir-Ir-X	$\delta$ (ppm)
H-Ir-Ir-Cl	-13.5
H-Ir-Ir-I	-13.1
H-Ir-Ir-CH(CN) <sub>2</sub>	-11.9
H-Ir-Ir-C $\equiv$ CC <sub>6</sub> H <sub>5</sub>	-10.8
H-Ir-Ir-H	-10.6

**Table 4.2.** Electronic absorption data for  $\text{Ir}_2(\text{TMB})_4\text{XY}^{2+}$  complexes.

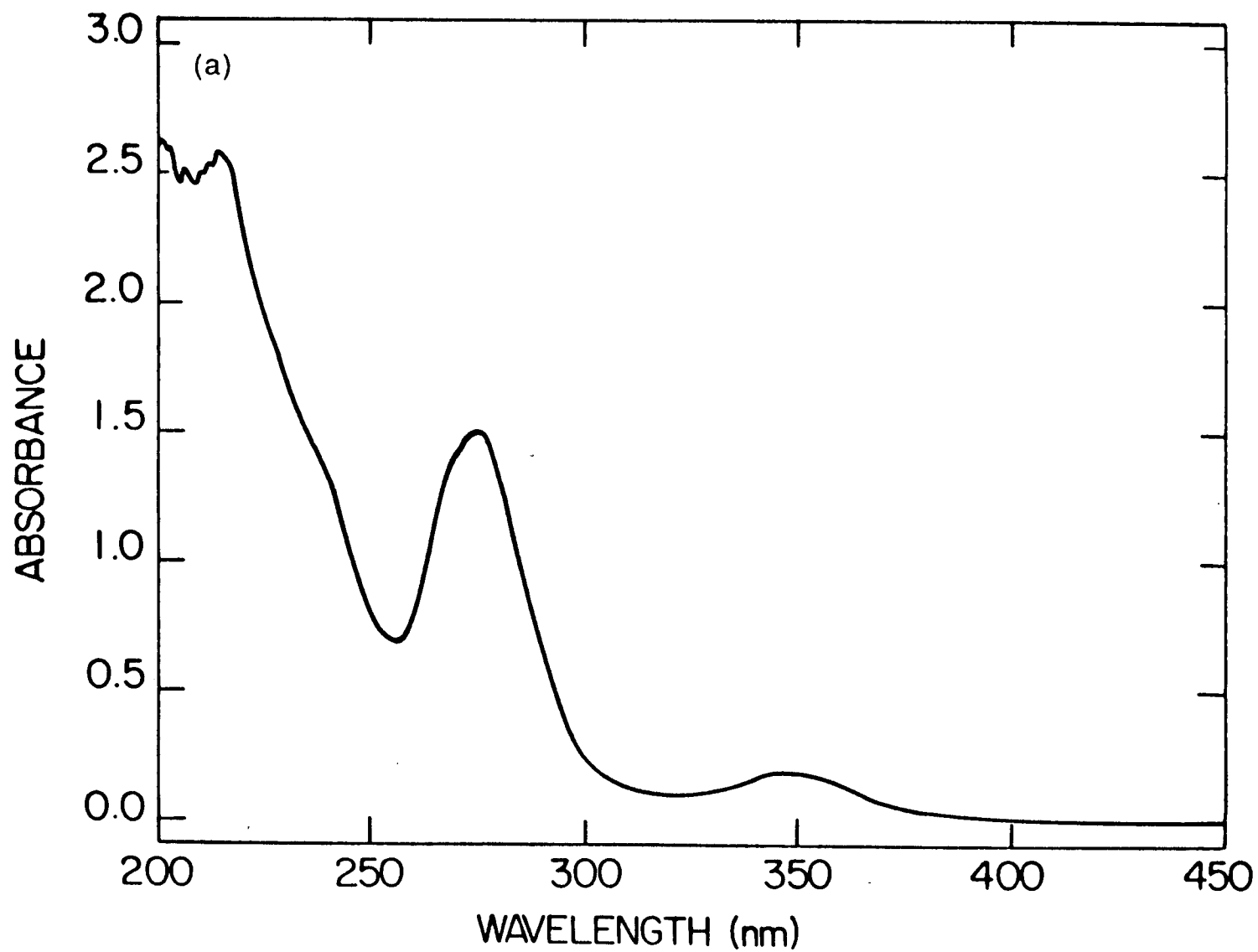
	$d\sigma \rightarrow d\sigma^*$ , nm ( $\epsilon$ , $\text{M}^{-1} \text{cm}^{-1}$ )	" $d\pi$ " $\rightarrow d\sigma^*$ , nm ( $\epsilon$ , $\text{M}^{-1} \text{cm}^{-1}$ )	Ref.
$\text{Ir}_2(\text{TMB})_4\text{Cl}_2^{2+}$	275 (50520)	349 (5910)	4
$\text{Ir}_2(\text{TMB})_4\text{Br}_2^{2+}$	297 (52569)	330 (4510)	4
$\text{Ir}_2(\text{TMB})_4\text{I}_2^{2+}$	330 (49800)	375 (8110)	4
$\text{Ir}_2(\text{TMB})_4(\text{H})\text{Cl}^{2+}$	336	a	3
$\text{Ir}_2(\text{TMB})_4(\text{H})(\text{CH}(\text{CN})_2)^{2+}$	330	a	b
$\text{Ir}_2(\text{TMB})_4(\text{CH}_3)\text{I}^{2+}$	346	a	b
$\text{Ir}_2(\text{TMB})_4(\text{CH}_3\text{CO})\text{Cl}^{2+}$	366	a	b
$\text{Ir}_2(\text{TMB})_4\text{H}_2^{2+}$	320	a	b

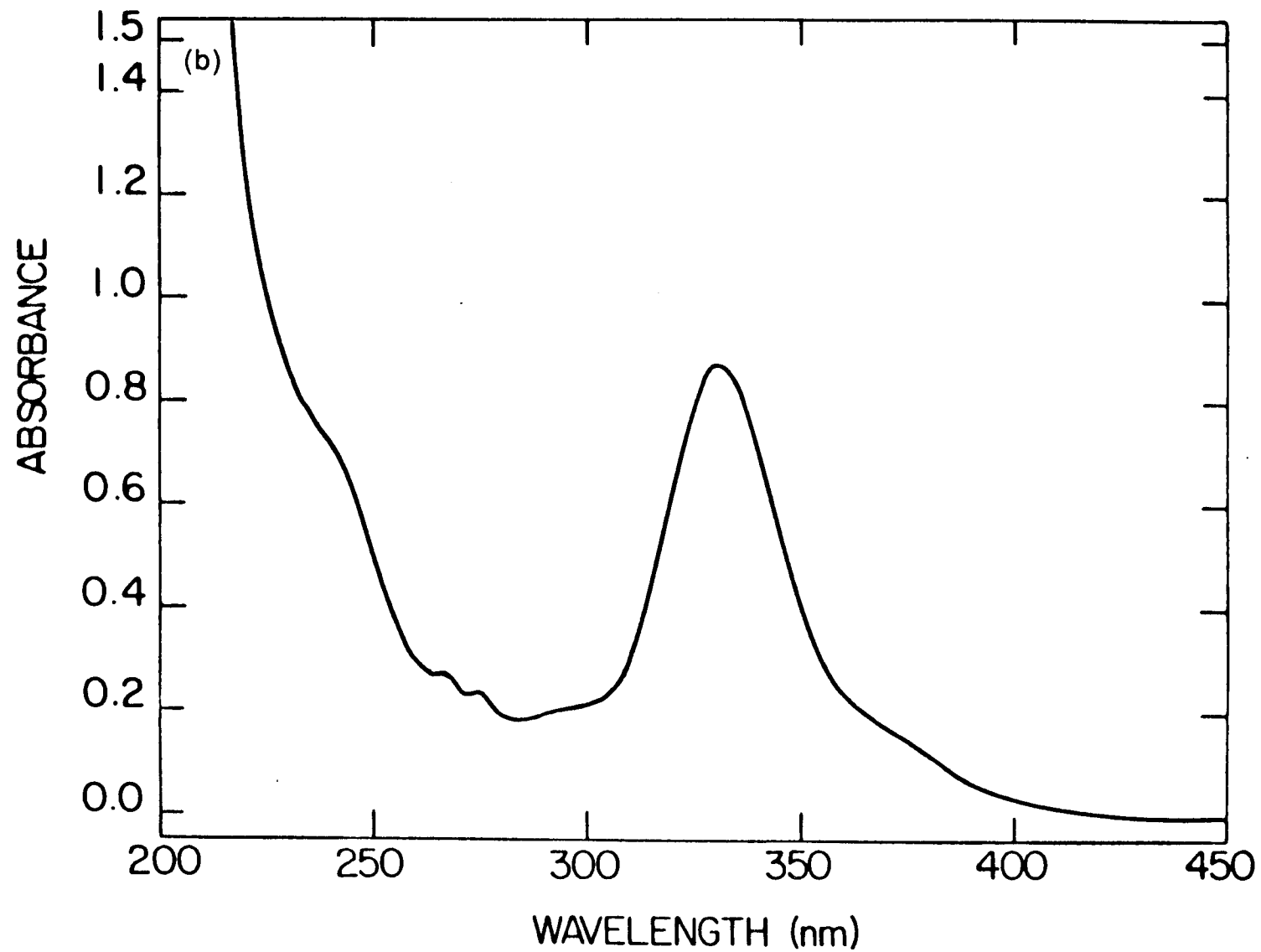
a. Not observed.

b. This work.

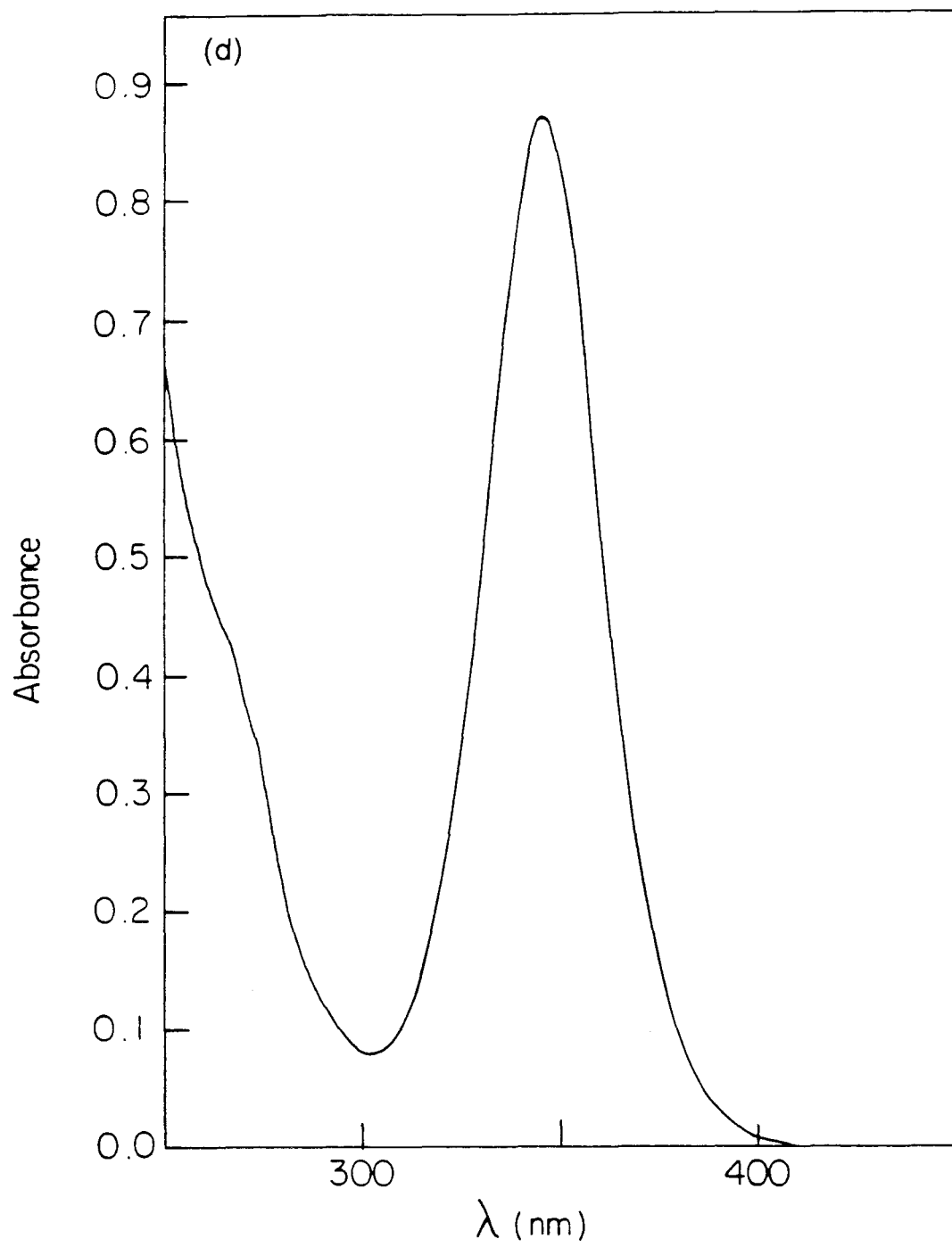
**Figure 4.5.** Electronic absorption spectra of (a)  $[\text{Ir}_2(\text{TMB})_4\text{Cl}_2](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in  $\text{CH}_3\text{CN}$ , (b)  $[\text{Ir}_2(\text{TMB})_4\text{I}_2](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in  $\text{CH}_3\text{CN}$ , (c)  $[\text{Ir}_2(\text{TMB})_4(\text{H})\text{Cl}](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in  $\text{CH}_3\text{CN}$ , (d)  $[\text{Ir}_2(\text{TMB})_4(\text{CH}_3)\text{I}](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in  $\text{CH}_3\text{CN}$ , (e)  $[\text{Ir}_2(\text{TMB})_4(\text{CH}_3\text{CO})\text{Cl}](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in  $\text{CH}_3\text{CN}$ , and (f)  $[\text{Ir}_2(\text{TMB})_4\text{H}_2](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in  $\text{CH}_3\text{CN}$ .

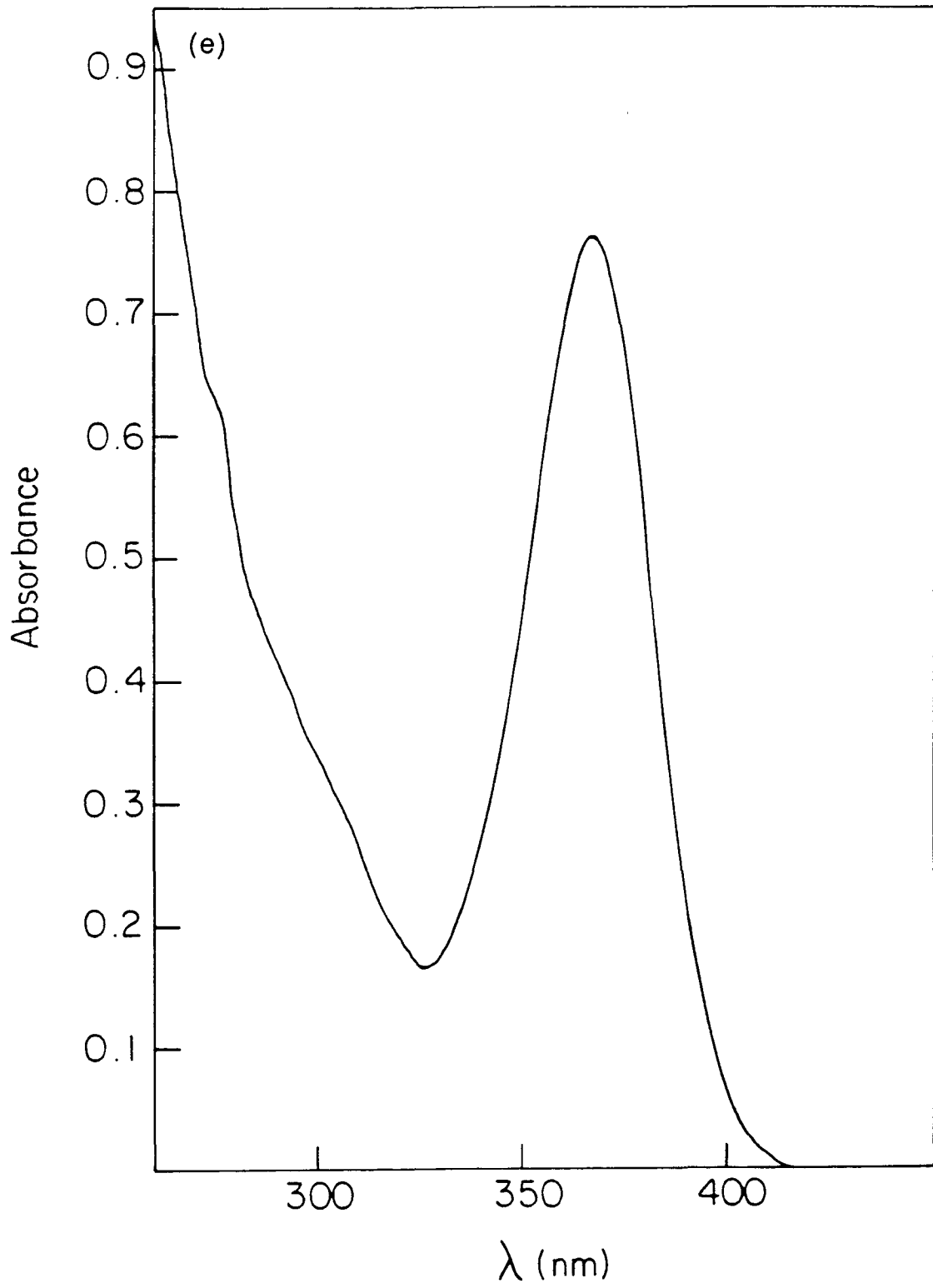


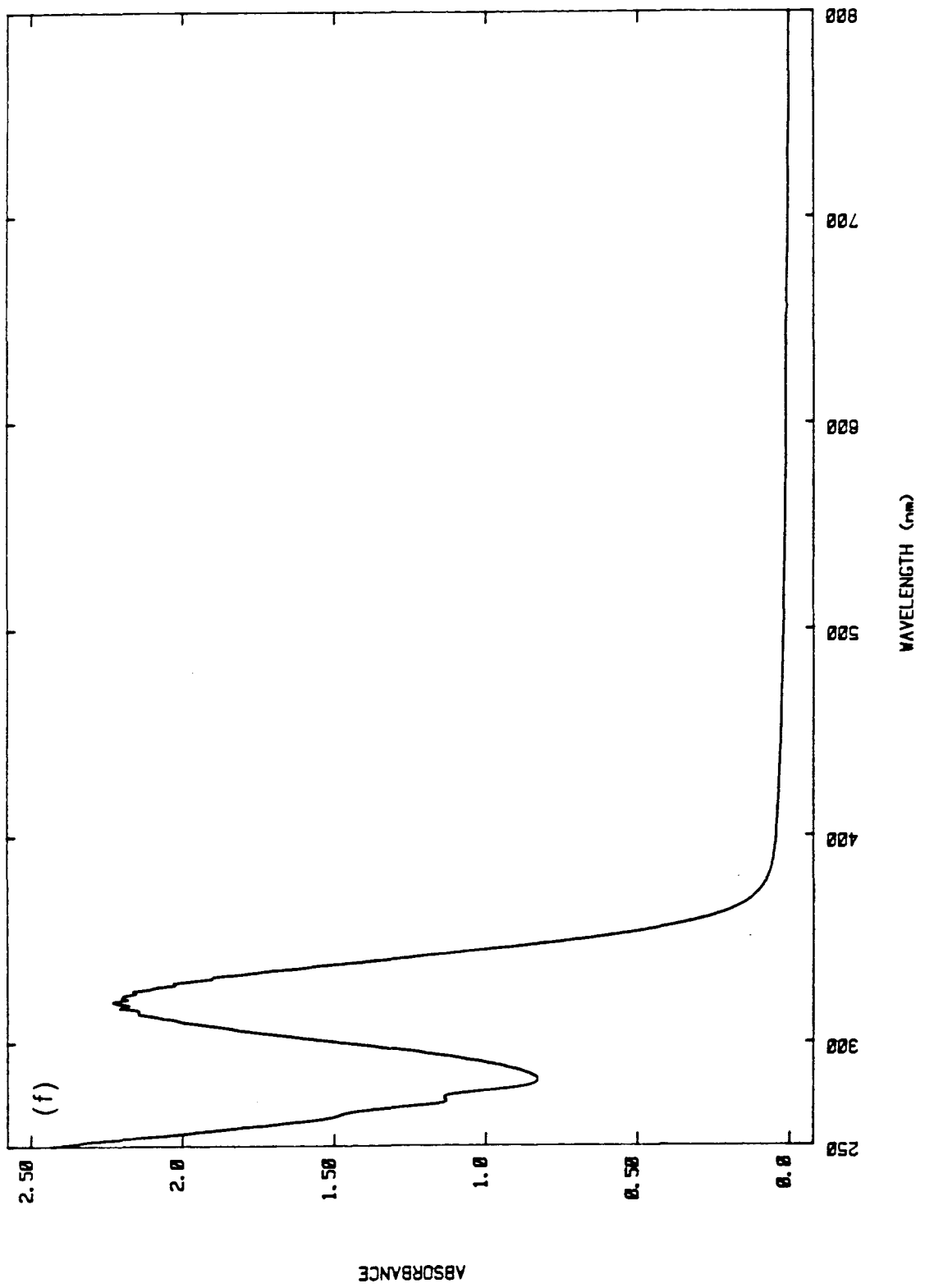












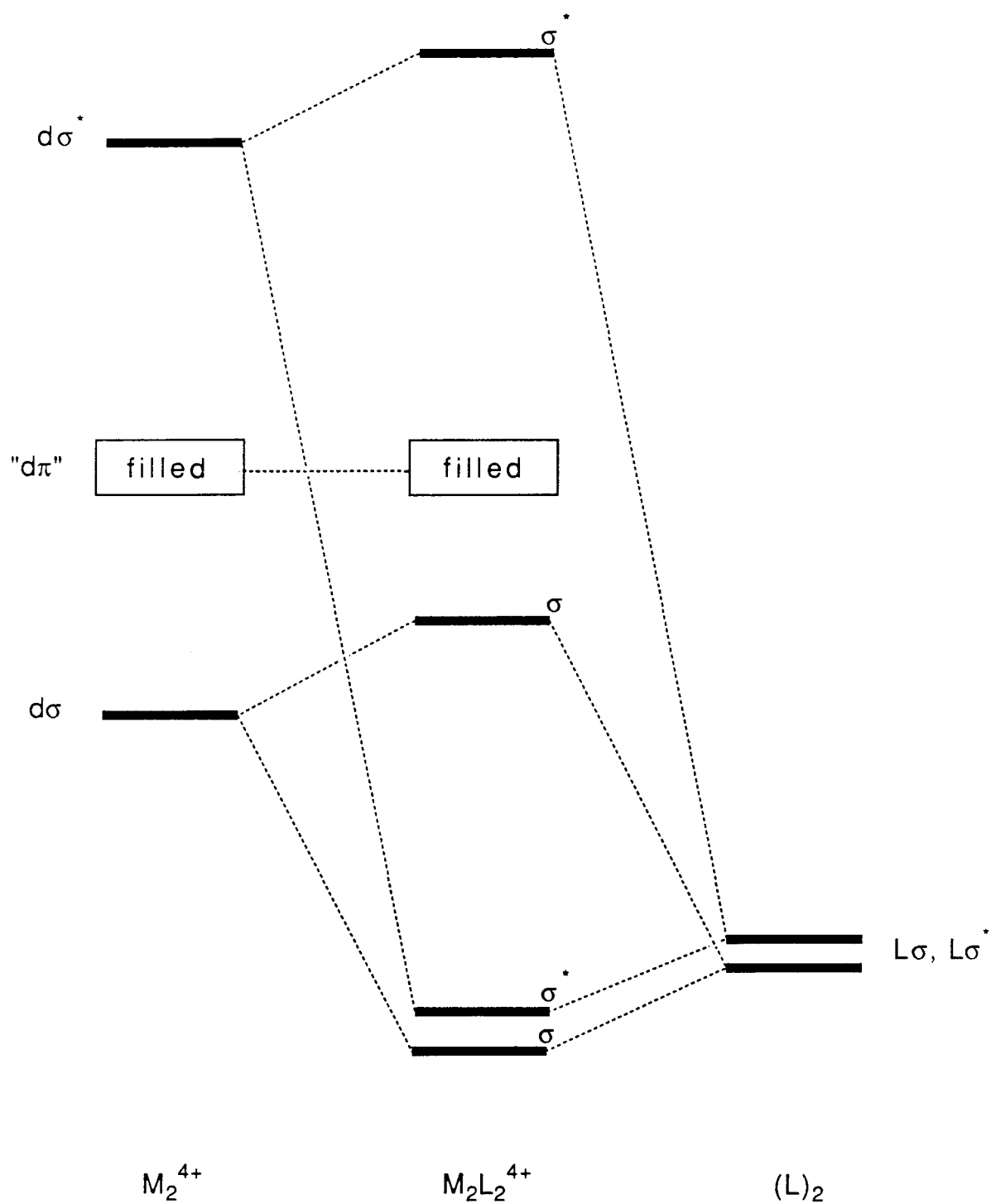
complexes. Earlier work on the dihalide adducts has assigned the intense high-energy feature to the  $d\sigma \rightarrow d\sigma^*$  transition in analogy to  $Mn_2(CO)_{10}$ . The weak lower-energy feature is attributed to the " $d\pi$ "  $\rightarrow d\sigma^*$  transition.

Work on axially substituted  $M_2(CO)_{10}$  complexes had suggested a correlation between the  $d\sigma \rightarrow d\sigma^*$  transition energy and the metal-metal bond strength.<sup>37,38</sup> Additional work on the adducts of the binuclear  $d^8$  complexes suggested a more reasonable explanation for the shift in the  $d\sigma \rightarrow d\sigma^*$  energy involving mixing of the metal-metal  $d\sigma \rightarrow d\sigma^*$  transition with axial ligand-to-metal charge transfer (LMCT).<sup>4</sup> This is indicated by the orbital mixing diagram in Figure 4.6. For the dihalides, a systematic red shift of the  $d\sigma \rightarrow d\sigma^*$  transition is observed through the series Cl, Br, I, a shift with decreasing ionization potential of the axial ligand, an increase in the amount of axial LMCT mixing. Substitution of a hydride for a chloride,  $Ir_2(TMB)_4(H)(Cl)^{2+}$ , results in a dramatic red shift of the  $d\sigma \rightarrow d\sigma^*$  transition,  $6700\text{ cm}^{-1}$ . Within the model, this shift can be understood to arise from a significantly greater charge transfer contribution for  $H^-$  in comparison to  $Cl^-$ . The increased interaction of the hydride ligand with the  $\sigma$ -symmetry combination is evident in comparing the transition energy of  $Ir_2H_2$ . In comparison to  $Ir_2Cl_2$ , a  $5200\text{ cm}^{-1}$  red shift of the  $d\sigma \rightarrow d\sigma^*$  transition is observed. The large red shift of the  $d\sigma \rightarrow d\sigma^*$  transition with introduction of a strong  $\sigma$ -donating ligand is general for the series of complexes listed in Table 4.2 and can be accounted for within the model presented earlier to describe the electronic structure of  $M_2^{4+}$  complexes with two  $\sigma$ -donor axial ligands.

A low-energy feature assignable to the " $d\pi$ "  $\rightarrow d\sigma^*$  transition is not observed for the hydride and alkyl substituted complexes. This transition may now be roughly isoenergetic with the  $d\sigma \rightarrow d\sigma^*$  transition, with the large red shift observed for the latter. Also, the absence of  $\pi$ -symmetry electrons on the ligand may increase the energy of the " $d\pi$ "  $\rightarrow d\sigma^*$  transition because of the loss of  $\pi$ -CT mixing, a simultaneous blue shift of the " $d\pi$ "  $\rightarrow d\sigma^*$  transition and red shift of the  $d\sigma \rightarrow d\sigma^*$  transition. For  $[Ir(\mu-$

**Figure 4.6.** Molecular orbital diagram for interaction of  $M_2^{4+}$  with two  $\sigma$ -donor axial ligands.





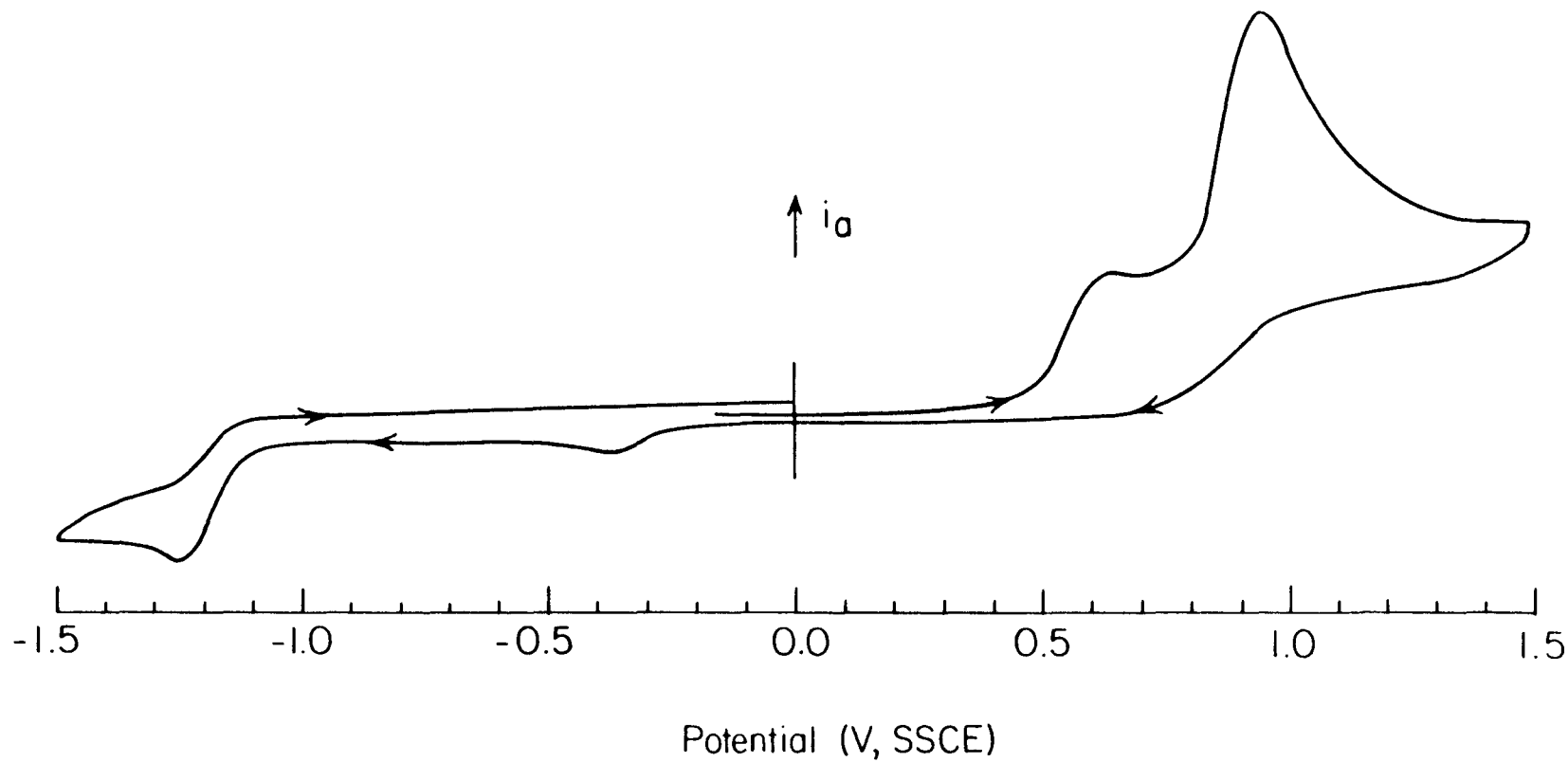
pz)COD(X)]<sub>2</sub> complexes, it has been suggested that the  $d\sigma \rightarrow p\sigma^*$  transition is to lower energy of the " $d\pi$ "  $\rightarrow d\sigma^*$  transition.<sup>39</sup>

## Electrochemistry

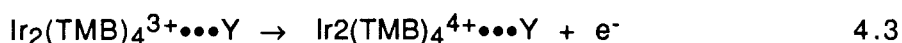
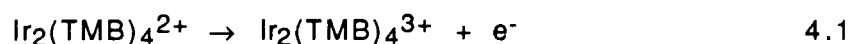
Little is known about the electrochemistry of Ir(I) isocyanide complexes. Much of the work in the literature has focused on Ir(I) phosphine complexes.<sup>40-42</sup> Earlier work on the  $\text{Ir}_2(\text{TMB})_4^{2+}/\text{Ir}_2(\text{TMB})_4^+$  couple found a quasi-reversible reduction at -1.15 V (SCE) in  $\text{CH}_3\text{CN}$ .<sup>3</sup> This reduction was found to be complicated by follow-up chemistry, which was thought to involve interaction between dimer units. No reference to the oxidation of  $\text{Ir}_2(\text{TMB})_4^{2+}$  was made in this work.

The cyclic voltammogram (CV) of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in  $\text{CH}_2\text{Cl}_2$  (0.1 M TBAPF<sub>6</sub>) is shown in Figure 4.7.  $\text{CH}_2\text{Cl}_2$  was chosen over  $\text{CH}_3\text{CN}$  as a solvent because earlier work had shown that the oxidation chemistry of binuclear  $d^8$  complexes was complicated by ECE processes in coordinating solvents.<sup>43</sup> The wave at +0.95 V (SSCE) is associated with the oxidation of the  $\text{B}(\text{C}_6\text{H}_5)_4^-$  counter ion. This observation was confirmed by independent study of  $\text{NaB}(\text{C}_6\text{H}_5)_4$ . The broad wave at +0.63 V (SSCE) is associated with the oxidation of  $\text{Ir}_2(\text{TMB})_4^{2+}$ . This feature appears very irreversible on the electrochemical time scale with no reversible component observed up to scan rates of 100 V/s. The fast-sweep experiments are, however, not conclusive because of experimental difficulties. The value of the  $\text{Ir}_2(\text{TMB})_4^{2+}$  oxidation was found to vary 100 mV between experiments. The weak feature at -0.4 V (SSCE) is associated with the oxidation at +0.63 V. This observation was made by reversing the scan direction and by probing only the first oxidation wave. No features were observed in this region for  $\text{NaB}(\text{C}_6\text{H}_5)_4$ . The wave at -1.25 V (SSCE) is associated with the reduction of  $\text{Ir}_2(\text{TMB})_4^{2+}$ . There was no change in the CV with variation in solvent ( $\text{CH}_2\text{Cl}_2$  and  $(\text{CH}_3)_2\text{CO}$ ).

**Figure 4.7.** Cyclic voltammogram of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  (1 mM) in  $\text{CH}_3\text{CN}$  (0.1 M  $\text{TBAPF}_6$ ).



A possible explanation as to why this couple is so irreversible is suggested by the work of Geiger.<sup>44,45</sup> The oxidation of  $\text{Ir}_2(\text{TMB})_4^{2+}$  will result in a contraction of the metal-metal distance. If the structural rearrangements between the reactants and the products are concerted with the actual electron transfer, an increased energy barrier to electron transfer is expected. If the reorganization energy is of moderate or large magnitude, a quasi-reversible or irreversible electrode reaction will be observed. Another more likely explanation for the irreversible electrochemistry observed for  $\text{Ir}_2(\text{TMB})_4^{2+}$  is a rapid associative chemical step following the initial oxidation.



The intermediate formed (Equation 4.2), being more easily oxidized than  $\text{Ir}_2(\text{TMB})_4^{2+}$ , is rapidly oxidized (Equation 4.3), an ECE process. At present no information toward this end is available.

With knowledge of the ground-state thermodynamics and values for the excited-state energies, an estimate of the excited-state redox potentials can be made.<sup>46,47</sup>

$$E^0(\text{M}^+/\text{}^*\text{M}) = E^0(\text{M}^+/\text{M}) - E_{00}(\text{M}/\text{}^*\text{M}) \quad 4.4$$

$$E^0(\text{}^*\text{M}/\text{M}^-) = E^0(\text{M}/\text{M}^-) + E_{00}(\text{M}/\text{}^*\text{M}) \quad 4.5$$

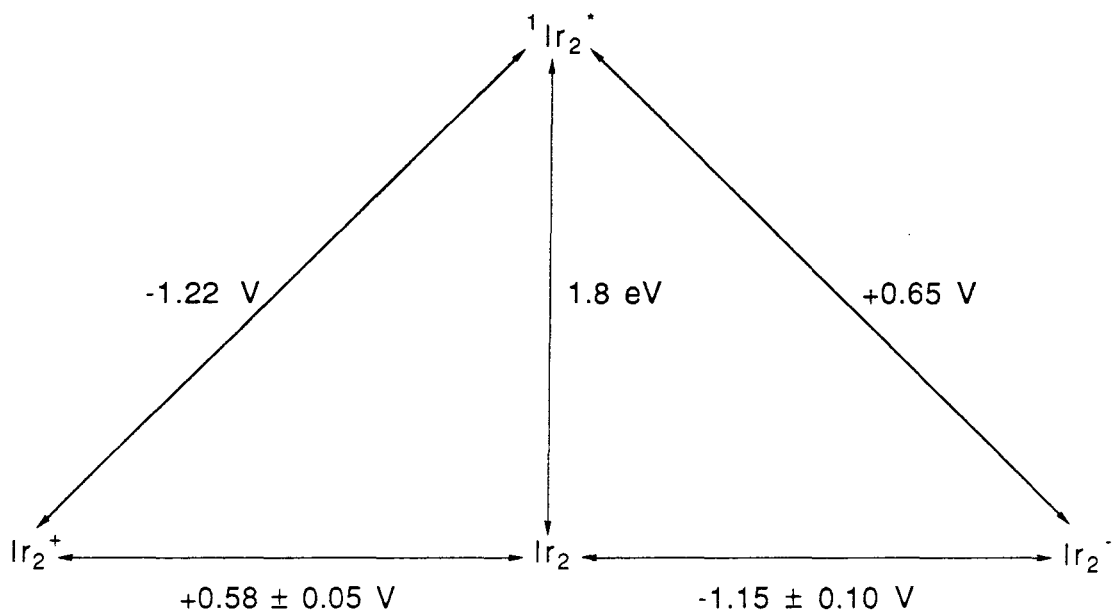
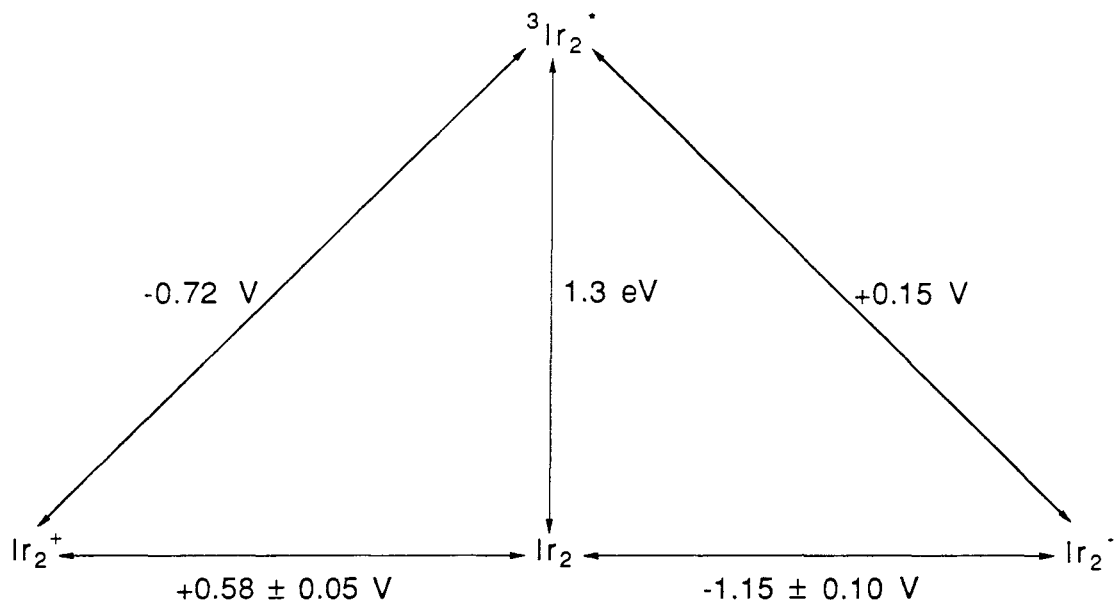
Excited-state potentials can also be estimated from kinetic studies of electron-transfer quenching reactions involving a series of acceptors and/or donors with varying potentials.

Because of the irreversibility of the  $\text{Ir}_2(\text{TMB})_4^{3+}/\text{Ir}_2(\text{TMB})_4^{2+}$  couple, attempts to estimate the excited-state reduction potentials are very problematic. Using

an average value of the  $E_p(\text{oxidation})$ , a modified Latimer diagram can be completed, which describes the excited-state redox potentials (Figure 4.8). The  $^1A_{2u}$  and  $^3A_{2u}$  excited states are, as expected, much more powerful reductants and oxidants than the ground state. One of the major assumptions in this analysis is that the entropy changes associated with the redox processes and excitation are small compared to the enthalpy changes. The excited-state reduction potentials, because of the uncertainty of the  $\text{Ir}_2(\text{TMB})_4^{3+}/\text{Ir}_2(\text{TMB})_4^{2+}$  couple, are only upper bounds of the true values. In comparison to other binuclear  $d^8$  complexes, the excited-state reduction potential of  $\text{Ir}_2(\text{TMB})_4^{2+}$  is on the low end of the energy scale (Table 4.3 and Figure 4.9).

In subsequent sections, the ability of the  $^3(d\sigma^*p\sigma)$  excited state of  $\text{Ir}_2(\text{TMB})_4^{2+}$  to carry out various chemical conversions will be studied. It will be found that while the excited state is a far superior reductant than the ground state, the photochemical reactions follow an atom-transfer pathway. This conclusion will be found to be general for binuclear  $d^8$  complexes.

**Figure 4.8.** Modified Latimer diagrams for  $\text{Ir}_2(\text{TMB})_4^{2+}$  describing the reduction energies for the  $^1(\text{d}\sigma^*\text{p}\sigma)$  and  $^3(\text{d}\sigma^*\text{p}\sigma)$  excited states.



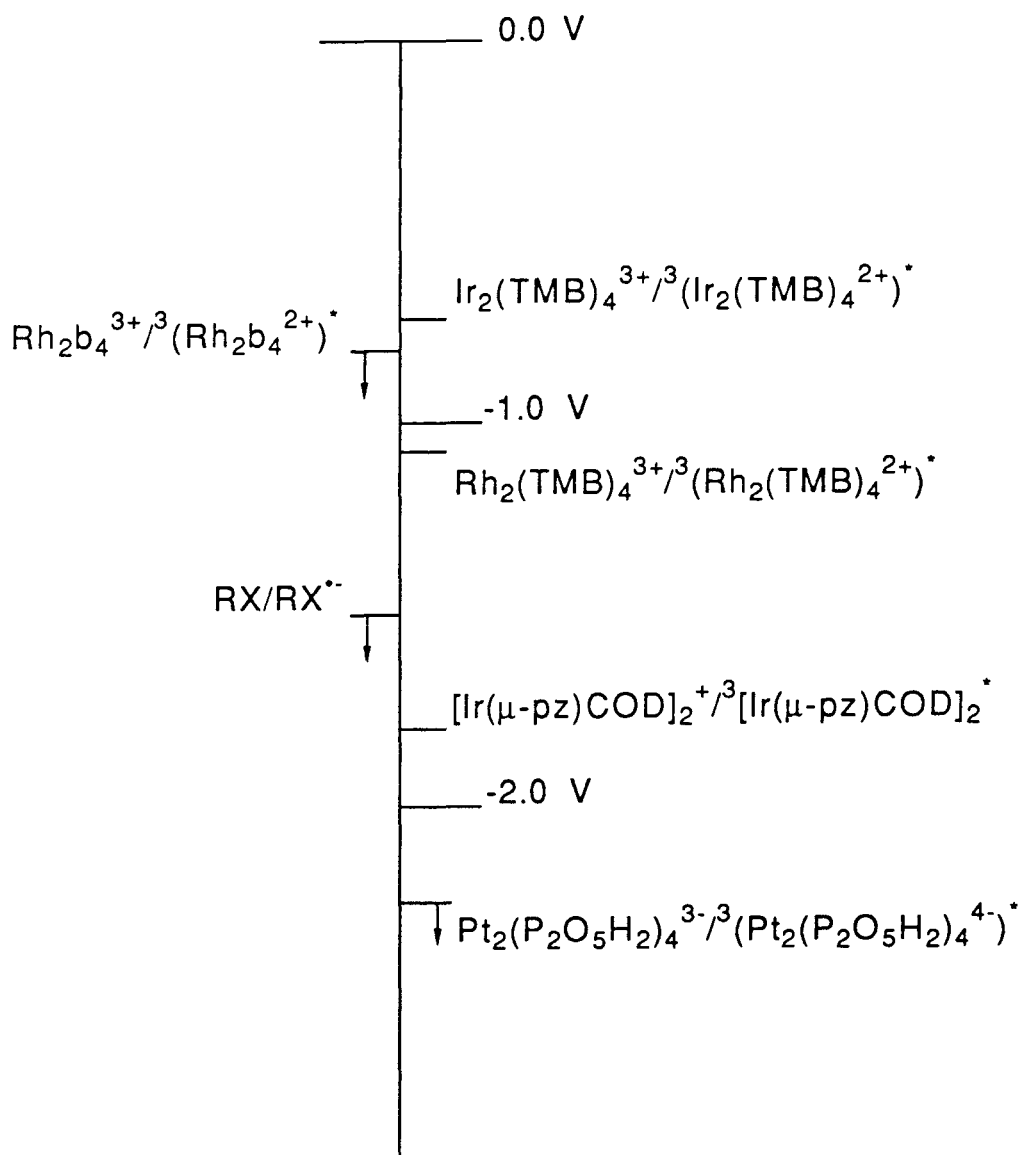


**Table 4.3.** Excited state reduction potentials for binuclear  $d^8$  complexes.

Metal Complex	$E^0(M_2^{+}/^3M_2^*), V (SSCE)$	$E^0(M_2^{+}/^1M_2^*), V (SSCE)$	Ref.
$Ir_2(TMB)_4^{2+}$	-0.72	-1.22	a
$Rh_2b_4^{2+}$	< -0.8	< -1.3	48
$Rh_2(TMB)_4^{2+}$	-1.03	-1.4	a
$[Ir(\mu\text{-pz})COD]_2$	-1.75	-2.2	49
$Pt_2(P_2O_5H_2)_4^{4-}$	-2.24	< -2.7	49

a. This work.

**Figure 4.9.** Reduction potentials for the  $^3(d\sigma^*p\sigma)$  excited state of binuclear  $d^8$  complexes and alkyl halides.



## **Photochemical and Electrocatalytic Reactions of Binuclear $d^8$ Complexes**

## Introduction

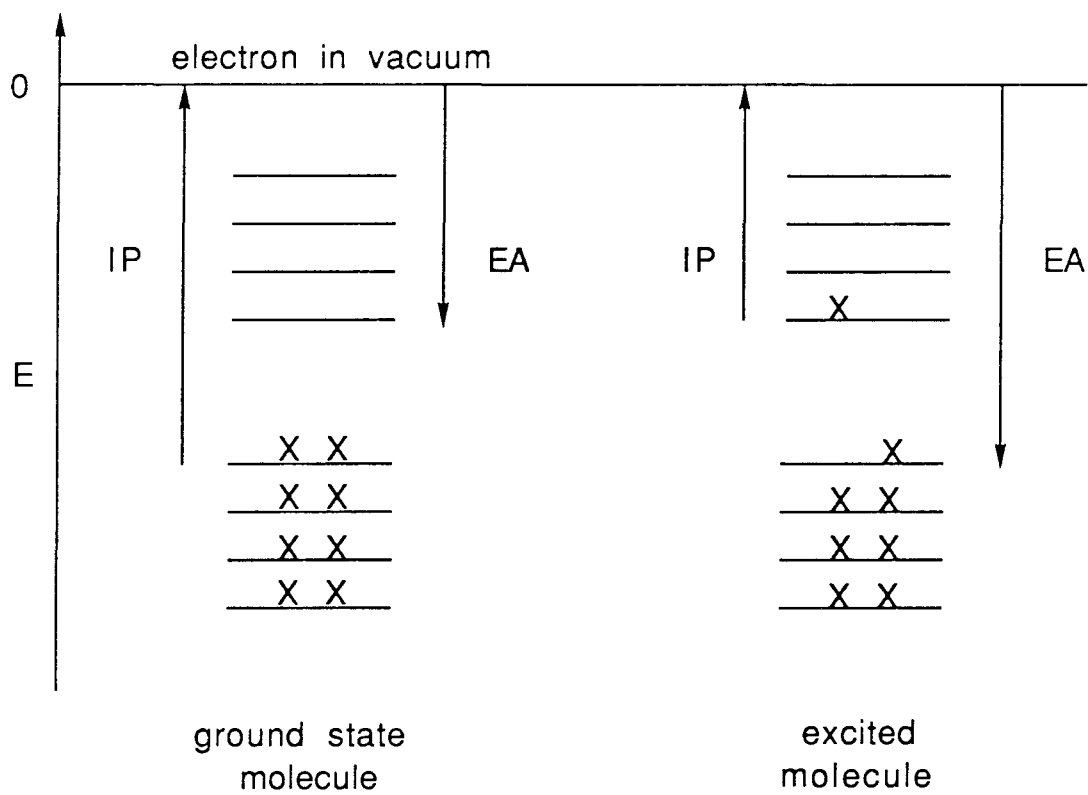
An electronic excited state of a metal complex is both a stronger reductant and a stronger oxidant than the ground state.<sup>46</sup> Excitation of an electron from a low-energy to a high-energy orbital reduces the ionization potential and increases the electron affinity of a molecule (Figure 4.10). Therefore, complexes with relatively long-lived excited states can participate in intermolecular electron-transfer reactions that are uphill for the corresponding ground-state species. Such excited-state electron-transfer reactions often play key roles in multistep schemes for the conversion of light to chemical energy.<sup>50</sup>

While electron-transfer processes are common in inorganic photochemistry, excited-state atom transfer is limited to a small class of inorganic complexes. For  $\text{UO}_2^{2+}$ , the diradical excited state ( $\bullet\text{U}-\text{O}\bullet$ ) is active in alcohol oxidation.<sup>51</sup> The primary photoprocess is hydrogen-atom abstraction by the oxygen-centered radical. Photoaddition to a metal center via atom transfer has been observed for binuclear metal complexes such as  $\text{Re}_2(\text{CO})_{10}$ .<sup>52-55</sup> The primary photoprocess is metal-metal bond homolysis. The photogenerated metal radical undergoes thermal atom-abstraction reactions. Until recently, atom transfer to a metal-localized excited state had not been observed.

Atom transfer to a metal complex is facilitated if localized electron or hole generation occurs at one or more open coordination sites. Binuclear  $d^8$  complexes have been found to undergo photochemical atom transfer to one of the metal centers.<sup>2,30,49,56</sup> These complexes possess open coordination sites in addition to an electronic structure that localizes the electron (hole) necessary for atom transfer to the metal center.

The lowest energy optical transition for the binuclear systems is  $d\sigma^* \rightarrow p\sigma$ . This excitation results in formation of a formal metal-metal single bond in the excited state. The transition is metal-localized and can be viewed as movement of an electron from an

**Figure 4.10.** Schematic orbital diagram showing the decrease in the ionization potential and increase in the electron affinity upon excitation of a molecule.



orbital localized on the exterior of the  $M_2$  unit (the  $d\sigma^*$  orbital) to an orbital localized in the interior of the dimer cage (the  $p\sigma$  orbital). The excitation results in hole formation localized on a metal center at an open coordination site (Figure 4.11). The lifetime of the  $^3(d\sigma^*p\sigma)$  excited state, normally in between 100 ns and 10  $\mu$ s, makes  $M_2$  systems attractive for bimolecular photoprocesses.

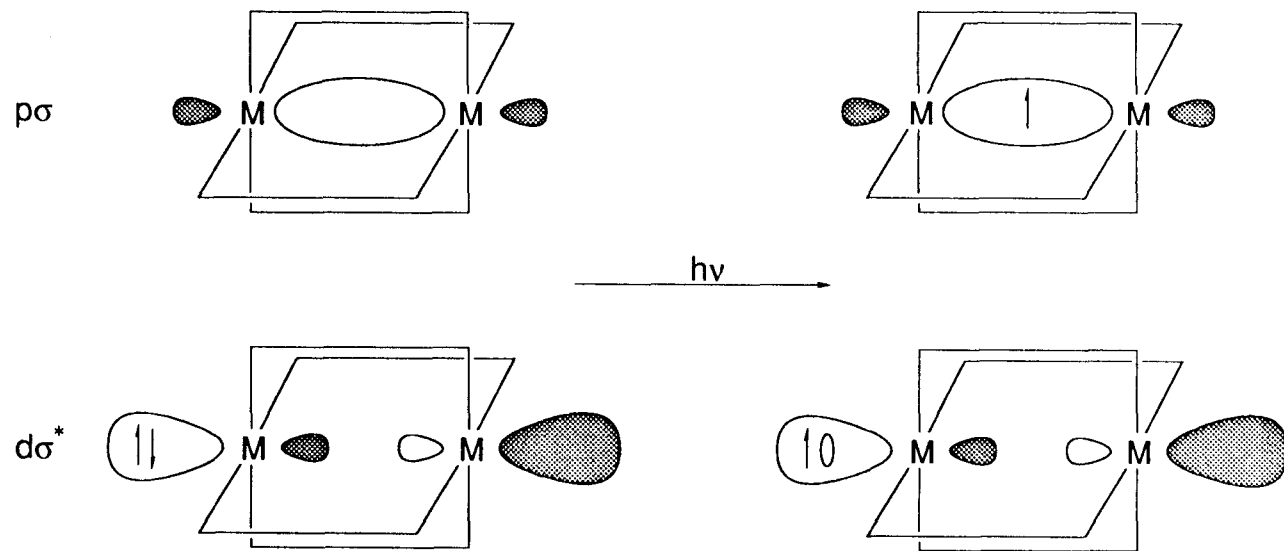
The initial interest in these systems was stimulated by observations of their photochemical electron-transfer reactivity.<sup>48,49</sup> From spectroscopic and electrochemical studies, the  $^3(d\sigma^*p\sigma)$  excited state is predicted to be a powerful reductant, with  $E^0(M_2^+/^3M_2^*)$  estimated to range from -0.8 to -2.0 V vs SSCE in  $CH_3CN$ . That this state is a powerful reductant has been confirmed by investigation of the electron-transfer quenching of  $^3M_2^*$  by a series of pyridinium acceptors with varying reduction potentials.<sup>57</sup> For several binuclear complexes, the excited-state reduction potential cannot be calculated accurately because of the irreversibility of the ground-state electrochemistry; but it can be estimated from bimolecular electron-transfer quenching experiments.

For systems that are powerful excited-state reductants, photoreduction of alkyl halides is observed.<sup>49,58</sup> This reaction was initially interpreted to be an outer-sphere electron transfer to form an alkyl halide radical anion, which rapidly decomposes to yield  $R\cdot$  and  $X^-$ . Subsequent thermal reactions give the observed products, an  $S_{RN}1$  mechanism (Figure 4.12). While such a mechanism,  $S_{RN}1$ , appears plausible for a metal complex with  $E^0(M_2^+/^3M_2^*) < -1.5$  V (SSCE), it seems unlikely for complexes with  $E^0(M_2^+/^3M_2^*) > -1.0$  V (SSCE). Reduction potentials for alkyl halides of interest are generally more negative than -1.5 V (SSCE).<sup>59</sup> Alkyl-halide photoreduction is observed for binuclear  $d^8$  complexes whose excited-state reduction potentials are more positive than -1.0 V (SSCE) in  $CH_3CN$ .

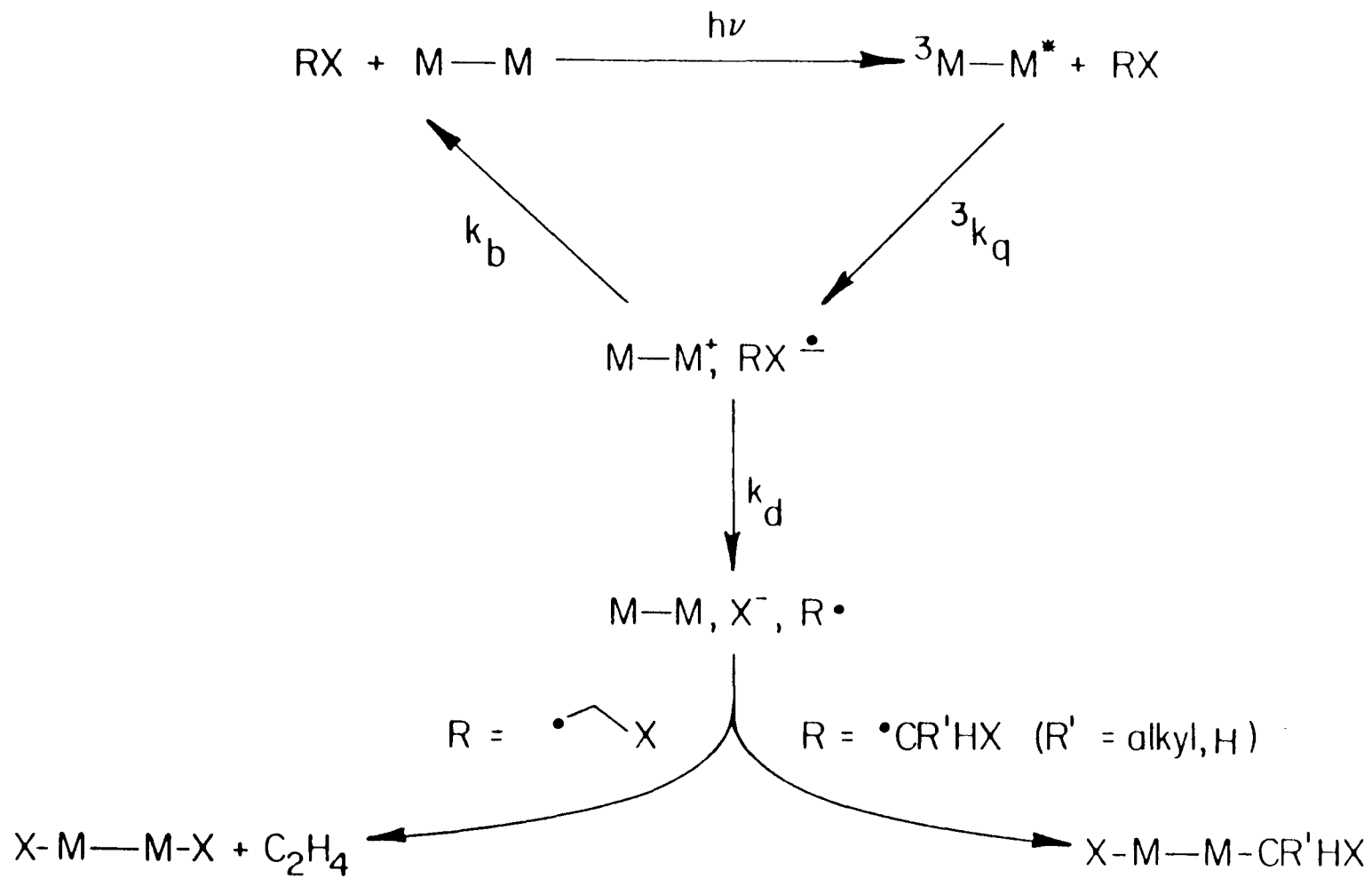
An alternative pathway to outer-sphere electron transfer, which yields similar photoredox products with alkyl halides, is excited-state atom transfer (Figure 4.13).



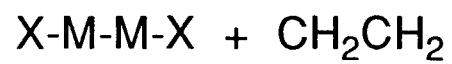
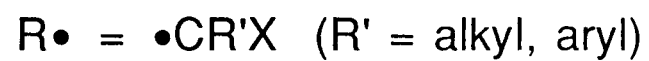
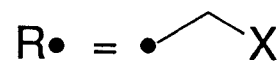
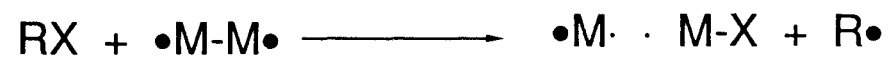
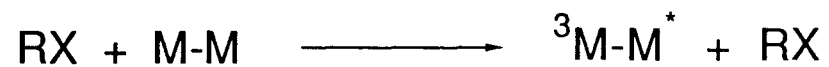
**Figure 4.11.** Pictorial representation of the  $M_2$ -localized hole in a  $^3(d\sigma^*p\sigma)$  state.



**Figure 4.12.**  $S_{RN}1$  mechanistic scheme for halocarbon photooxidative addition to binuclear  $d^8$  complexes.



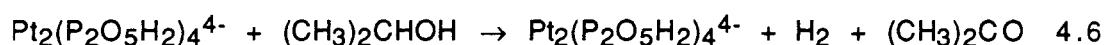
**Figure 4.13.** Atom-transfer mechanism for halocarbon photooxidative addition.



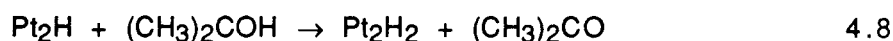
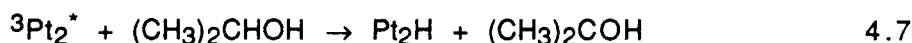
Data obtained for  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$  indicate that alkyl and aryl halides react with the  $^3(\text{d}\sigma^*\text{p}\sigma)$  excited state via halogen-atom transfer.<sup>2</sup>

Although the primary photoprocess for alkyl-halide photoreduction may not be atom transfer in all cases,  $^3(\text{d}\sigma^*\text{p}\sigma)$  excited-state hydrogen-atom transfer has been established as the mechanism of the reactions between several binuclear  $\text{d}^8$  complexes and a number of organic and organometallic substrates.<sup>2,29,49,56</sup> Initial work in this area focused on  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$ , for which the catalytic conversion of isopropanol to acetone (Equation 4.6) had been first observed.<sup>60</sup>

$h\nu$



From detailed studies of this system, it was concluded that the primary photoprocess is abstraction of the  $\alpha$ -hydrogen by the  $^3\text{Pt}_2^*$  to form a monohydride species (directly observed by transient absorption spectroscopy for a number of substrates) and the organic radical (Equation 4.7), with the final photoproduct being  $\text{Pt}_2\text{H}_2$  and acetone (Equation 4.8). The  $\text{Pt}_2\text{H}_2$  complex has been characterized by NMR, UV-Vis, and IR (but has not been successfully isolated).<sup>61</sup>



The relatively long lifetime of the  $^3(\text{d}\sigma^*\text{p}\sigma)$  excited state of  $\text{Ir}_2(\text{TMB})_4^{2+}$  in fluid solution suggests that it should be able to participate in bimolecular photochemical reactions. In light of the favorable absorption and emission properties of this complex, it is of interest to examine the possibility of the conversion of visible light to net chemical energy through a photochemical process, either electron transfer or atom

transfer. The excited-state electron-transfer reactivity of  $\text{Ir}_2(\text{TMB})_4^{2+}$  with pyridinium acceptors is reported and analyzed in the context of classical Marcus theory for outer-sphere electron transfer. The photochemical reactions of  $\text{Ir}_2(\text{TMB})_4^{2+}$  with alkyl halides and hydrogen atom donors are reported.



## Experimental

### Synthesis

#### *Pyridinium Hexafluorophosphates*

All solvents were from freshly opened bottles with no further purification. All chemicals were of reagent grade or comparable quality and used as received. The  $^1\text{H}$  NMR spectra were obtained on a 90 MHz EM-390 spectrometer. Elemental analyses were obtained at the Cal Tech Analytical Laboratory.

Several of the pyridinium hexafluorophosphates were prepared and recrystallized by Dr. Janet L. Marshall<sup>39</sup> and used as received. Methyl violegen di(hexafluorophosphate) was received from Dr. Miriam A. Heinrichs.

#### *4-CN-N-Benzylpyridinium*

A slight excess of benzyl bromide was added to an acetone solution of 4-CN-pyridine. After refluxing for 10 hrs, the solution was cooled to room temperature and the precipitated pyridinium halide was filtered off. The yellow powder was dissolved in a minimum of distilled water and metathesized with  $\text{KPF}_6$ . The precipitated pyridinium hexafluorophosphate was filtered off and recrystallized from acetone. Calculated for  $\text{C}_{13}\text{H}_{11}\text{NPF}_6$ : C, 45.9; H, 3.3; N, 8.2. Found: C, 45.8; H, 3.4; N, 8.4.  $^1\text{H}$  NMR:  $\delta$  ( $\text{CD}_3\text{CN}$ , 20 °C); 5.8 (singlet, 2H,  $\text{CH}_2$ ), 7.5 (singlet, 5H,  $\text{C}_6\text{H}_5$ ), 8.4 (multiplet, 2H, 3,5-H), 8.9 (doublet, 2H, 2,6-H).

#### *3-CN-N-Methylpyridinium*

A tenfold excess of methyl iodide was added to a stirred THF solution of 3-cyanopyridine. The solution was stirred for 48 hrs. The precipitated pyridinium iodide was filtered off and washed with THF. The yellow powder was dissolved in distilled water and metathesized with  $\text{NH}_4\text{PF}_6$ . The white precipitate was recrystallized from

H<sub>2</sub>O/acetone/ethanol (2:1:1). Calculated for C<sub>7</sub>H<sub>7</sub>N<sub>2</sub>PF<sub>6</sub>: C, 31.8; H, 2.7; N, 10.6.

Found: C, 32.0; H, 2.6; N, 10.6.

### *3-Carbomethoxy-N-Benzylpyridinium*

A fivefold excess of benzyl chloride was combined with methyl nicotinate in a 1:1 ethanol/acetone solution. The solution was refluxed for 10 hrs. No solid material precipitated upon cooling to room temperature. The solvent was removed under vacuum, yielding a viscous oil. The oil was dissolved in a minimum of distilled water and metathesized with KPF<sub>6</sub>. The white precipitate was recrystallized from methanol. Calculated for C<sub>14</sub>H<sub>14</sub>NO<sub>2</sub>PF<sub>6</sub>: C, 45.1; H, 3.8; N, 3.8. Found: C, 45.2; H, 3.8; N, 3.8. <sup>1</sup>H NMR: δ (CD<sub>3</sub>CN, 20 °C); 3.9 (singlet, 3H, CH<sub>3</sub>), 5.8 (singlet, 2H, CH<sub>2</sub>), 7.5 (singlet, 5H, C<sub>6</sub>H<sub>5</sub>), 8.1 (triplet, 1H, 5-H), 8.9 (multiplet, 2H, 4,6-H), 9.3 (singlet, 1H, 2-H).

### *3-Carbomethoxy-N-Methylpyridinium*

A tenfold excess of methyl iodide was added to a THF solution of methyl nicotinate. After stirring the solution overnight, the precipitated pyridinium iodide was filtered off and washed with THF. The yellow powder was dissolved in a minimum of distilled water and metathesized with NH<sub>4</sub>PF<sub>6</sub>. The white precipitate was recrystallized from H<sub>2</sub>O/acetone/ethanol (2:1:1). Calculated for C<sub>8</sub>H<sub>10</sub>NO<sub>2</sub>PF<sub>6</sub>: C, 32.2; H, 3.4; N, 4.7. Found: C, 32.2; H, 3.1; N, 4.8. <sup>1</sup>H NMR: δ (CD<sub>3</sub>CN, 20 °C); 4.0 (singlet, 3H, OCH<sub>3</sub>), 4.4 (singlet, 3H, NCH<sub>3</sub>), 8.1 (triplet, 1H, 5-H), 8.8 (multiplet, 2H, 4,6-H), 9.2 (singlet, 1H, 2-H).

### *Additional Compounds*

Standard synthetic procedures were used to prepare

[Ir<sub>2</sub>(TMB)<sub>4</sub>](B(C<sub>6</sub>H<sub>5</sub>)<sub>4</sub>)<sub>2</sub>,<sup>62</sup> [Rh<sub>2</sub>b<sub>4</sub>](B(C<sub>6</sub>H<sub>5</sub>)<sub>4</sub>)<sub>2</sub>,<sup>63</sup> and [Rh<sub>2</sub>(TMB)<sub>4</sub>](PF<sub>6</sub>)<sub>2</sub>.<sup>64</sup>

## Physical Measurements

### Materials

All solvents were purified, if necessary, degassed with a minimum of five freeze-pump-thaw cycles on a high-vacuum line, and bulb-to-bulb distilled into glass, round-bottomed storage flasks equipped with Teflon vacuum valves. Acetonitrile was used as received and stored over activated alumina. 1,2-Dichloroethane was distilled from  $\text{CaH}_2$  and stored over activated molecular sieves. Dichloromethane was distilled from  $\text{CaH}_2$  and stored under argon. The halocarbon solvents' storage flasks were wrapped in foil to prevent exposure to room light. 1,4-Cyclohexadiene was distilled from  $\text{NaBH}_4$  under argon, freeze-pump-thaw degassed, and stored under vacuum, protected from light. Cyclohexene was distilled from either  $\text{CaH}_2$  or Na, freeze-pump-thaw degassed, and stored under vacuum, protected from the light. 9,10-Dihydroanthracene was recrystallized three times from absolute ethanol. Cyclohexene- $d_{10}$  was used as received. All other hydrocarbon quenchers were purified by standard methods,<sup>22,23</sup> freeze-pump-thaw degassed, and stored, protected from light, under vacuum or argon.

Tetra-*n*-butyl ammonium hexafluorophosphate was prepared from tetra-*n*-butyl ammonium iodide and ammonium hexafluorophosphate by addition of a hot, saturated, acetone solution of  $\text{NH}_4\text{PF}_6$  to a hot, saturated, acetone solution of  $\text{Bu}_4\text{NI}$ . To this hot solution was added a large excess of distilled water. The solution was cooled and filtered. The precipitated  $\text{Bu}_4\text{NPF}_6$  ( $\text{TBAPF}_6$ ) was recrystallized three times from 95% ethanol.

### Electrochemistry

Cyclic voltammograms (CVs) and constant potential bulk electrolysis were performed using a Princeton Applied Research (PAR) model 173 potentiostat/galvanostat, a model 175 universal programmer, and a model 179 digital

coulometer. The CVs were plotted on a Houston Instruments Omnigraphic 2000 x,y-recorder. Volatile organics obtained from bulk electrolysis were analyzed on a Hewlett-Packard 5890A gas chromatograph. The UV-Vis spectra were measured on a Shimadzu UV-260 spectrophotometer.

### *Pyridinium Hexafluorophosphates*

The cell geometry used to measure the CVs is described elsewhere.<sup>24</sup> CVs were measured at a BAS Pt button electrode, with a Pt wire auxiliary electrode, and a sodium saturated calomel electrode (SSCE) as the reference. CVs of the pyridinium hexafluorophosphates were measured for acetonitrile solutions under argon that contained 0.1 M TBAPF<sub>6</sub> as the supporting electrolyte and 10<sup>-3</sup> M compound. Acetonitrile was obtained from a freshly opened bottle and was used as received.

### *Rh<sub>2</sub>(TMB)<sub>4</sub><sup>2+</sup>*

The cell geometry used to measure the CVs for [Rh<sub>2</sub>(TMB)<sub>4</sub>](PF<sub>6</sub>)<sub>2</sub> was similar to that used to measure the CVs of the pyridinium hexafluorophosphates. CVs were measured at either a BAS Pt button or a BAS glassy carbon electrode, with a Pt wire auxiliary electrode, and an SSCE as the reference. CVs were measured of dichloromethane solutions under argon that contained 0.1 M TBAPF<sub>6</sub> as supporting electrolyte and 10<sup>-3</sup> M compound.

The Pt button electrode was prepared following the procedure described by Anson, *et al.*<sup>65</sup> The glassy carbon electrode was polished with 0.3 μm α-alumina slurry, sonicated in purified water for 15 mins, and rinsed with dichloromethane.

Constant potential bulk electrolysis experiments were done in an H-tube type cell at either a Pt basket or a pyrolytic graphite working electrode. The H-tube cell was modified to allow CVs to be measured of the solution in the compartment with the working electrode. CVs were measured at either a BAS Pt button or a BAS glassy carbon

electrode. The auxiliary electrode was a Pt basket and was separated from the working electrode by a medium-porosity, sintered glass frit. Potentials were measured with respect to an SSCE. Measurements were made of dichloromethane solutions under argon, which contained 0.1 M TBAPF<sub>6</sub> as the supporting electrolyte and 10<sup>-3</sup> M compound. Substrates were added to the electrolysis cell via syringe.

A twentyfold excess of 1,4-cyclohexadiene (or cyclohexene) was added to a dichloromethane solution of [Rh<sub>2</sub>(TMB)<sub>4</sub>](PF<sub>6</sub>)<sub>2</sub> and TBAPF<sub>6</sub>. The initial solution was clear, red. Exhaustive electrolysis yielded a clear, light-red solution. This solution was transferred to a round-bottom flask and the volatiles were vacuum-transferred to a second round-bottom flask. Benzene was identified as the electrolysis product by gas chromatography. Production of protons was confirmed from the CV of the final solution. This CV compared well to a CV of a dichloromethane solution of HBF<sub>4</sub>. Extraction of the dichloromethane electrolysis solution with distilled water resulted in an increase in pH of the water phase. UV-Vis spectra of the electrolysis solution during the course of the reaction showed the slow appearance of Rh<sub>2</sub>(TMB)<sub>4</sub>Cl<sub>2</sub><sup>2+</sup> with loss of Rh<sub>2</sub>(TMB)<sub>4</sub><sup>2+</sup>. Rh<sub>2</sub>(TMB)<sub>4</sub>Cl<sub>2</sub><sup>2+</sup> was isolated as a yellow powder from the final electrolysis solution.

### Electronic Emission Spectroscopy

Electronic emission spectra were measured using an emission spectrophotometer constructed at Cal Tech, which has been described previously.<sup>66</sup> Spectra recorded at ambient temperature were obtained for solutions that were prepared on a high-vacuum line in a cell consisting of a 10 ml Pyrex bulb, a 1 cm pathlength quartz cell, and a Teflon vacuum valve. Solvent was bulb-to-bulb distilled into the cell from the appropriate solvent storage flask.

## Stern-Volmer Quenching

For the experiments with methyl viologen di(hexafluorophosphate) and the pyridinium hexafluorophosphate quenchers, acetonitrile solutions of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  and  $\text{TBAPF}_6$  ( $\mu = 0.1 \text{ M}$ ) were prepared on a high-vacuum line in a two-compartment spectrophotometric cell consisting of a 1 cm pathlength square cuvette and a 10 ml bulb. Acetonitrile was vacuum-distilled from its storage flask to a calibrated volumetric cylinder, then to the evacuated cell containing the iridium complex and the electrolyte. Addition of solid quencher to one compartment of the cell was followed by re-evacuation of the cell to maintain the oxygen-free conditions of the experiment. Rate constants for the quenching of the triplet excited state of  $\text{Ir}_2(\text{TMB})_4^{2+}$  were determined by measuring its lifetime as a function of quencher concentration. The emission lifetime measurements were conducted with a Nd:YAG pulsed laser system that has been described previously, using 532 nm excitation.<sup>67</sup> All emission intensity decays exhibited first-order kinetics over at least three half-lives. Nonlinear-least-squares fits of  $\text{RTln}(k_q')$  ( $k_q' = \text{activation-controlled rate constant}$ ) versus  $E^0(\text{A}^{+/0})$  using Equations 4.20 and 4.23 were accomplished using the program MARQUARDT.

Rate constants for the quenching of the triplet excited of  $\text{Ir}_2(\text{TMB})_4^{2+}$  with hydrogen-atom donors were determined by the same methods as described for the study using the pyridinium hexafluorophosphates. Acetonitrile solutions were prepared in a similar fashion except with no added electrolyte. The hydrocarbon quenchers were either added via syringe in an inert atmosphere and the solution freeze-pump-thaw degassed after each addition, or vacuum-transferred from an appropriate storage flask. At least two independently prepared lots of hydrocarbon quencher were used for each kinetic analysis. For cyclohexene, two different purification procedures were also used. Unlike the electron-transfer quenching, irreversible photochemical reactions were observed with some hydrocarbon quenchers. For these substrates, each data point

required a separate sample to be prepared. Low laser powers along with a limited number of laser pulses (20 to 50) were used to minimize the amount of photoreaction. All emission-intensity decays exhibited first-order kinetics over at least three half-lives.

Absorption spectra of the solutions were measured with a Hewlett-Packard 8450A spectrophotometer or a Shimadzu UV-260 spectrophotometer before and after the quenching experiments to insure that the quenching was reversible and that no reaction, either thermal or photochemical, was occurring between  $\text{Ir}_2(\text{TMB})_4^{2+}$  and the quenchers during the course of the experiments.

### Steady-State Photolysis

Steady-state photolysis experiments were conducted on solutions prepared on a high-vacuum line in a two-compartment cell. Solvent was bulb-to-bulb distilled into the cell containing the metal complex from the appropriate storage flask followed by the addition of the reactant. Liquid reactants were bulb-to-bulb distilled into the cell, and solid reactants were added to one compartment of the cell, followed by re-evacuation of the cell to maintain oxygen-free conditions.

$[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in neat 1,2-dichloroethane was irradiated with  $\lambda_{\text{ex}} > 604$  nm and showed formation of  $\text{Ir}_2(\text{TMB})_4\text{Cl}_2^{2+}$ . Solutions kept in the dark showed no formation of the oxidative addition product. Solutions of  $[\text{Rh}_2\text{b}_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in neat 1,2-dichloroethane were irradiated with  $\lambda_{\text{ex}} > 500$  nm and showed rapid formation of  $\text{Rh}_2\text{b}_4\text{Cl}_2^{2+}$ . Solutions kept in the dark showed no formation of the oxidative addition products. The  $\text{d}^7\text{-d}^7$  complexes were identified by their absorption spectra; the organic products were not identified.

Acetonitrile solutions of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  and the hydrogen-atom donors were photolyzed with  $\lambda_{\text{ex}} > 550$  nm. The oxidative addition product was identified by its absorption spectrum and NMR spectrum. The organic products from the

photoreaction of 9,10-dihydroanthracene and 1,4-cyclohexadiene were identified by their NMR spectra.

A 1000 W high-pressure Hg/Xe arc lamp equipped with Corning cut-off filters was used for the irradiations. Spectrophotometric monitoring of both photolysis and thermal-blank experiments was done with either a Hewlett-Packard 8450A spectrophotometer or a Shimadzu UV-260 spectrophotometer.  $^1\text{H}$  NMR spectra were obtained on a 400 MHz JNM-GX400 FT NMR spectrometer.



## Results and Discussion

### Homogeneous Electron Transfer

Earlier work has shown that the  $^3(d\sigma^*p\sigma)$  excited state of  $\text{Ir}_2(\text{TMB})_4^{2+}$  can be reductively quenched with *N,N,N',N'*-tetramethylphenylenediamine (TMPD), yielding  $\text{TMPD}^{+\bullet}$  and presumably  $\text{Ir}_2(\text{TMB})_4^{+}$ .<sup>3</sup> Kinetic data were not reported. Similar reductive quenching has been reported for  $\text{Rh}_2(\text{TMB})_4^{2+}$ ,  $\text{Rh}_2\text{b}_4^{2+}$ , and  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$  with spectroscopic characterization of the reduced products.<sup>48</sup> No discussion of the oxidative quenching of  $\text{Ir}_2(\text{TMB})_4^{2+}$  was presented in the earlier work.

It has been reported that  $\text{Ir}_2(\text{TMB})_4^{2+}$  undergoes a thermal reaction with *N,N'*-dimethyl-4,4'-pyridinium ( $\text{MV}^{2+}$ ) or with chloranil.<sup>3</sup> No characterization of the final products was presented. Tetracyanoethylene and chloranil react thermally with  $\text{Ir}_2(\text{TMB})_4^{2+}$ . The final products appear to be the one-electron reduced acceptors and the two-electron oxidized metal complex. Contrary to the earlier result, there is no thermal reaction between  $\text{MV}(\text{PF}_6)_2$  and  $\text{Ir}_2(\text{TMB})_4^{2+}$ .

While no thermal ground-state reaction with  $\text{MV}^{2+}$  is observed for  $\text{Ir}_2(\text{TMB})_4^{2+}$ ,  $\text{MV}^{2+}$  efficiently quenches the phosphorescence of the metal complex. Electron transfer is the expected quenching process with a net driving force for reaction of the  $^3(d\sigma^*p\sigma)$  excited state of  $\text{Ir}_2(\text{TMB})_4^{2+}$  with  $\text{MV}^{2+}$  of 0.27 V ( $E_{1/2}(\text{MV}^{2+}/\text{MV}^+) = -0.45$  V (SSCE)).<sup>68</sup> The triplet energy of  $\text{MV}^{2+}$  (3.10 eV)<sup>69</sup> is too high for energy-transfer quenching to be competitive with electron-transfer quenching of  $^3(\text{Ir}_2^{2+})^*$ , and energy transfer can be ruled out. Additional support for excited-state electron transfer is obtained from the measurement of the rate of the quenching reaction (Equation 4.9).



The reaction obeys Stern-Volmer quenching kinetics;<sup>70</sup> the quenching rate constant is determined by measuring the change in the  $^3(\text{d}\sigma^*\text{p}\sigma)$  excited state lifetime as a function of the concentration of  $\text{MV}^{2+}$ .

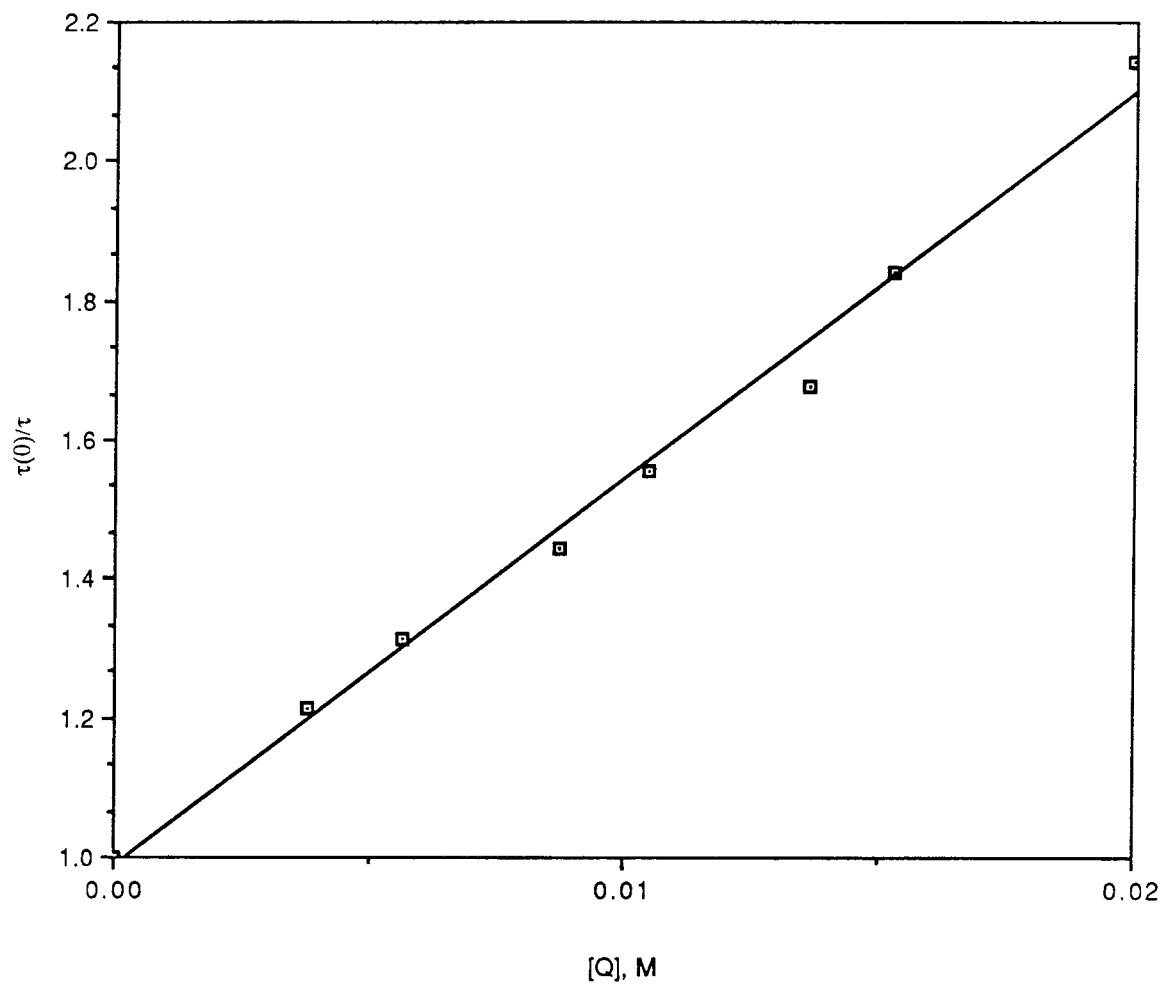
$$\frac{\tau_0}{\tau} = 1 + k_q\tau_0[\text{Q}] \quad 4.10$$

A plot of  $\tau_0/\tau$  versus  $[\text{Q}]$ , where  $\tau_0$  is the unquenched  $^3(\text{d}\sigma^*\text{p}\sigma)$  lifetime and  $\tau$  is the quenched  $^3(\text{d}\sigma^*\text{p}\sigma)$  lifetime, yields a line of slope  $k_q\tau_0$  (Figure 4.14). The measured rate constant for the quenching by  $\text{MV}^{2+}$  ( $2.7 \pm 0.3 \times 10^8 \text{ M}^{-1}\text{s}^{-1}$ ) is approximately two orders of magnitude below the calculated diffusion-limited rate constant in  $\text{CH}_3\text{CN}$  ( $k_d \sim 2 \times 10^{10} \text{ M}^{-1}\text{s}^{-1}$ ).<sup>71</sup> Attempts to characterize the redox products from the photoreaction using conventional microsecond flash photolysis were unsuccessful, presumably because back electron transfer occurs too quickly.

A small change in the optical absorption spectrum of a  $\text{CH}_3\text{CN}$  solution of  $\text{Ir}_2(\text{TMB})_4^{2+}$  and  $\text{MV}^{2+}$  was observed over an extended period of irradiation. This change may be attributed to the formation of trace amounts of  $\text{MV}^{+\bullet}$ .<sup>72</sup> No significant build-up of products is observed since the rapid back reaction of the quenching products  $\text{Ir}_2^{3+}$  and  $\text{MV}^{+\bullet}$  gives ground-state  $\text{Ir}_2^{2+}$  and  $\text{MV}^{2+}$ . Consequently, this photoinduced electron transfer does not lead to net chemistry.

The measured rate constant for electron-transfer quenching of  $^3(\text{Ir}(\text{TMB})_4^{2+})^*$  by  $\text{MV}^{2+}$  is compared to the measured rate constants for several related reactions in Table 4.4. The variation in the observed rates may be understood as a result of the increased driving force (more negative  $\Delta G^0$ ) for electron transfer through the series of complexes. Such an argument appears reasonable for the comparison to  $\text{Rh}_2(\text{TMB})_4^{2+}$  and  $[\text{Ir}(\mu\text{-pz})\text{COD}]_2$ . These systems should have comparable electronic and nuclear factors (for the same intermolecular separation) that determine the reaction rate; redox

**Figure 4.14.** Stern-Volmer plot for the oxidative quenching of the  $^3(d\sigma^*p\sigma)$  excited state of  $\text{Ir}_2(\text{TMB})_4^{2+}$  by methyl viologen di(hexafluorophosphate) (0.1 M  $\text{TBAPF}_6$ ,  $\text{CH}_3\text{CN}$ ).



**Table 4.4.** Measured rate constant for the reaction  $^3M_2^+ + MV^{2+} \rightarrow M_2^+ + MV^+$  or  $M^+ + MV^{2+} \rightarrow M^+ + MV^+$ .

Metal Complex	$k_q, M^{-1}s^{-1}$	$E^0(M^+/M^*), V$ (SSCE)	Ref.
$Ir_2(TMB)_4^{2+}$	$2.7 \times 10^8$	-0.72	a
$Ru(bpy)_3^{2+}$	$2.4 \times 10^9$	-0.81	73
$Rh_2(TMB)_4^{2+}$	$1.1 \times 10^9$	-1.25	48
$[Ir(\mu\text{-pz})COD]_2$	$8.7 \times 10^9$	-1.75	39

a. This work.

orbitals are the same, and oxidation involves comparable contractions of the metal-metal distance. For  $\text{Ru}(\text{bpy})_3^{2+}$ , the measured rate constant is an order of magnitude larger than that of  $\text{Ir}_2(\text{TMB})_4^{2+}$  and comparable to that for  $\text{Rh}_2(\text{TMB})_4^{2+}$ . However, the driving force for the excited-state electron transfer for  $\text{Ru}(\text{bpy})_3^{2+}$  is comparable to that for  $\text{Ir}_2(\text{TMB})_4^{2+}$ . Variations in factors other than driving force have a dramatic effect on the rate of reaction. The donor orbital for excited  $\text{Ru}(\text{bpy})_3^{2+}$  is located on one of the ligands,<sup>74</sup> which means that it is more accessible to the acceptor than the metal-localized orbital of the binuclear  $d^8$  complexes. (In other words, the electron-transfer reaction for  $\text{Ir}_2(\text{TMB})_4^{2+}$  may be more nonadiabatic than that one for  $\text{Ru}(\text{bpy})_3^{2+}$ .) To further understand the electron-transfer reactivity of the  $^3(d\sigma^*p\sigma)$  excited state of  $\text{Ir}_2(\text{TMB})_4^{2+}$ , a study of electron-transfer quenching by a series of pyridinium acceptors of variable reduction potential was undertaken. This study has allowed a better determination of the excited-state reduction potential ( $\text{M}_2^{3+} + e^- \rightarrow ^3(\text{M}_2^{2+})^*$ ); it has also elucidated the factors controlling the excited-state reactivity and has enabled comparisons of the electron-transfer reactivity of  $\text{Ir}_2(\text{TMB})_4^{2+}$  to be made with other inorganic systems.

Previous work has made available a series of pyridinium compounds that are well-suited for a systematic study.<sup>57</sup> The list of pyridinium quenchers and the measured  $E_{1/2}(\text{A}^{+/0})$  or  $E_{p,c}(\text{A}^{+/0})$  values for one-electron reductions of these acceptors is given in Table 4.5. Discussion of the possible complications of using a series of acceptors that are electrochemically irreversible is presented elsewhere.<sup>39</sup> No complications that are due to the dimerization of the pyridinium radical were observed.

Addition of the pyridiniums to an acetonitrile solution of  $\text{Ir}_2(\text{TMB})_4^{2+}$  does not alter the optical absorption spectrum of the metal complex. Emission spectra showed only quenching of the triplet excited state (Figure 4.15). No evidence for singlet

**Table 4.5.** Reduction Potentials of the Pyridinium Quenchers.

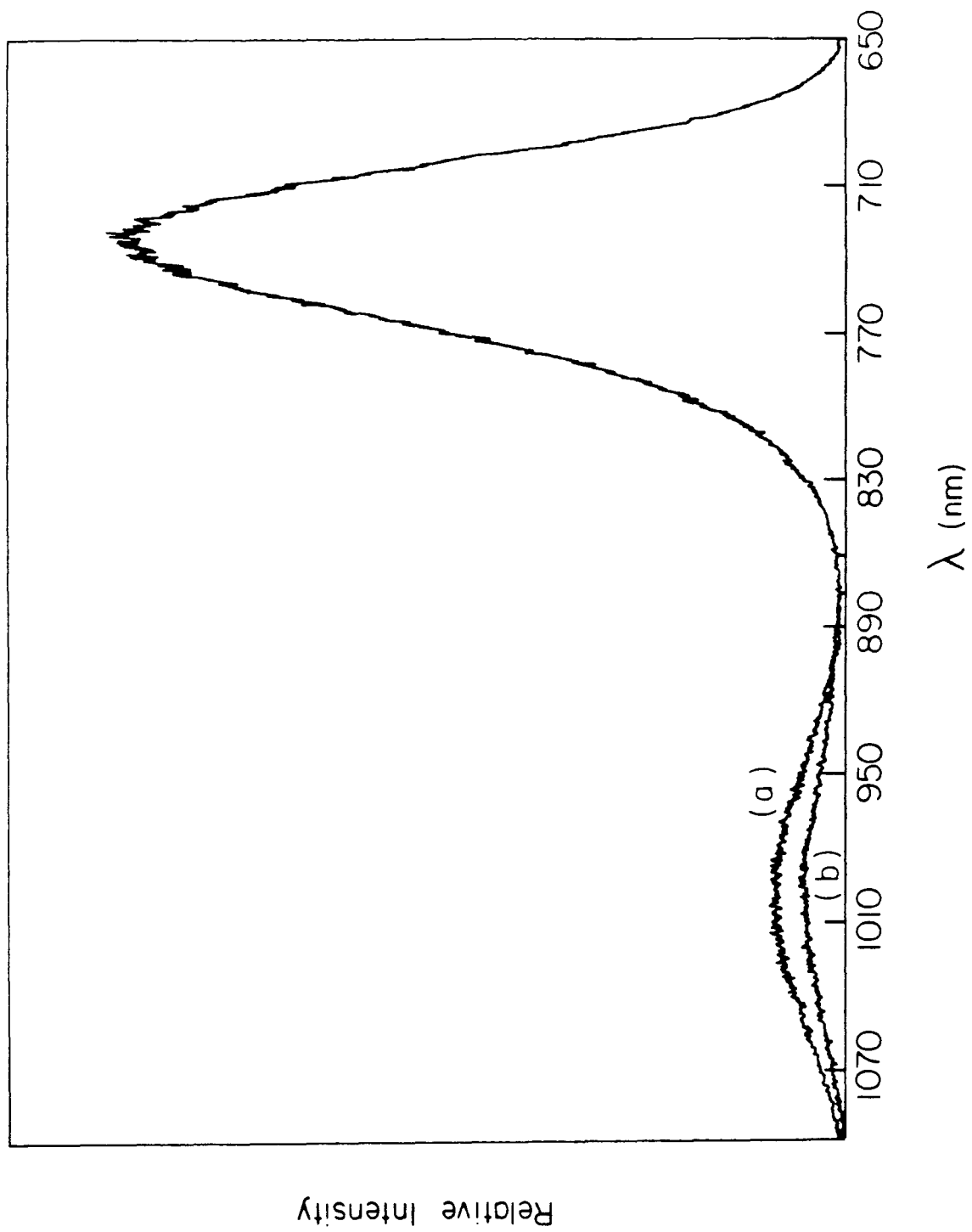
Quencher <sup>a</sup>		$E(A^{+/0})$ , V (SSCE) <sup>b</sup>
( 1 )	4-cyano-N-benzylpyridinium	-0.58
( 2 )	4-cyano-N-methylpyridinium	-0.65
( 3 )	4-carbomethoxy-N-methylpyridinium	-0.78
( 4 )	3-cyano-N-methylpyridinium	-0.90
( 5 )	4-amido-N-ethylpyridinium	-0.93
( 6 )	3-carbomethoxy-N-benzylpyridinium	-0.96
( 7 )	3-carbomethoxy-N-methylpyridinium	-1.01
( 8 )	3-amido-N-benzylpyridinium	-1.04

a. All of the compounds are hexafluorophosphate salts.

b. For Pyridiniums 1-3,  $E(A^{+/0}) = E_{1/2}(A^{+/0})$ . For pyridiniums 4-8 the reductions are irreversible; therefore the values of  $E(A^{+/0})$  are the cathodic peak potentials,  $E_{p,c}(A^{+/0})$  measured at a scan rate of 200 mV/s. Both  $E_{1/2}(A^{+/0})$  and  $E_{p,c}(A^{+/0})$  were measured by cyclic voltammetry (0.1 M TBAPF<sub>6</sub>) in acetonitrile solution.

**Figure 4.15.** (a) Emission spectrum of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in  $\text{CH}_3\text{CN}$  (0.1 M  $\text{TBAPF}_6$ ). (b) Emission spectrum of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in  $\text{CH}_3\text{CN}$  (0.1 M  $\text{TBAPF}_6$ ) with 4-CN-*N*-methylpyridinium.





reactivity was observed. Presumably, the singlet state does not react because of its short lifetime.

The kinetics of the excited state quenching of  $\text{Ir}_2(\text{TMB})_4^{2+}$  with the pyridinium acceptors was studied, using the same techniques as in the study of the quenching by  $\text{MV}^{2+}$ . All systems were found to obey Stern-Volmer kinetics over the quencher concentration range studied. A representative plot of the data for 4-CN-N-benzylpyridinium hexafluorophosphate is shown in Figure 4.16. The quenching rate constants, listed in Table 4.6, were obtained from the slopes of linear-least-squares fits to the data.

Excited-state electron transfer has been the subject of many reviews in the recent literature.<sup>46,47,50,75-77</sup> The general scheme presented to describe the kinetics is given in Figure 4.17. The quenching rate constant for the system is

$$k_q = \frac{k_d}{1 + \frac{k_{-d}}{k_e} \left( 1 + \frac{k_{-e}}{k_{-e}(g)} \right)} . \quad 4.11$$

For  $k_{-e} \ll k_{-e}(g)$ , the expression reduces to

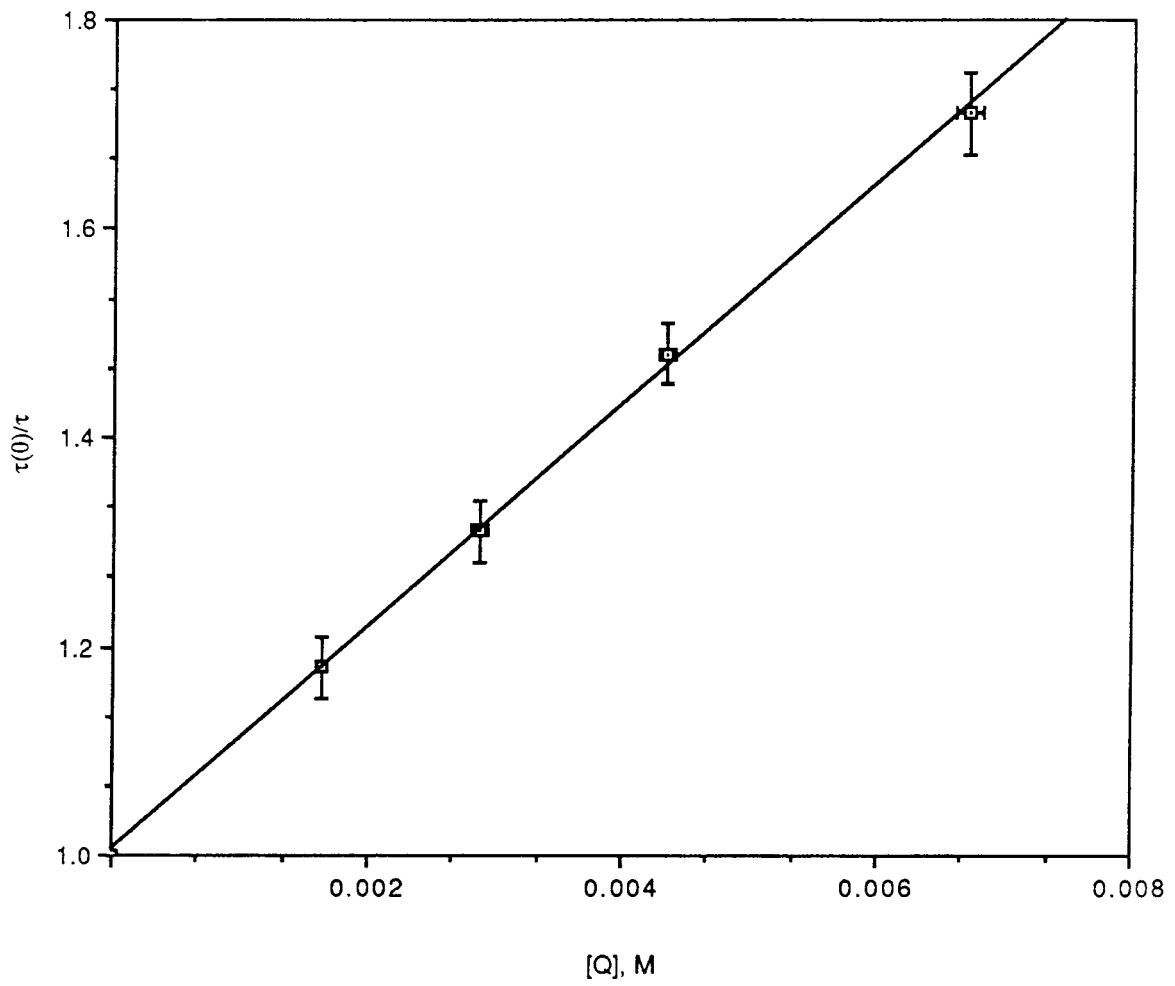
$$k_q = \frac{k_d}{1 + \frac{k_{-d}}{k_e}} \quad 4.12$$

$$\frac{1}{k_q} = \frac{1}{k_d} + \frac{1}{Kk_e} . \quad 4.13$$

Correcting the observed rate constants for diffusion, one obtains the activation-controlled rate constant ( $k_q'$ ) for the bimolecular reaction

$$k_q' = Kk_e , \quad 4.14$$

**Figure 4.16.** Stern-Volmer plot of oxidative quenching of the  $^3(d\sigma^*p\sigma)$  excited state of  $\text{Ir}_2(\text{TMB})_4^{2+}$  by 4-CN-*N*-benzylpyridinium (0.1 M TBAPF<sub>6</sub>, CH<sub>3</sub>CN).

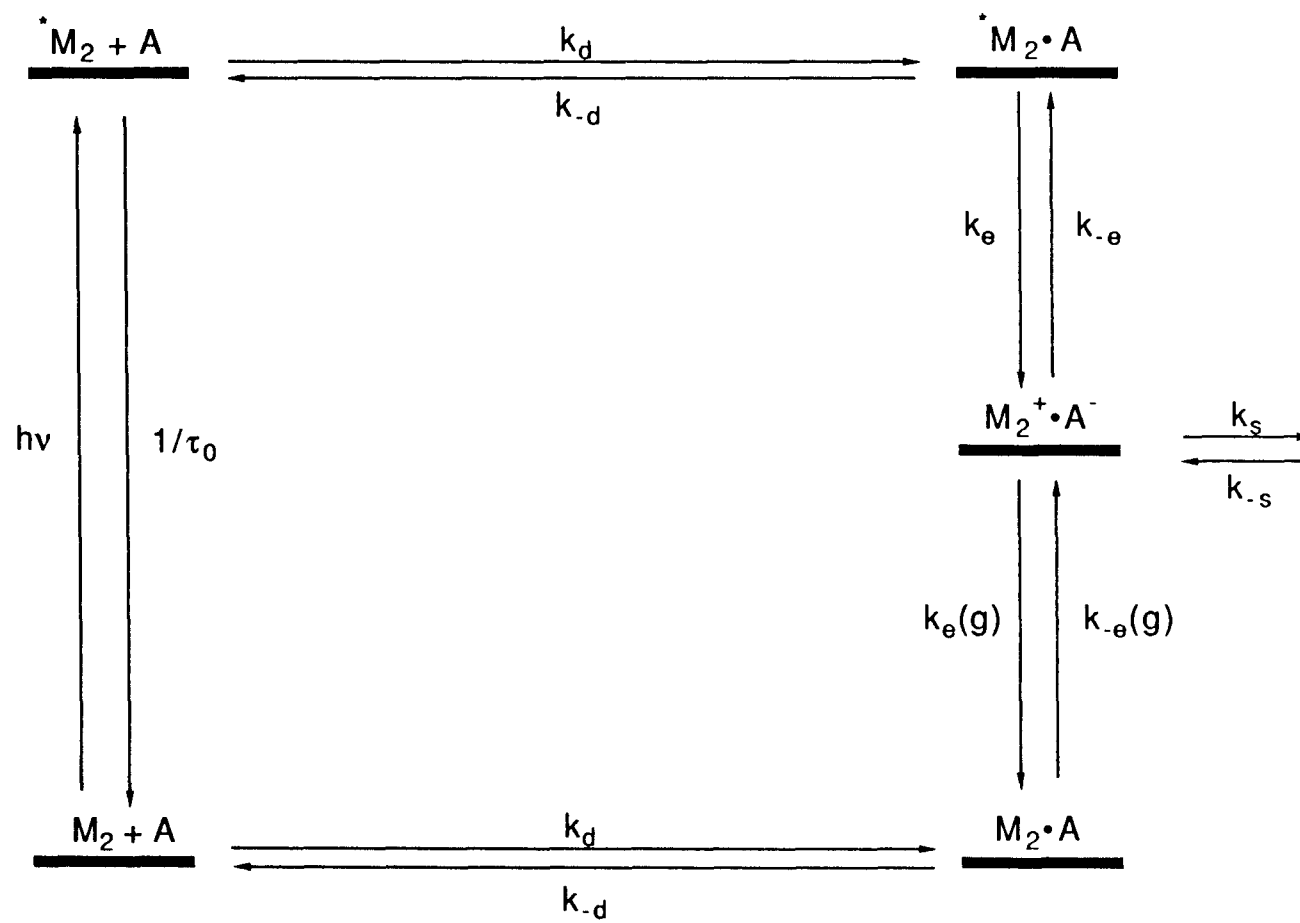


**Table 4.6.** Rate constants for the electron-transfer quenching of  $^3(\text{Ir}_2(\text{TMB})_4^{2+})^*$  by pyridinium acceptors in acetonitrile solution.

Pyridinium Acceptors <sup>a</sup>	$k_q, \text{M}^{-1}\text{s}^{-1}$ <sup>b</sup>	$k_q', \text{M}^{-1}\text{s}^{-1}$ <sup>c</sup>
1	$4.9 \pm 0.5 \times 10^8$	$5.0 \pm 0.5 \times 10^8$
2	$4.5 \pm 0.4 \times 10^8$	$4.6 \pm 0.5 \times 10^8$
3	$3.3 \pm 0.3 \times 10^8$	$3.4 \pm 0.3 \times 10^8$
4	$3.2 \pm 0.3 \times 10^8$	$3.3 \pm 0.3 \times 10^8$
5	$1.9 \pm 0.2 \times 10^8$	$1.9 \pm 0.2 \times 10^8$
6	$2.8 \pm 0.3 \times 10^7$	$2.8 \pm 0.3 \times 10^7$
7	$9.6 \pm 1.0 \times 10^6$	$9.6 \pm 1.0 \times 10^6$
8	$4.6 \pm 0.5 \times 10^6$	$4.6 \pm 0.5 \times 10^6$

- The acceptors are the pyridinium hexafluorophosphates listed in Table 4.5.
- The measured quenching rate constants ( $k_q$ ) are not corrected for diffusional effects.
- The quenching rate constants ( $k_q'$ ) are corrected for diffusional effects as discussed in the text.

**Figure 4.17.** Kinetic scheme for the electron-transfer quenching of  $^*M_2$  by pyridinium acceptors.



where  $K = \frac{k_d}{k_{-d}}$ . From a classical approach, the rate constant for the electron-transfer step is

$$k_e = \kappa \nu \exp\left(\frac{-\Delta G^*}{RT}\right), \quad 4.15$$

where  $\kappa$  is the transmission coefficient or average probability for electron transfer per passage of the system through the intersection region,  $\nu$  is the nuclear frequency,  $\Delta G^*$  is the free energy of activation that is related to the reorganization energy of the system.

To analyze the kinetic data, it is necessary to adopt a relationship for  $\Delta G^*$ .

Several free-energy relationships, FERs, have been presented in the literature with the pros and cons of each formalism discussed in some detail.<sup>50</sup> Two FERs are used to analyze the bimolecular electron-transfer kinetics. From theoretical considerations, Marcus has suggested a relationship for  $\Delta G^*$  in terms of  $\Delta G^{0'}$ , the standard free energy of the reaction,  $\lambda$ , the reorganization energy associated with the inner and outer coordination spheres,  $w_r$  and  $w_p$ , the work of bring reactants or products to the mean separation for reaction.<sup>78</sup>

$$\Delta G^* = \frac{\lambda}{4} \left( 1 + \left( \frac{\Delta G^{0'}}{\lambda} \right)^2 \right) \quad 4.16$$

$$\Delta G^{0'} = \Delta G^0 + w_p - w_r \quad 4.17$$

$$\Delta G^0 = E^0(\text{Ir}_2^{3+/3}(\text{Ir}_2^{2+})^*) - E^0(\text{A}^{+/0}) - w_r \quad 4.18$$

$$k_q' = K \kappa \nu \exp\left(\frac{\lambda}{4RT} \left( 1 + \frac{\Delta G^{0'}}{\lambda} \right)^2 \right) \quad 4.19$$

$$RT \ln(k_q') = RT \ln(K \kappa \nu) + \frac{\lambda}{4} \left( 1 + \frac{\Delta G^{0'}}{\lambda} \right)^2. \quad 4.20$$



Rehm and Weller, studying the fluorescence quenching of organic compounds, have suggested an empirical FER.<sup>79</sup>

$$\Delta G^* = \frac{\Delta G^{0'}}{2} + \left[ \left( \frac{\Delta G^{0'}}{2} \right)^2 + \left( \frac{\lambda}{4} \right)^2 \right]^{1/2} \quad 4.21$$

$$k_q' = K_{kv} \exp \left\{ \frac{\Delta G^{0'}}{2RT} + \frac{1}{RT} \left[ \left( \frac{\Delta G^{0'}}{2} \right)^2 + \left( \frac{\lambda}{4} \right)^2 \right]^{1/2} \right\} \quad 4.22$$

$$RT \ln(k_q') = RT \ln(K_{kv}) + \frac{\Delta G^{0'}}{2} + \left[ \left( \frac{\Delta G^{0'}}{2} \right)^2 + \left( \frac{\lambda}{4} \right)^2 \right]^{1/2} \quad 4.23$$

Using either the Marcus or the Rehm-Weller expression, the rate of electron transfer can be analyzed as a function of the driving force of the reaction,  $-\Delta G^0$ . An analysis of  $RT \ln(k_q')$  as a function of the ease of reduction of the pyridinium acceptor gives the excited-state reduction potential, the reorganization energy associated with the electron-transfer reaction, and the preexponential factor, assuming that the associative equilibrium constant can be estimated.<sup>47,50,80</sup>

The two FERs predict different behavior for the rate of electron transfer as a function of  $\Delta G^{0'}$ . The most pronounced difference occurs at high exoergicity. The Marcus relationship predicts that the reaction rate will reach a maximum for  $\Delta G^{0'} = -\lambda$  and decrease for increasing exoergicity. For the Rehm-Weller empirical FER, the rate constant tends asymptotically to its maximum value with increasing exoergicity.

Rate constants corrected for the diffusional contribution are listed in Table 4.6. The diffusion-controlled rate constant can be calculated using the Debye-Smoluchowski equation<sup>81</sup>

$$k_d = \frac{216}{1000} \frac{4\pi N_0 (D_A + D_B) a f}{1000} \quad 4.24$$

$N_0$  = Avogadro's number.

$a$  = distance of closest approach,  $r_A + r_B$ .

$D_A, D_B$  = diffusion coefficient for the ions calculated, using the Stokes-Einstein equation,

where  $f$  is the integral

$$\left[ a \int_a^\infty \exp\left(\frac{U}{k_b T}\right) \frac{dr}{r^2} \right]^{-1} \quad 4.25$$

$$U = \left( \frac{z_A z_B e^2}{D_s} \right) \left[ \frac{\exp(Ka)}{(1 + Ka)} \right] \left( \frac{\exp(-Kr)}{r} \right)$$

$r$  = distance of separation of the ions.

$z_A, z_B$  = charge of the ions.

$e$  = unit electron charge.

$D_s$  = static dielectric of the medium.

$$K = \left[ \frac{8\pi N_0 e^2 \mu}{1000 D_s k_b T} \right]^{1/2}$$

$\mu$  = ionic strength.

$k_b$  = Boltzmann's constant.

$T$  = temperature.

$k_d$  may be evaluated using the Eigen equation.<sup>82</sup>

$$k_d = \left[ \frac{3(D_A + D_B)f}{a^2} \right] \exp\left[ \frac{b}{(1 + Ka)} \right] \quad 4.26$$

$$b = \frac{z_A z_B e^2}{a D_s k_b T}$$

The equilibrium constant for association is then

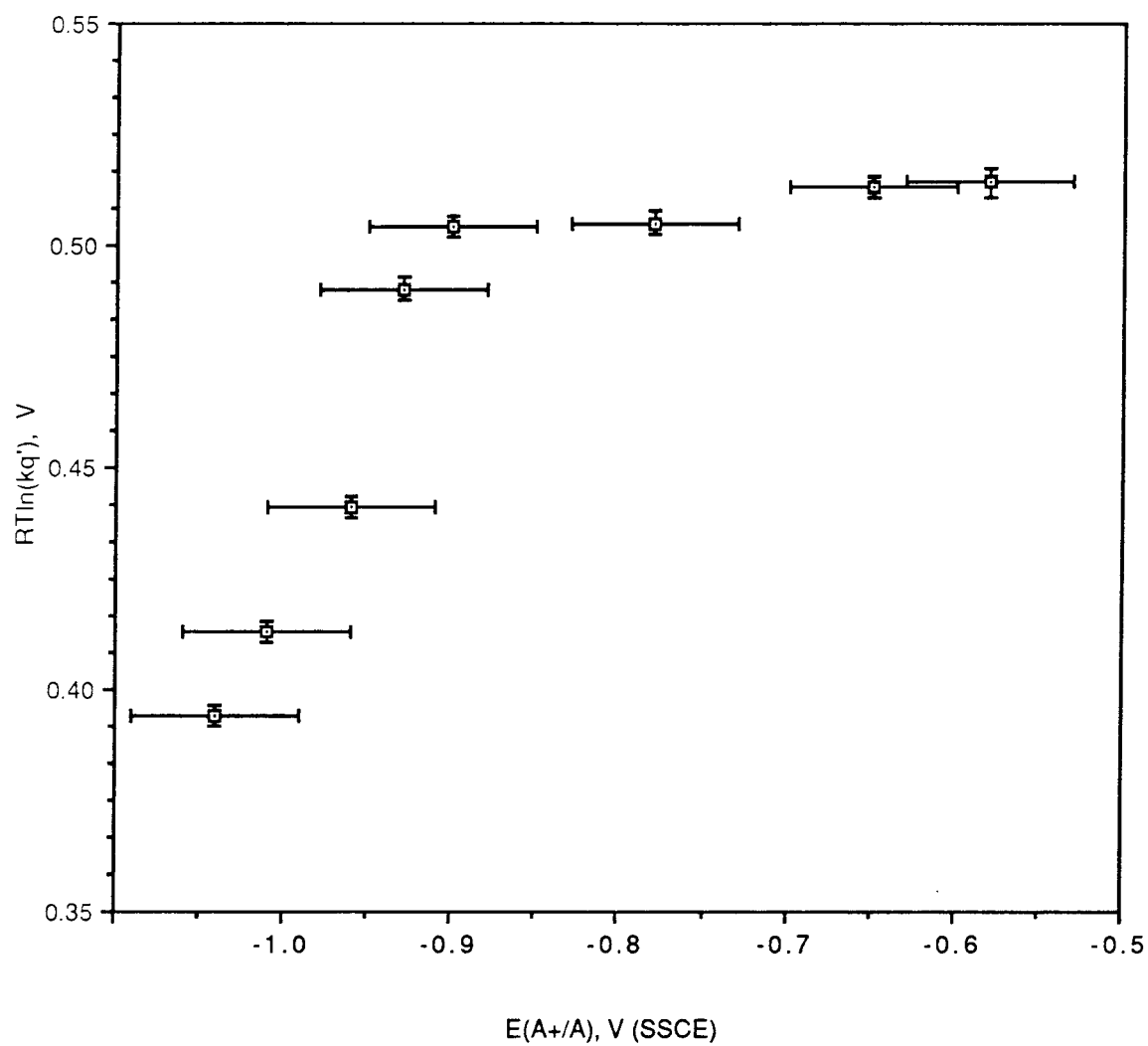
$$\frac{k_d}{k_{-d}} = \left( \frac{4\pi a^3 N_0}{3000} \right) \exp \left[ \frac{-b}{(1 + K a)} \right] . \quad 4.27$$

The equilibrium constant for association calculated for  $\text{Ir}_2(\text{TMB})_4^{2+}$  and the pyridinium quenchers is approximately 1 to 2  $\text{M}^{-1}$ . See Appendix 2 for a numerical integration routine to evaluate  $f$ . Also included in Appendix 2 are the calculated constants for the various pyridinium acceptors.<sup>83</sup>

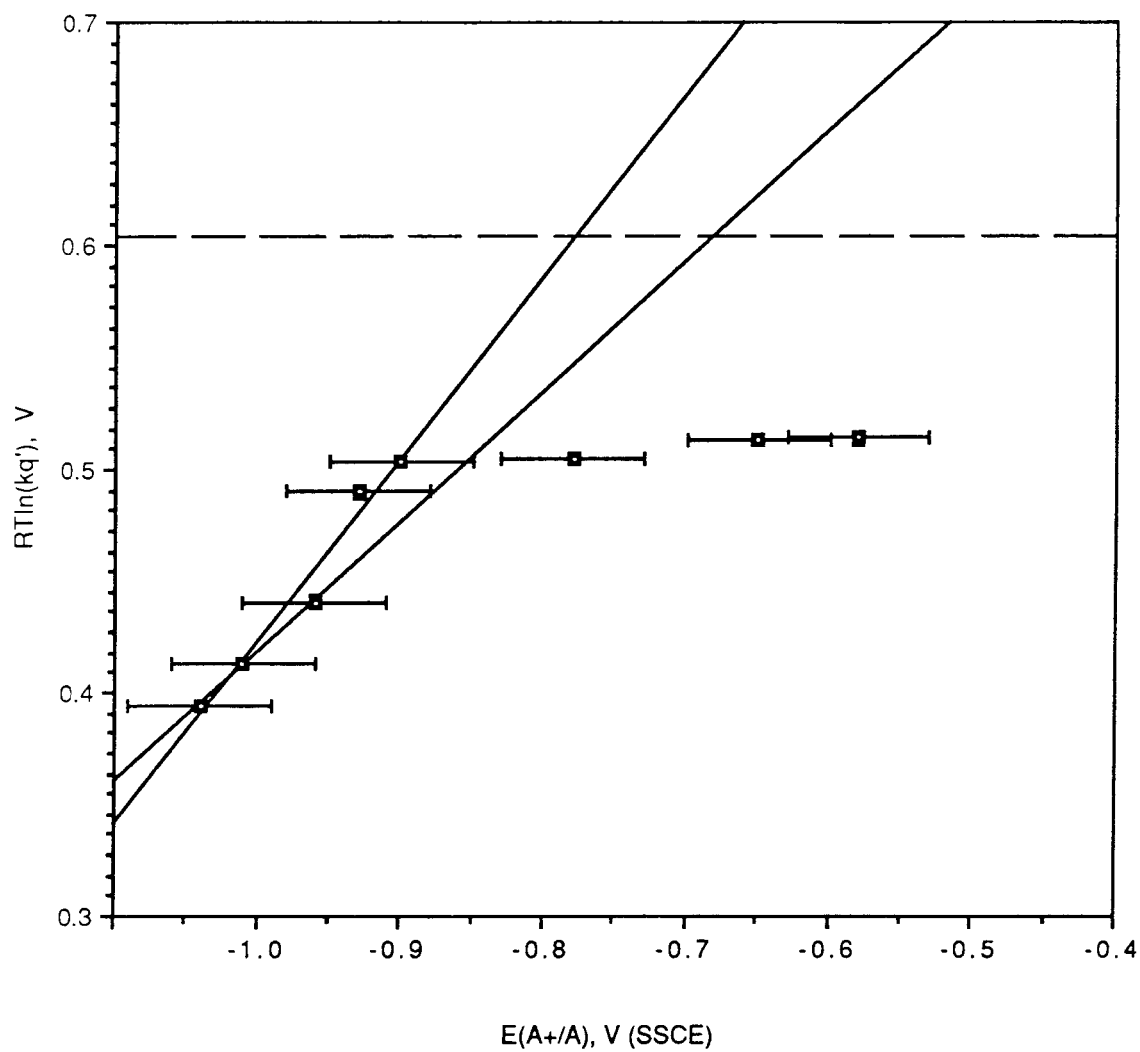
A plot of  $RT \ln(k_q')$  versus  $E^0(\text{A}^{+/0})$  is shown in Figure 4.18. Two regions are observed; (i) an activated region that shows Arrhenius-type linear behavior, and (ii) a plateau region where the rate shows no change for a decrease in  $E^0(\text{A}^{+/0})$ . Lacking in these data is a region, predicted by both the Marcus and Rehm-Weller FERs, intermediate between (i) and (ii), where  $k_q'$  increases nonlinearly as  $E^0(\text{A}^{+/0})$  decreases.

Within the Marcus formalism, a plot of  $RT \ln(k_q')$  versus  $E^0(\text{A}^{+/0})$  should yield a line of slope 0.5 in the region where  $\Delta G^0' \ll 2\lambda$ .<sup>73</sup> A plot of the data with linear-least-squares fits to the activated region is shown in Figure 4.19. The slope is  $0.7 \pm 0.1$ . This poor agreement to the predicted value of 0.5 is not an indication that the Rehm-Weller formalism is a better description of the system.<sup>50</sup> Earlier work for  $[\text{Ir}(\mu\text{-pz})\text{COD}]_2$  found excellent agreement between the experimentally determined slope (0.48) and the theoretical value of 0.5.<sup>57</sup> The uncertainty in the value of the slope determined for  $\text{Ir}_2(\text{TMB})_4^{2+}$  arises from the limited amount of data available in the activated region and the large uncertainty in  $E^0(\text{A}^{+/0})$ .

**Figure 4.18.** Plot of  $RT\ln(k_q')$  (V) versus  $E^0(A^{+/0})$  (V) for the electron-transfer quenching of  ${}^3(\text{Ir}_2(\text{TMB})_4{}^{2+})^*$  by the pyridinium acceptors in Table 4.5 ( $\mu = 0.1 \text{ M TBAPF}_6, \text{CH}_3\text{CN}$ ).



**Figure 4.19.** Plot of  $RT\ln(k_q')$  (V) versus  $E^0(A^{+/0})$  (V) for the electron-transfer quenching of  $^3(\text{Ir}_2(\text{TMB})_4^{2+})^*$  by the pyridinium acceptors in Table 4.5 ( $\mu = 0.1 \text{ M TBAPF}_6$ ,  $\text{CH}_3\text{CN}$ ). Diffusion-limited rate shown as the horizontal dashed line. Two linear fits to the activated region are shown.



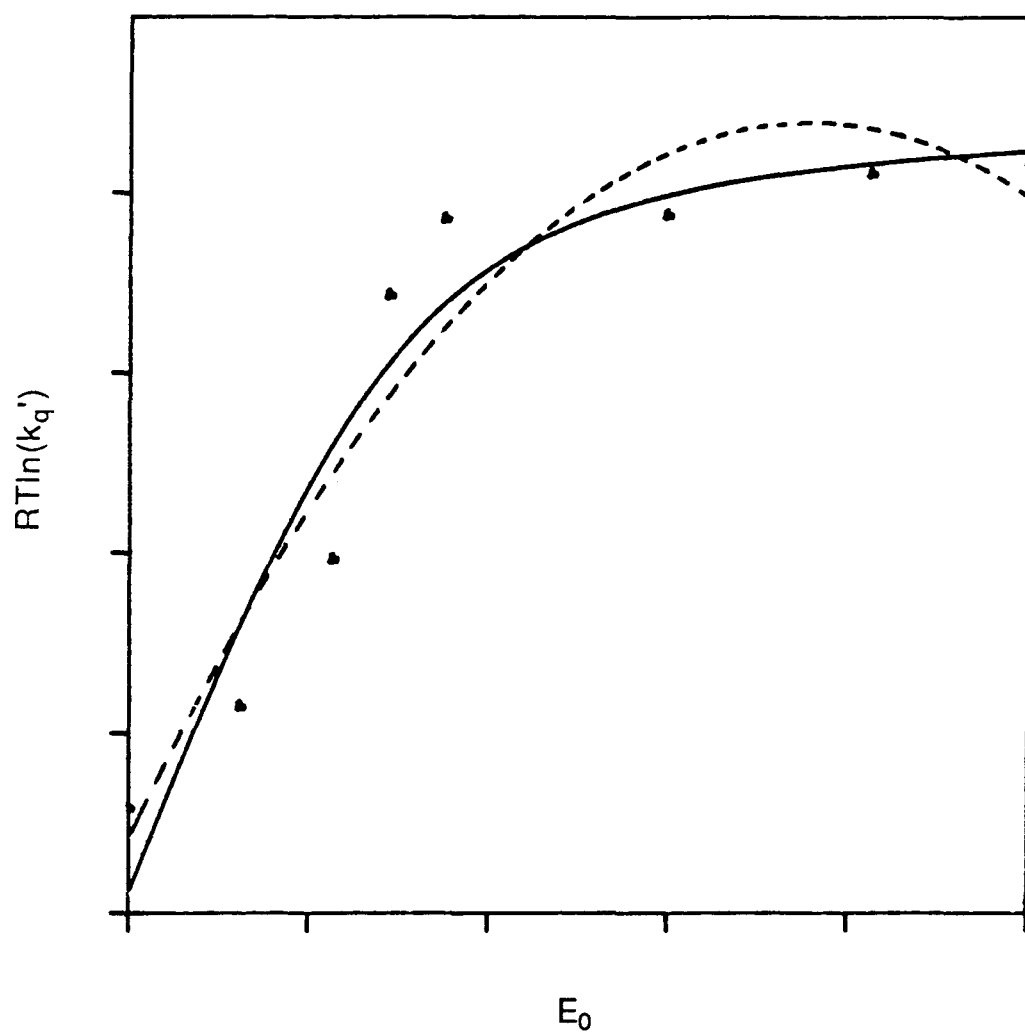
Three-parameter, nonlinear-least-squares fits of  $RT\ln(k_q')$  versus  $E^0(A^{+/0})$  using Equations 4.20 and 4.23 are shown in Figure 4.20. The work term associated with bringing  $\text{Ir}_2(\text{TMB})_4^{2+}$  and a pyridinium acceptor together is 0.01 to 0.03 eV.<sup>46</sup> This contribution is neglected in the analysis of the electron-transfer rate data.

The fitted values for  $RT\ln(K_{\kappa\nu})$ ,  $E^0(\text{Ir}_2^{3+/3}(\text{Ir}_2^{2+})^*)$ , and  $\lambda$  are listed in Table 4.7. The fitted values for the two FERs are almost identical with the largest variation observed for  $\lambda$ . It is satisfying to find that the fitted value of  $E^0(\text{Ir}_2^{3+/3}(\text{Ir}_2^{2+})^*)$  agrees with the earlier estimate of -0.72 V (SSCE). Variances of 200- 300 mV in reduction potentials are observed for other systems.<sup>57</sup> The surprising result from the nonlinear-least-squares fits is the low value for the reorganization energy ( $\lambda \sim 0.25$  eV). Earlier studies for analogous binuclear systems found  $\lambda$ 's of approximately 1 eV.<sup>57,84</sup> For small  $\lambda$  values, the intermediate curved region will decrease, and for very small values of  $\lambda$ , the intermediate range will be almost unnoticeable (Figure 4.21).<sup>50</sup> For  $[\text{Ir}(\mu\text{-pz})\text{COD}]_2$ , a broad, nonlinear region is observed (Figure 4.22). The curvature observed for  $\text{Ir}_2(\text{TMB})_4^{2+}$  is in stark contrast to that for  $[\text{Ir}(\mu\text{-pz})\text{COD}]_2$  (Figure 4.22); however, this may be a consequence of the low limiting rate, not a small value of  $\lambda$ . In Figure 4.23, several two-parameter, nonlinear-least-squares fits to the data for  $\text{Ir}_2(\text{TMB})_4^{2+}$  are shown. The value of  $\lambda$  is fixed, and  $E^0(\text{Ir}_2^{3+/3}(\text{Ir}_2^{2+})^*)$  and  $RT\ln(\kappa\nu)$  are optimized according to Equation 4.20. Reasonable fits to the observed data are obtained for  $\lambda$  from 0.5 to 0.8 eV (Table 4.8). Thus, the reorganization energy for the oxidative electron-transfer quenching of  $^3(\text{Ir}_2(\text{TMB})_4^{2+})^*$  by pyridinium acceptors may be comparable to that determined for other binuclear  $d^8$  complexes. Also, for a large value of  $\lambda$ , the fitted Marcus curve does not show an artificial downward curvature at higher driving force as for  $\lambda \sim 0.2$  eV.

The limiting rate observed for  $\text{Ir}_2(\text{TMB})_4^{2+}$  is approximately two orders of magnitude below the calculated diffusional rate constant. For the plateau region two limiting values are predicted. As  $\Delta G^{0'} \rightarrow -\infty$ ,



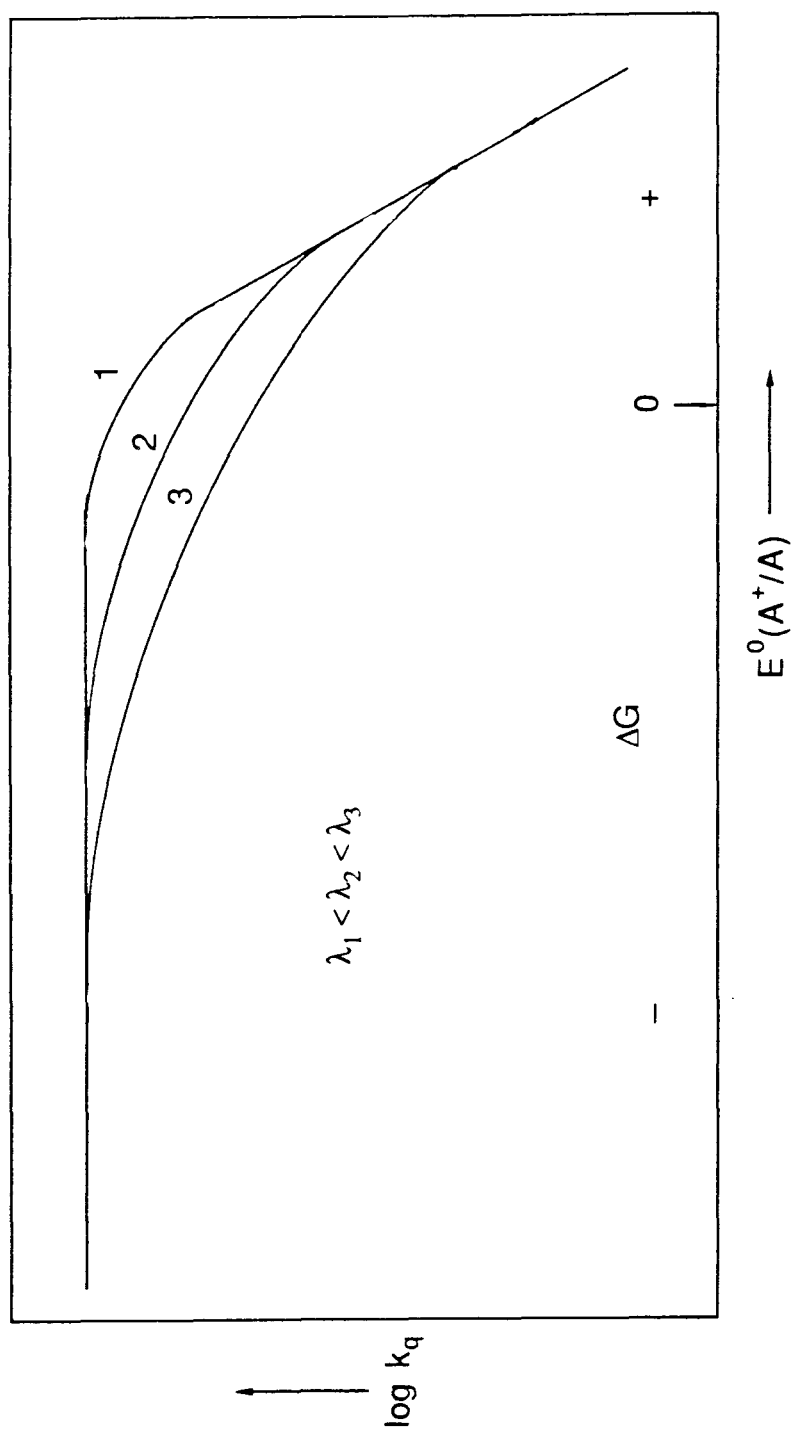
**Figure 4.20.** Plot of  $RT\ln(k_q')$  (V) versus  $E^0(A^{+/0})$  (V) for the electron-transfer quenching of  ${}^3(\text{Ir}_2(\text{TMB})_4{}^{2+})^*$  by the pyridinium acceptors in Table 4.5 ( $\mu = 0.1 \text{ M TBAPF}_6$ ,  $\text{CH}_3\text{CN}$ ). Three-parameter, nonlinear-least-squares fit using Marcus formalism (dashed curve) and Rehm-Weller (solid curve).



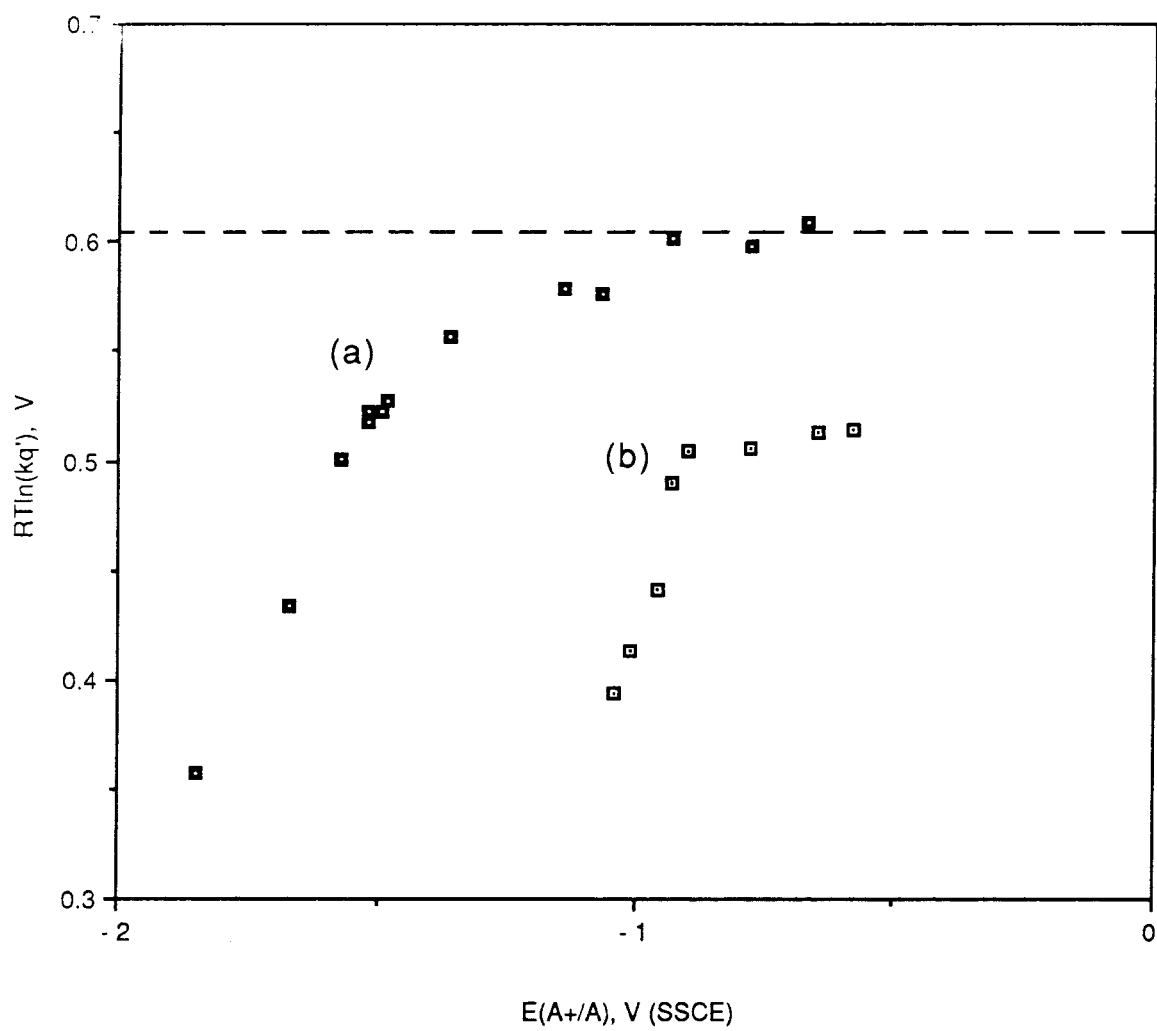
**Table 4.7.** Values for  $RT\ln(K_{\text{Kv}})$ ,  $E^0(\text{Ir}_2^{3+}/^3(\text{Ir}_2^{2+})^*)$ , and  $\lambda$  for  $\text{Ir}_2(\text{TMB})_4^{2+}$  obtained from three-parameter, nonlinear-least-squares fits.

	Marcus	Rehm-Weller
$RT\ln(K_{\text{Kv}})$ , V (SSCE)	0.52	0.52
$E^0(\text{Ir}_2^{3+}/^3(\text{Ir}_2^{2+})^*)$ , V (SSCE)	-0.95	-0.94
$\lambda$ , eV	0.25	0.21

**Figure 4.21.** Influence of  $\lambda$  on  $\log(k_q)$  versus  $\Delta G$  plots.

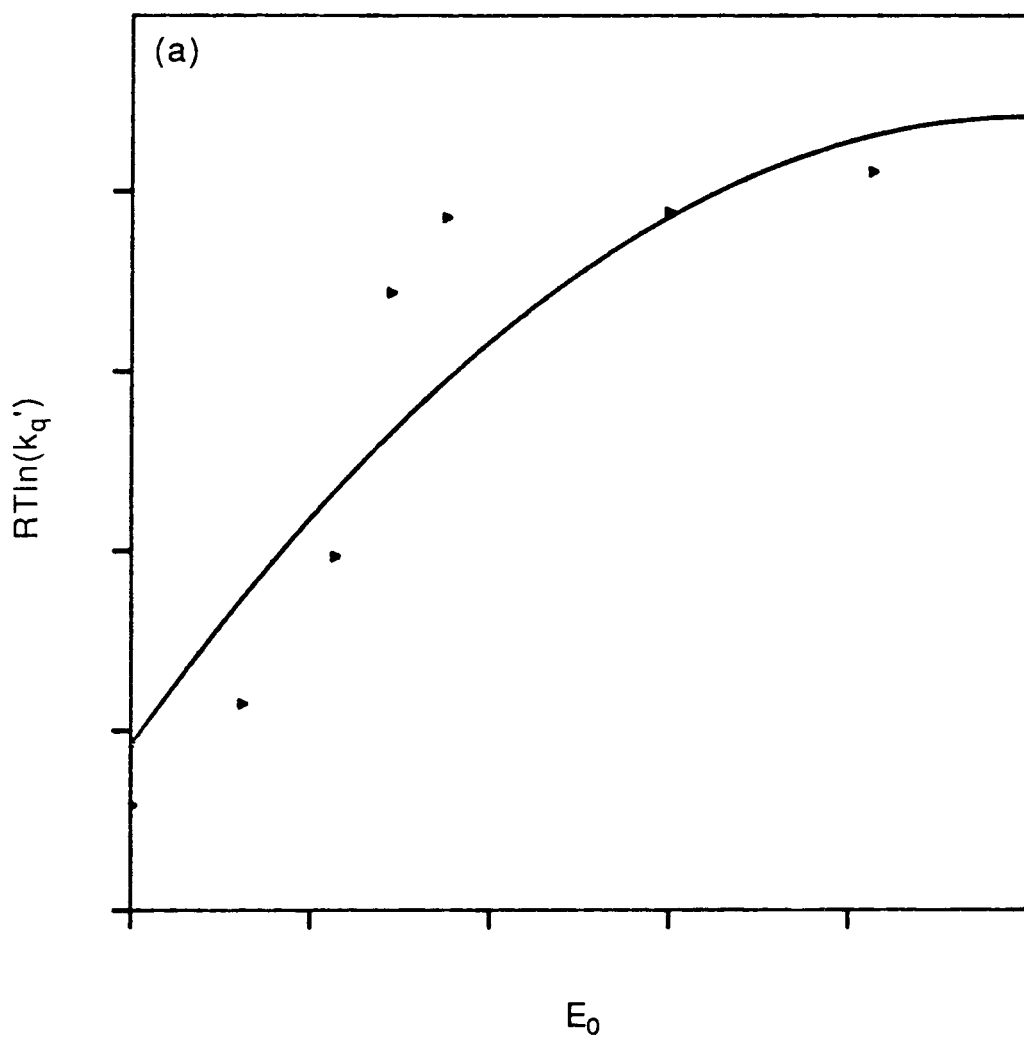


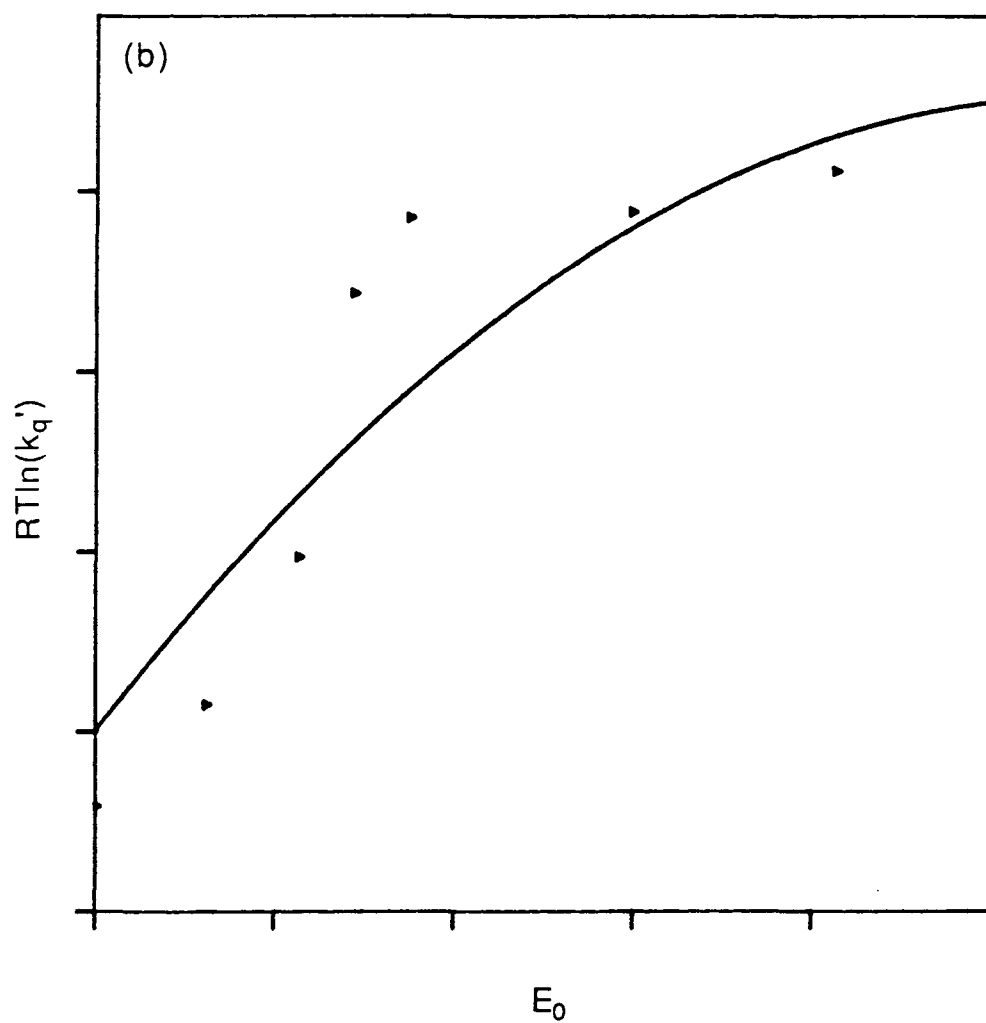
**Figure 4.22.** Plot of  $RT\ln(k_q')$  versus  $E^0(A^{+/0})$  for the electron-transfer quenching of  $^3[\text{Ir}(\mu\text{-pz})\text{COD}]_2^*$  (curve a) and  $^3(\text{Ir}_2(\text{TMB})_4^{2+})^*$  (curve b) by the pyridinium acceptors. Diffusion-limited rate horizontal dashed line.

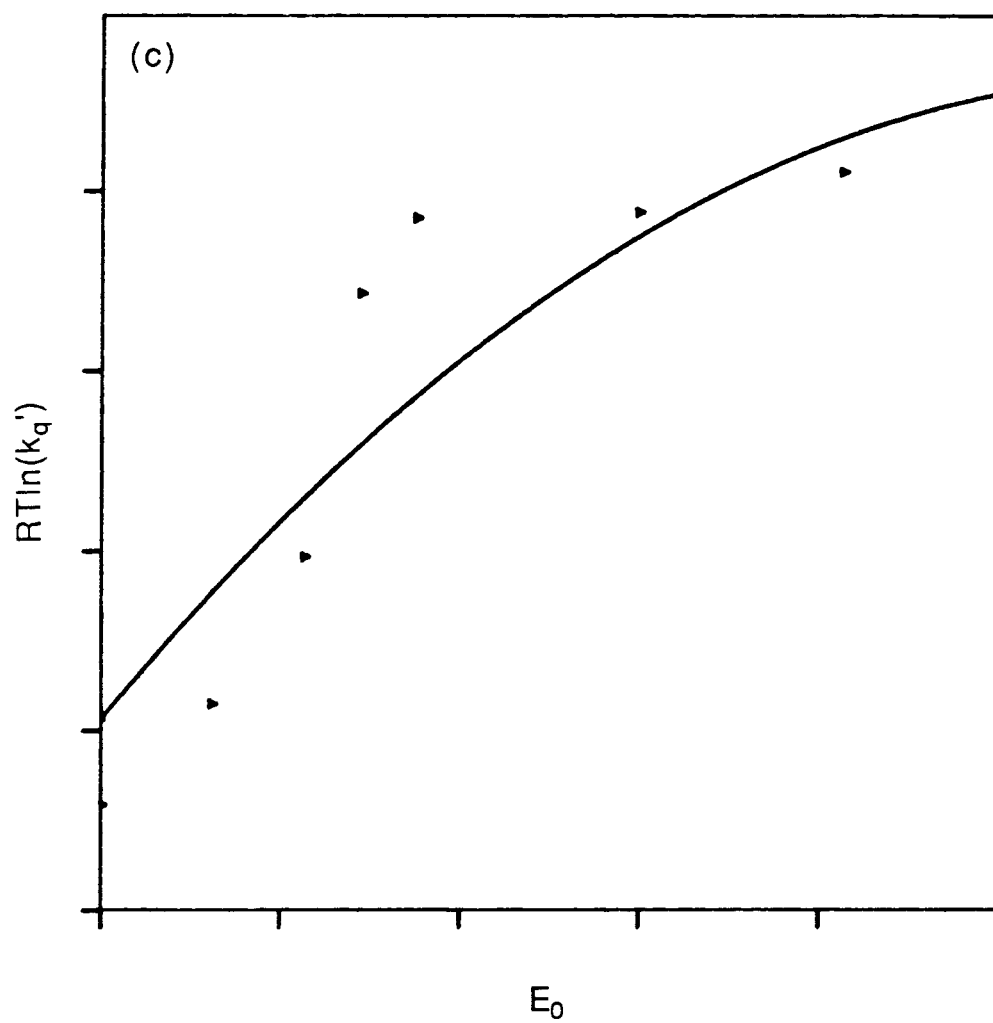


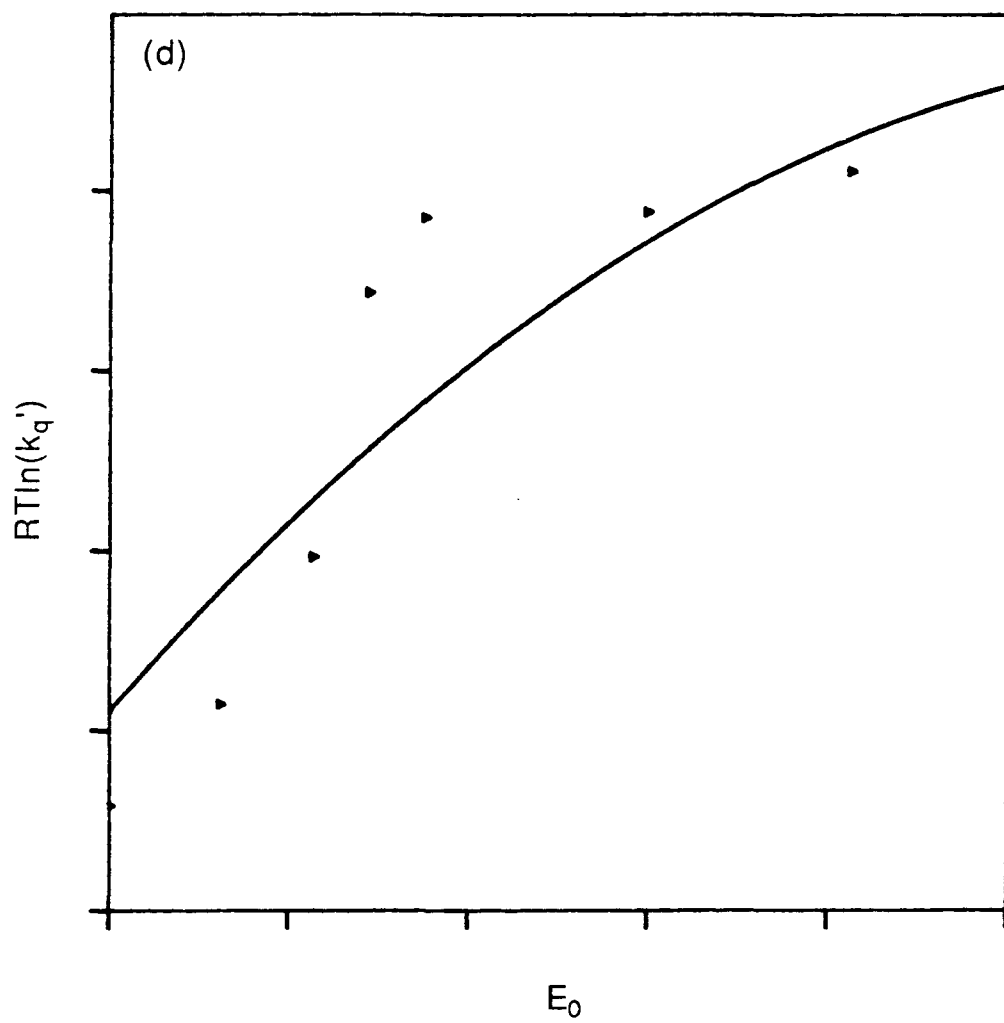
**Figure 4.23.** Plot of  $RT\ln(k_q')$  (V) versus  $E^0(A^{+/0})$  (V) for the electron-transfer quenching of  $^3(\text{Ir}_2(\text{TMB})_4^{2+})^*$  by the pyridinium acceptors in Table 4.5 ( $\mu = 0.1 \text{ M TBAPF}_6$ ,  $\text{CH}_3\text{CN}$ ). Two-parameter, nonlinear-least-squares fits using Marcus formalism for (a)  $\lambda = 0.5 \text{ eV}$ , (b)  $\lambda = 0.6 \text{ eV}$ , (c)  $\lambda = 0.7 \text{ eV}$ , and (d)  $\lambda = 0.8 \text{ eV}$ .











**Table 4.8.** Values for  $RT\ln(K_{\kappa\nu})$  and  $E^0(\text{Ir}_2^{3+}/3(\text{Ir}_2^{2+})^*)$  for  $\text{Ir}_2(\text{TMB})_4^{2+}$  obtained from two-parameter, nonlinear-least-squares fits to Marcus equation with fixed  $\lambda$  values.

$\lambda$ , eV	$RT\ln(K_{\kappa\nu})$ , V (SSCE)	$E^0(\text{Ir}_2^{3+}/3(\text{Ir}_2^{2+})^*)$ , V (SSCE)
0.5	0.52	-1.08
0.6	0.53	-1.14
0.7	0.53	-1.19
0.8	0.53	-1.24

$$k_q = \frac{k_d \kappa v}{(\kappa v + k_{-d})} \quad 4.28$$

for  $\kappa v > k_{-d}$

$$k_q = k_d$$

for  $\kappa v < k_{-d}$

$$k_q = K \kappa v.$$

The limiting rate for  $\text{Ir}_2(\text{TMB})_4^{2+}$  is  $K \kappa v$ , indicating a highly nonadiabatic process,  $\kappa \ll 1$ .

A value for the nonadiabaticity of the electron-transfer reaction can be estimated from the value of  $RT \ln(K \kappa v)$ . The values of  $RT \ln(K \kappa v)$  for  $[\text{Ir}(\mu\text{-pz})\text{COD}]_2$  determined for both the Marcus and Rehm-Weller FERs are listed in Table 4.9 (Figure 4.24). The electron-transfer reaction for  $[\text{Ir}(\mu\text{-pz})\text{COD}]_2$  is found to be nonadiabatic with  $\kappa \approx 0.01$ . A preexponential factor of  $5 \times 10^8 \text{ s}^{-1}$  is estimated for  $\text{Ir}_2(\text{TMB})_4^{2+}$ , yielding a  $\kappa \approx 0.0001$ , indicative of a highly nonadiabatic process. The small values of  $\kappa$  may be understood in terms of the semiclassical Landau-Zener theory.<sup>85</sup> The probability of the reactants' being converted to products per single passage through the intersection region is given by

$$P_{if} = 1 - \exp\left[\frac{-4\pi^2 H_{if}^2}{h v |s_i - s_f|}\right], \quad 4.29$$

where  $H_{if}$  is the electronic coupling matrix element between the initial reactant surface, and the final product surface.  $\kappa$  is given by

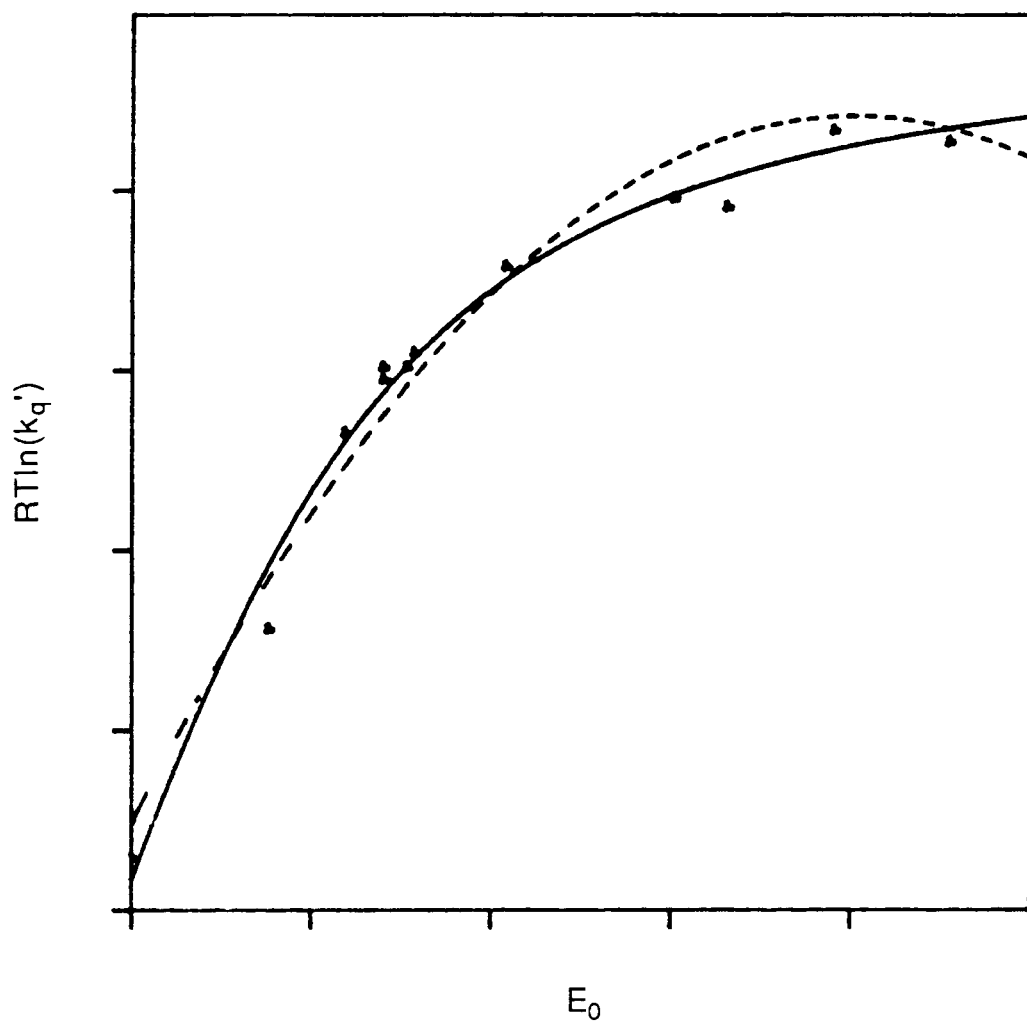
$$\kappa = \frac{2P_{if}}{(1 + P_{if})}, \quad 4.30$$

**Table 4.9.** Values for  $RT\ln(K_{\kappa\nu})$ ,  $E^0(\text{Ir}_2^{3+}/^3(\text{Ir}_2^{2+})^*)$ , and  $\lambda$  for  $[\text{Ir}(\mu\text{-pz})\text{COD}]_2$  obtained from three-parameter, nonlinear-least-squares fits.

	Marcus	Rehm-Weller
$RT\ln(K_{\kappa\nu})$ , V (SSCE)	0.61	0.64
$E^0(\text{Ir}_2^{3+}/^3(\text{Ir}_2^{2+})^*)$ , V (SSCE)	-1.84	-1.69
$\lambda$ , eV	0.94	0.78

**Figure 4.24.** Plot of  $RT\ln(k_q')$  (V) versus  $E^0(A^{+/0})$  (V) for the electron-transfer quenching of  $^3[\text{Ir}(\mu\text{-pz})\text{COD}]_2^*$  by the pyridinium acceptors. Three-parameter, nonlinear-least-squares fit using Marcus formalism (dashed curve) and Rehm-Weller (solid curve).





For strong coupling between the initial and final states,  $\kappa$  is unity; when  $H_{if}$  is small (weak coupling),  $\kappa$  is also small. The origin of the small value of  $\kappa$  for  $[\text{Ir}(\mu\text{-pz})\text{COD}]_2$  is not well understood. Weak electronic coupling might be expected for a donor and acceptor of like charge because of the Coulombic interaction that would keep them far apart, thus the very small value of  $\kappa$  observed for  $\text{Ir}_2(\text{TMB})_4^{2+}$ .

An estimate of the donor-acceptor separation for the electron-transfer event can be obtained from the value of  $\kappa$ .<sup>78</sup> In the nonadiabatic region,  $\kappa v$  is given by

$$\kappa v = 2\pi \frac{H_{ab}^2}{h(4\pi\lambda RT)^{1/2}} \quad 4.31$$

$H_{ab}^2$  decreases exponentially with separation distance for many systems, and so  $\kappa v$  can be written as

$$\kappa v = v_0 \exp[-\beta(r - r_0)] \quad , \quad 4.32$$

where  $r_0$  is the value of  $r$  at which  $\kappa v$  equals some preassigned value. For bimolecular reactions  $r_0$  is frequently assumed to correspond to close contact of the two reactants. Therefore,  $r - r_0$  corresponds to the donor-acceptor separation. Assuming a  $\beta$  of  $1.2 \text{ \AA}^{-1}$ ,  $r - r_0$  for the electron transfer reaction of  $[\text{Ir}(\mu\text{-pz})\text{COD}]_2$  is approximately  $4 \text{ \AA}$ . A closest contact ( $r = r_0$ ) encounter complex is not formed for the reaction between  $^3[\text{Ir}(\mu\text{-pz})\text{COD}]_2^*$  and a pyridinium acceptor. For  $\text{Ir}_2(\text{TMB})_4^{2+}$ ,  $r - r_0$  is calculated to be  $\sim 8 \text{ \AA}$ , a reasonable value considering the unfavorable Coulombic interaction between  $\text{Ir}_2(\text{TMB})_4^{2+}$  and the pyridinium cation.

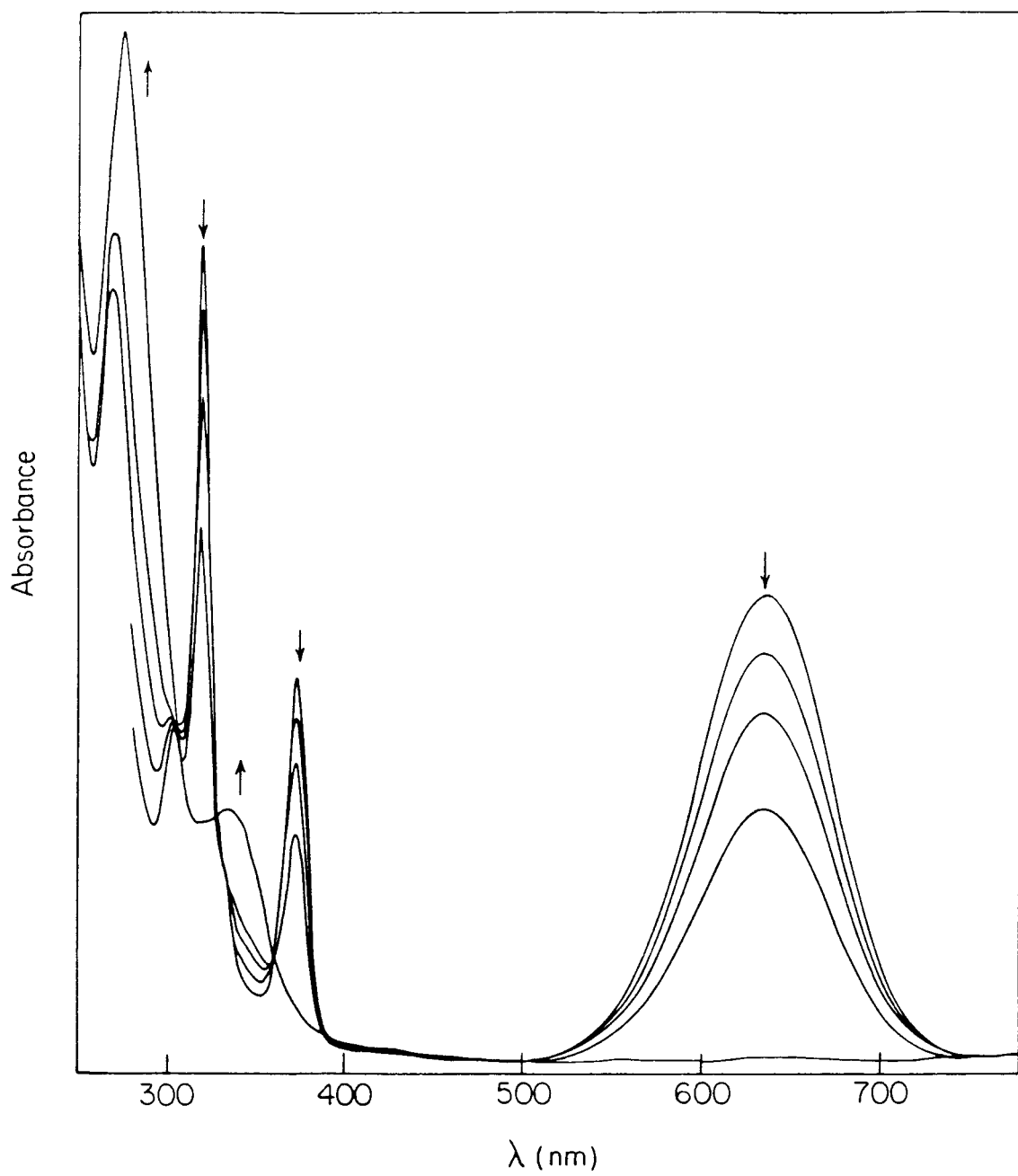
## Alkyl-Halide Photoredox

Although the conversion of light to chemical energy via the electron transfer reaction of binuclear  $d^8$  complexes is rather facile, the photochemical products return to starting materials because the back-electron-transfer reactions are very rapid. In order to utilize the strong reducing power of the  $^3(d\sigma^*p\sigma)$  excited state, the nonproductive back electron transfer must be inhibited. The first approach taken was to intercept the  $d^8$ - $d^7$  complex. This species, possessing open coordination sites, might be trapped by Lewis bases. Photolysis of  $[\text{Ir}(\mu\text{-pz})\text{COD}]_2$  with  $\text{MV}^{2+}$  in the presence of an excess of  $\text{I}^-$  did not yield  $[\text{Ir}(\mu\text{-pz})\text{COD}(\text{I})]_2$ .<sup>39</sup> Stern-Volmer quenching of the excited state was observed with rapid, back-electron transfer to yield the starting materials. The back reaction between  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{3-}$  and  $\text{A}^-$  could be inhibited by axial ligand binding.<sup>86</sup> A fine balance between productive scavenging of  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{3-}$  and nonproductive back reaction to  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$  and  $\text{A}$  was found.

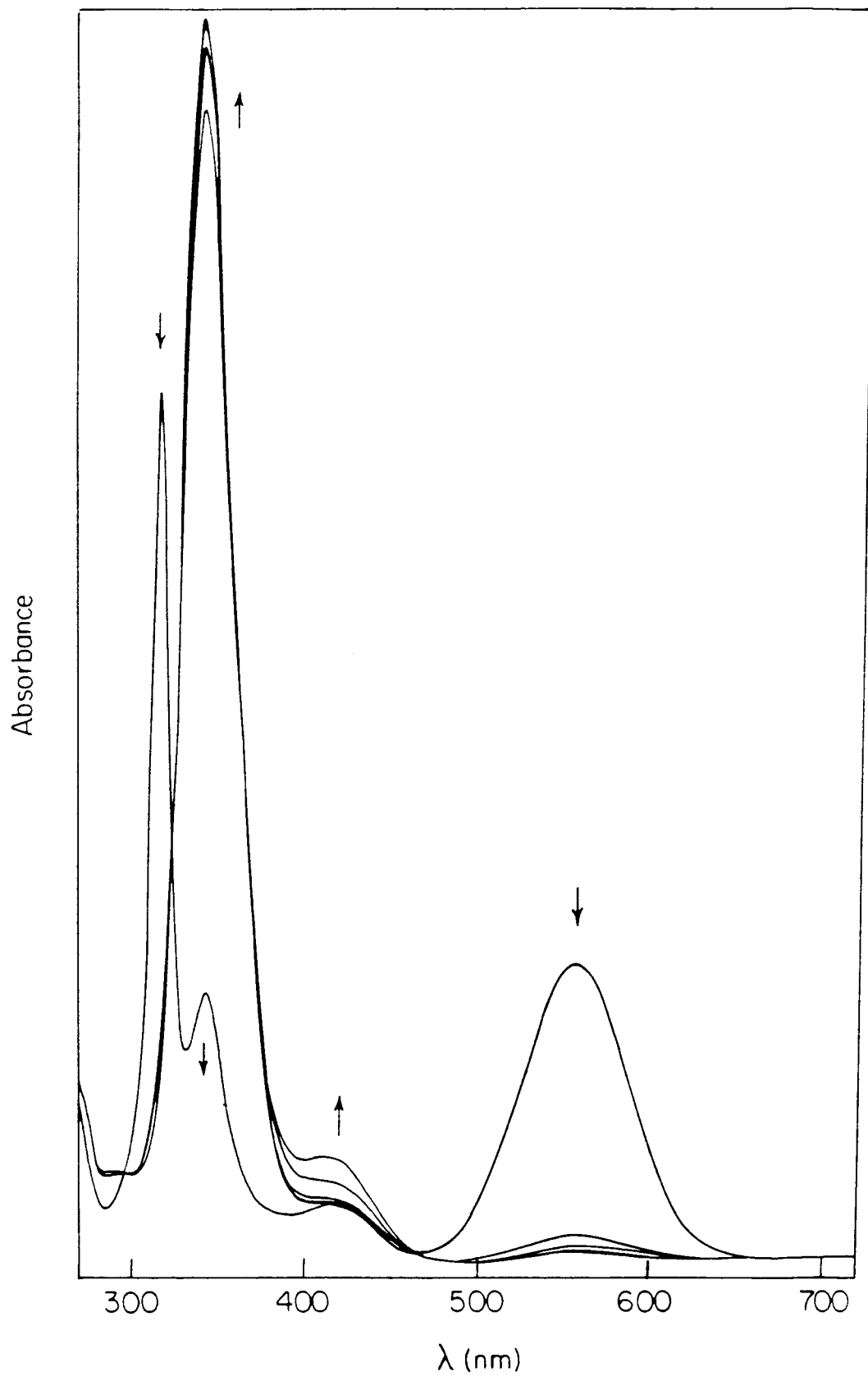
An alternative approach for inhibiting the back electron transfer is generation of a reduced acceptor, which is thermally unstable. The unimolecular decomposition of the reduced acceptor would compete with the bimolecular back-electron-transfer reaction.

Photolysis of the binuclear  $d^8$  complexes in the presence of alkyl halides yields products that can be rationalized to occur from the thermal decomposition of the reduced alkyl halide: an excited-state electron-transfer reaction followed by dissociative decomposition of the alkyl halide radical anion, and rapid scavenging of the fragmentation products.<sup>2,49,58,87,88</sup> The back-electron-transfer reaction is inhibited by the decomposition of the alkyl halide radical anion. Photolysis of  $\text{Ir}_2(\text{TMB})_4^{2+}$  in 1,2-dichloroethane cleanly generates  $\text{Ir}_2(\text{TMB})_4\text{Cl}_2^{2+}$  (Figure 4.25). The presumed organic product is ethylene. Similarly for  $\text{Rh}_2\text{b}_4^{2+}$ , a clean photoredox reaction is observed with 1,2-dichloroethane for  $\lambda_{\text{ex}} > 500 \text{ nm}$  (Figure 4.26). The inorganic product is  $\text{Rh}_2\text{b}_4\text{Cl}_2^{2+}$ .

**Figure 4.25.** Spectral changes upon irradiation of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in neat 1,2-dichloroethane,  $\lambda_{\text{ex}} > 604$  nm.



**Figure 4.26.** Spectral changes upon irradiation of  $[\text{Rh}_2\text{b}_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in neat 1,2-dichloroethane,  $\lambda_{\text{ex}} > 500 \text{ nm}$ .



An  $S_{RN}1$  mechanism (Figure 4.12) has been proposed to explain the photoredox reaction of alkyl halides with binuclear  $d^8$  complexes.<sup>49</sup> For  $[\text{Ir}(\mu\text{-pz})\text{COD}]_2$  and  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$ , whose excited-state reduction potentials are less than -1.5 V (SSCE), an outer-sphere electron-transfer reaction seems likely. Reduction potentials for alkyl halides of interest are generally more negative than -1.5 V (SSCE).<sup>59</sup> The quenching rate constant for the  $^3(d\sigma^*p\sigma)$  excited state of  $[\text{Ir}(\mu\text{-pz})\text{COD}]_2$  by 1,2-dichloroethane was found to agree with the rate expected for an outer-sphere electron-transfer reaction to an acceptor with a reduction potential of -2.0 V (SSCE).<sup>58</sup> However, an outer-sphere electron-transfer reaction seems unlikely for complexes with  $E^0(\text{M}_2^{2+}/^3\text{M}_2^*) > -1.0$  V (SSCE).

An alternate pathway to outer-sphere electron transfer, which yields similar photoredox products with alkyl halides, is excited-state atom transfer (Figure 4.13). Data obtained for  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$  indicate that alkyl and aryl halides react with the  $^3(d\sigma^*p\sigma)$  excited state via halogen-atom transfer,<sup>2</sup> the  $^3(d\sigma^*p\sigma)$  excited state abstracting the halide generating a  $d^8$ - $d^7$  monohalide species and an organic radical. An atom-transfer mechanism is favored for the photoredox reaction observed for  $\text{Ir}_2(\text{TMB})_4^{2+}$  and  $\text{Rh}_2\text{b}_4^{2+}$  with 1,2-dichloroethane.

Earlier work had suggested that the photoreaction of  $\text{Rh}_2\text{b}_4^{2+}$  with alkyl iodides proceeds by a radical chain mechanism initiated by an excited-state electron-transfer reaction.<sup>88</sup> An equivalent mechanism can be written with an atom-transfer initiation step. The photochemical reaction for  $\text{Rh}_2\text{b}_4^{2+}$ , unlike the reaction for  $\text{Ir}_2(\text{TMB})_4^{2+}$ , is complete in a matter of minutes (rather than hours). The relative rates of the photoredox reaction for  $\text{Rh}_2\text{b}_4^{2+}$  and  $\text{Ir}_2(\text{TMB})_4^{2+}$  can be understood from the varied lifetimes (and emission quantum yields) for the two complexes.<sup>89</sup>

The photoredox reaction for  $\text{Ir}_2(\text{TMB})_4^{2+}$  appears very clean for  $\lambda_{\text{ex}} > 500$  nm. As  $\lambda_{\text{ex}}$  is decreased to 400 nm, secondary photolysis occurs. The most likely reaction is photochemical degradation of  $\text{Ir}_2(\text{TMB})_4\text{Cl}_2^{2+}$ .<sup>4</sup> Higher-energy irradiation of  $\text{Rh}_2\text{b}_4^{2+}$



in 1,2-dichloroethane results in very complex behavior (Figure 4.27). Rapid disappearance of  $\text{Rh}_2\text{b}_4^{2+}$  with growth of  $\text{Rh}_2\text{b}_4\text{Cl}_2^{2+}$  is observed with slower loss of  $\text{Rh}_2\text{b}_4\text{Cl}_2^{2+}$  and appearance of bands attributable to oligomeric rhodium species.<sup>90</sup> The final solution is brown-orange. Removal of the solvent yields dark-purple material.

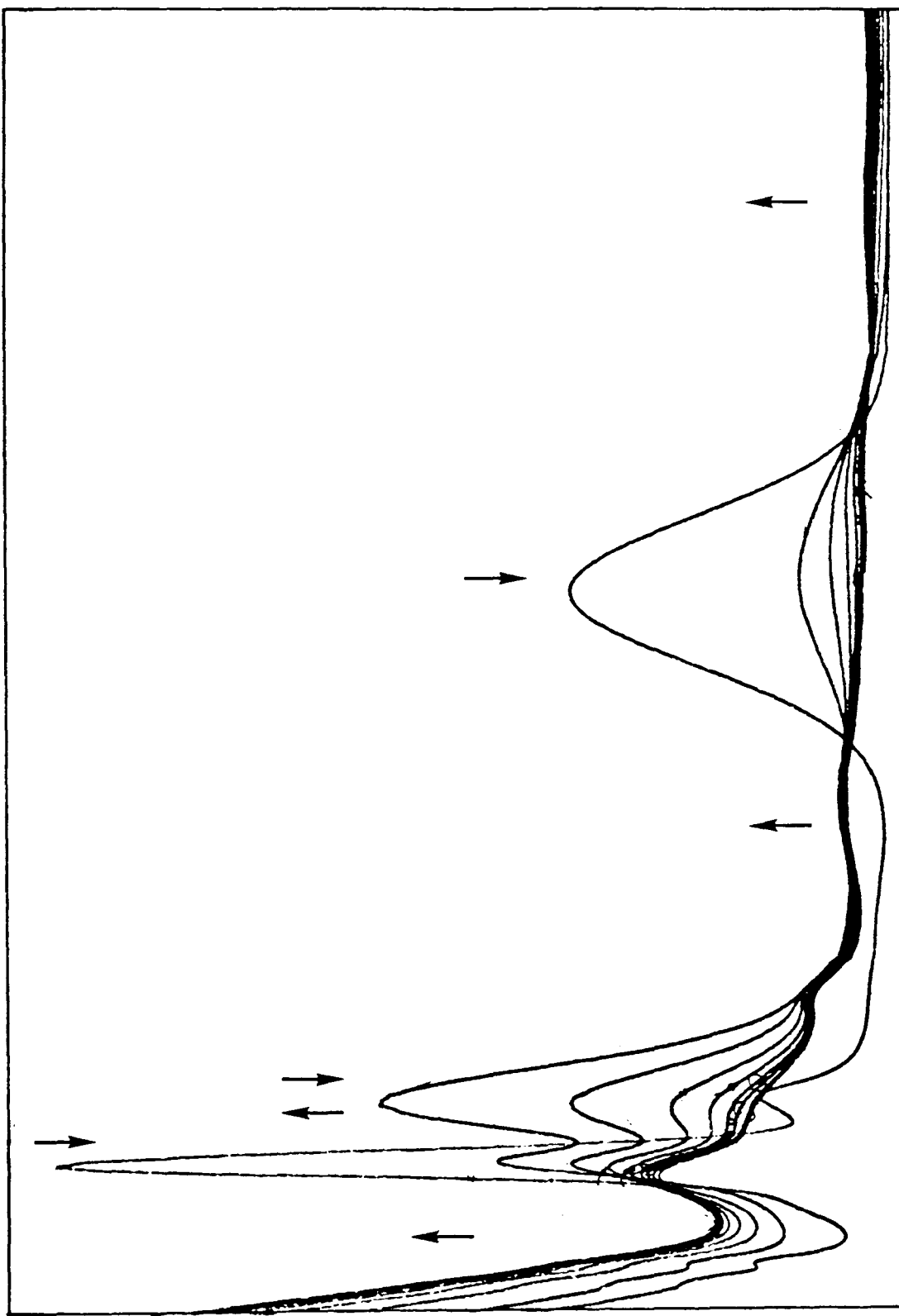
One concern for the photoreaction of  $\text{Rh}_2\text{b}_4^{2+}$  is the possibility of oligomer reaction rather than dimer reaction. The photochemical generation of  $\text{H}_2$  from  $\text{HCl}$  solutions of  $\text{Rh}_2\text{b}_4^{2+}$  is thought to result from photochemical cleavage of  $[\text{Rh}_2\text{b}_4\text{Cl}]_2^{4+}$ .<sup>91</sup> Generation of mixed valence rhodium complexes followed by thermal electron transfer or atom transfer might explain the observed alkyl-halide photoreduction.

Atom transfer to the  $^3(\text{d}\sigma^*\text{p}\sigma)$  excited state is the favored reaction mechanism for the alkyl-halide photoreduction reaction. While outer-sphere electron transfer appears plausible for a metal complex with  $E^0(\text{M}_2^{+}/^3\text{M}_2^*) < -1.5 \text{ V (SSCE)}$ , it seems unlikely for complexes with  $E^0(\text{M}_2^{+}/^3\text{M}_2^*) > -1.0 \text{ V (SSCE)}$ . One question remains to be answered. Why would systems, whose excited state reduction potentials are much less than  $-1.5 \text{ V (SSCE)}$ , react by an atom-transfer mechanism rather than by an outer-sphere electron-transfer pathway? One explanation may be obtained from the study of the electron-transfer quenching of binuclear  $\text{d}^8$  complexes. The oxidative quenching of the  $^3(\text{d}\sigma^*\text{p}\sigma)$  excited state is found to be quite nonadiabatic. In order to overcome the large nonadiabaticity, the metal complex coordinates the alkyl halide. Inner-sphere electron transfer coupled with C-X bond cleavage results in net halide-atom transfer to the metal complex.

### Hydrogen-Atom Transfer

$\text{Ir}_2(\text{TMB})_4^{2+}$  ( $\text{Ir}_2$ ) is found to react photochemically with 9,10-dihydroanthracene and 1,4-cyclohexadiene. A representative electronic absorption spectrum of the photochemical reaction between  $\text{Ir}_2(\text{TMB})_4^{2+}$  and 9,10-

**Figure 4.27.** Spectral changes upon irradiation of  $[\text{Rh}_2\text{b}_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in neat 1,2-dichloroethane,  $\lambda_{\text{ex}} > 400 \text{ nm}$ .



dihydroanthracene is shown in Figure 4.28. Rapid disappearance of the  $^1(d\sigma^*p\sigma)$  band with growth of higher-energy transitions is observed, indicating the formation of anthracene. Growth of resonances assignable to anthracene and Ir-H are observed in the  $^1\text{H}$  NMR (Figure 4.29). Isolation of the inorganic product from the reaction mixture yields a material whose electronic spectrum suggests the formation of a transdihydride  $d^7$ - $d^7$  adduct,  $\text{Ir}_2(\text{TMB})_4\text{H}_2^{2+}$  ( $\text{Ir}_2\text{H}_2$ ). A moderate Ir-H stretch ( $1940\text{ cm}^{-1}$ ) is observed in the IR, indicative of a terminal metal hydride. No appreciable thermal reaction with 9,10-dihydroanthracene is observed at room temperature; however, a reaction is observed at elevated temperatures. Anthracene and Ir-H resonances are observed in the  $^1\text{H}$  NMR.

For 1,4-cyclohexadiene, no thermal reaction is observed, even at elevated temperatures. Photolysis of  $\text{Ir}_2(\text{TMB})_4^{2+}$  and 1,4-cyclohexadiene results in rapid disappearance of the  $d^8$ - $d^8$  metal complex. With no complication because of overlapping product spectra for the reaction with 1,4-cyclohexadiene, clean isosbestic conversion to the product spectrum is observed (Figure 4.30). Growth of resonances attributable to benzene and Ir-H are observed in the  $^1\text{H}$  NMR (Figure 4.31).

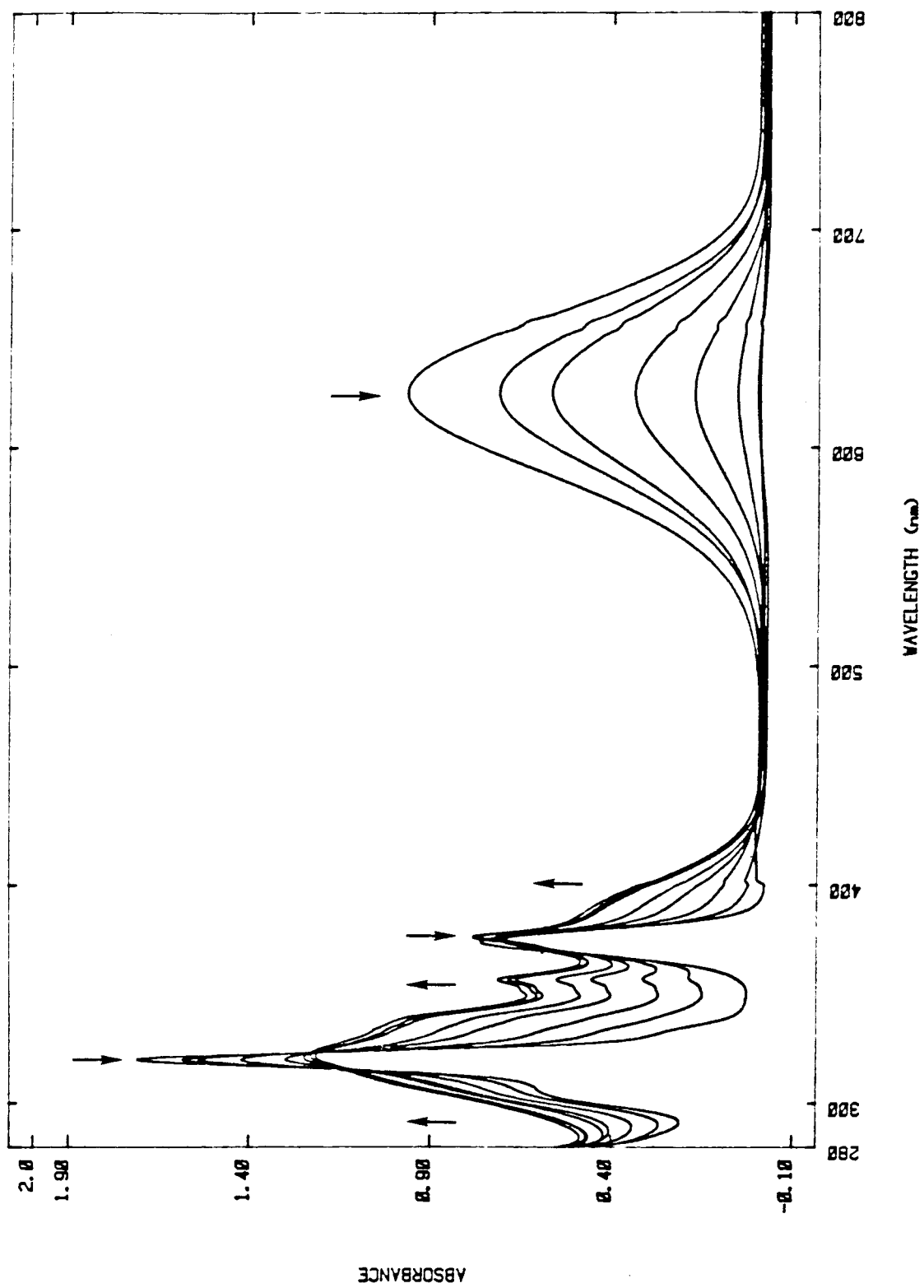
No dark reaction is observed after partial photolysis for either 9,10-dihydroanthracene or 1,4-cyclohexadiene (no indication of a photoinitiated process).

The net photoreaction observed for  $\text{Ir}_2(\text{TMB})_4^{2+}$  with a number of hydrocarbon substrates is hydrogen-atom transfer to the metal center.

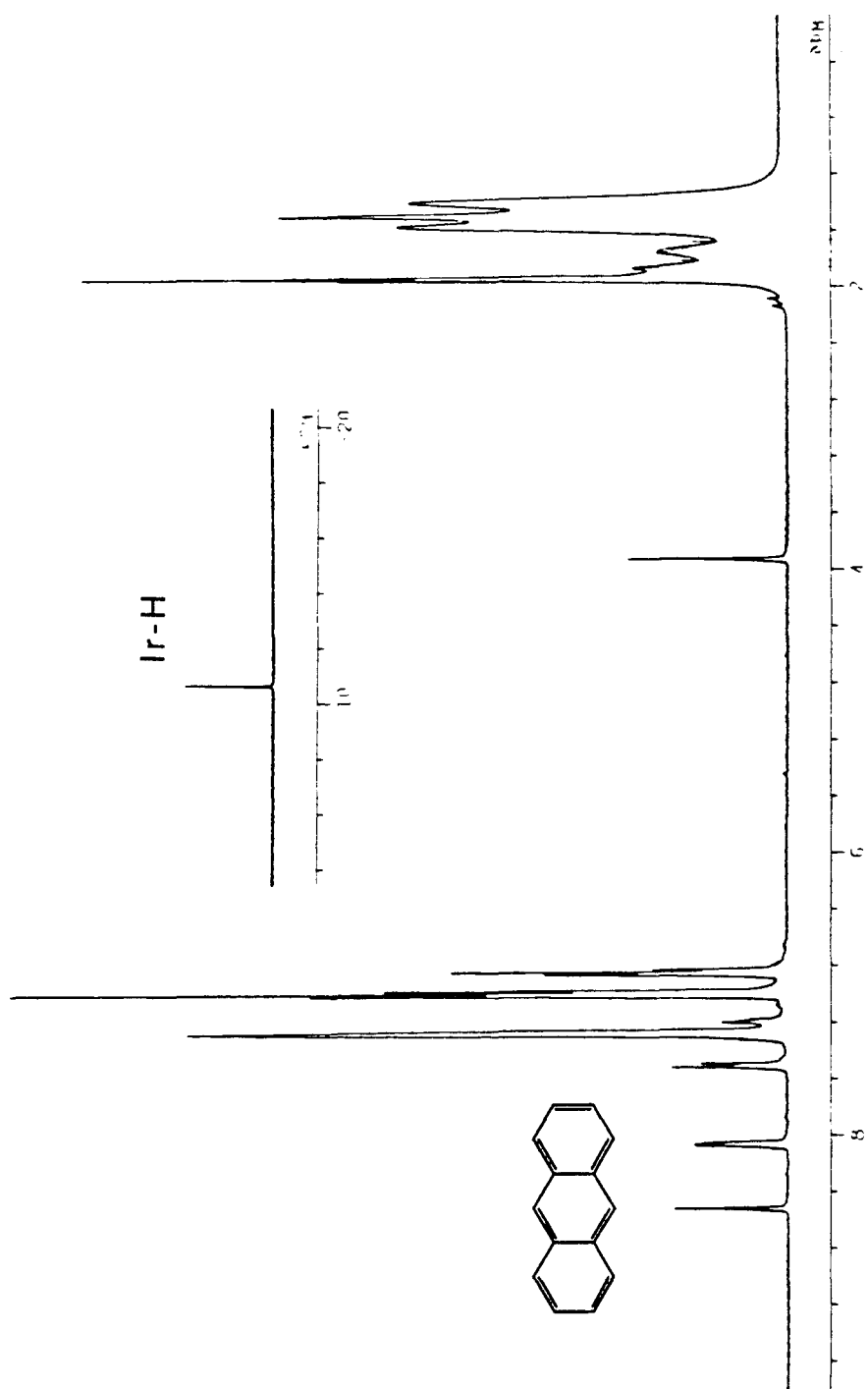


No evidence of organic radical-coupling products ( $(\text{RH})_2$ ) or alkyl-hydride metal complexes ( $\text{RH-Ir-Ir-H}$ ) is observed. The  $\text{Ir}_2\text{H}_2$  complex has been isolated and characterized (see Chapter 5).

**Figure 4.28.** Spectral changes upon irradiation of an acetonitrile solution of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  and 9,10-dihydroanthracene.

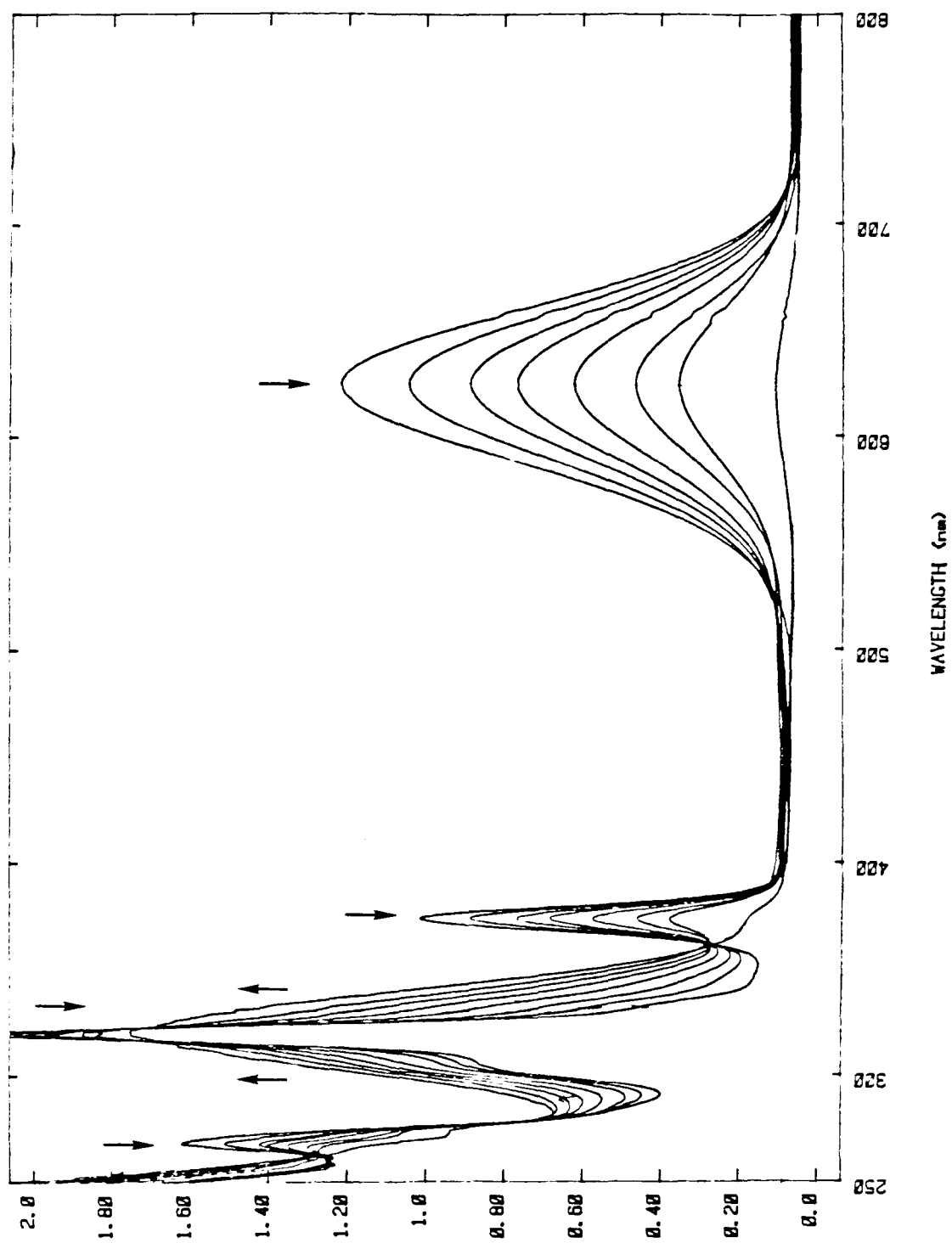


**Figure 4.29.** NMR spectrum of an acetonitrile- $d_3$  photolysis solution of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  and 9,10-dihydroanthracene.





**Figure 4.30.** Spectral changes upon irradiation of an acetonitrile solution of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  and 1,4-cyclohexadiene.



**Figure 4.31.** NMR spectrum of an acetonitrile- $d_3$  photolysis solution of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  and 1,4-cyclohexadiene.

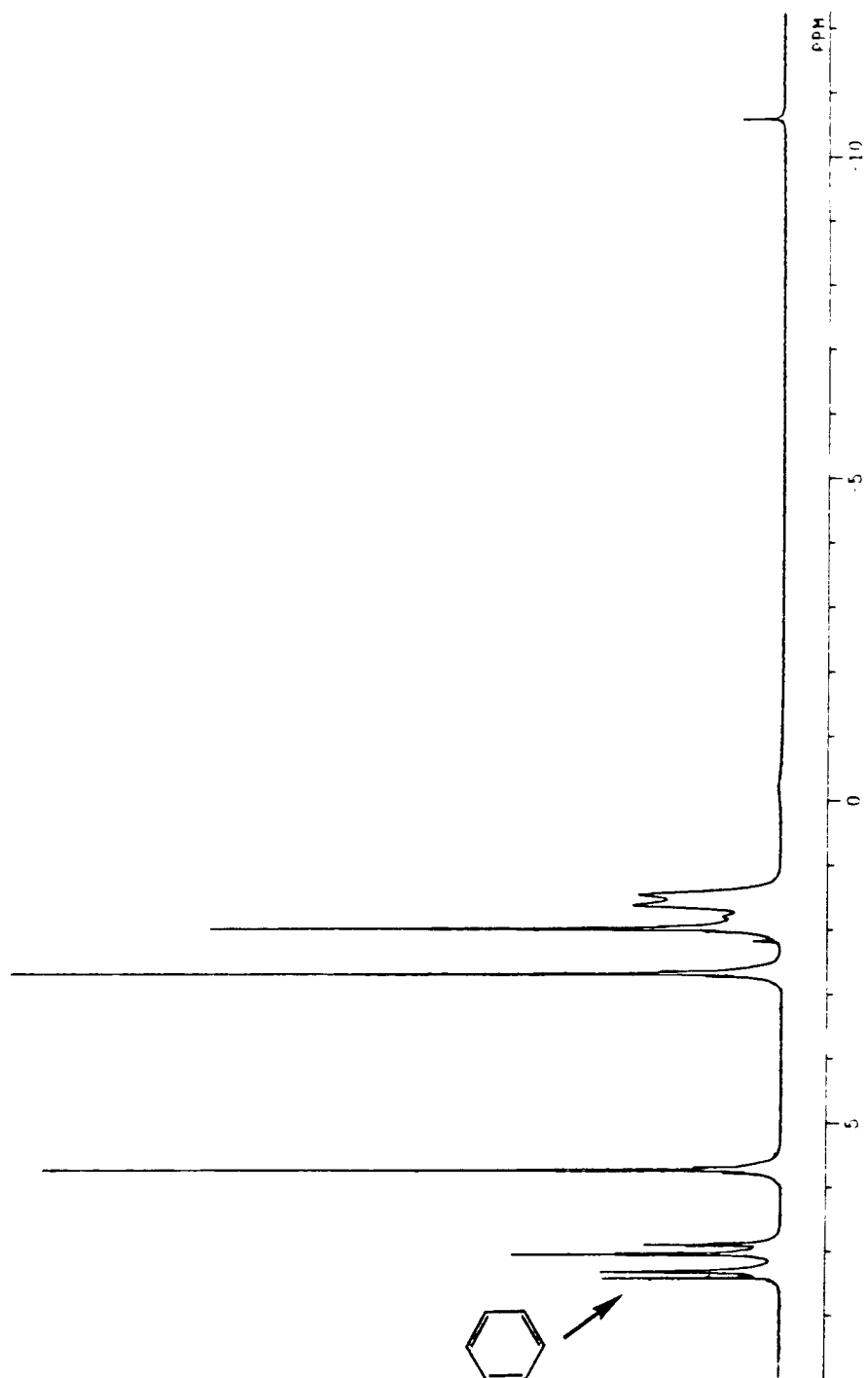


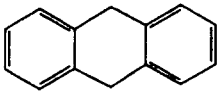
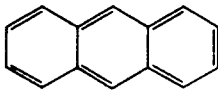



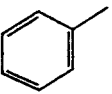
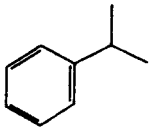
Table 4.10 lists a number of hydrocarbon substrates and their observed photochemical reactivity with  $\text{Ir}_2$ . The conclusions are based on conversion to  $\text{Ir}_2\text{H}_2$ . While conversion to  $\text{Ir}_2\text{H}_2$  is not observed for all substrates, all but triphenylmethane quench the  $^3(\text{d}\sigma^*\text{p}\sigma)$  excited state (*vide infra*). Photochemical conversion of the organic substrates may occur with no build-up of inorganic product. For  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$  ( $\text{Pt}_2$ ), excited-state quenching by toluene and cumene is observed; steady-state photolysis ( $\lambda_{\text{ex}}$  370 nm) produces  $\text{H}_2$  and organic radical products (for toluene, bibenzyl; for cumene, 2,3-dimethyl-2,3-diphenylbutane) with no formation of  $\text{Pt}_2\text{H}_2$ .<sup>92</sup>

To better understand the hydrogen-atom transfer reactivity of the  $^3(\text{d}\sigma^*\text{p}\sigma)$  excited state of  $\text{Ir}_2(\text{TMB})_4^{2+}$ , a study of the hydrogen-transfer quenching by a series of hydrogen-atom donors with variable homolytic C-H bond strengths was undertaken. All systems were found to obey Stern-Volmer kinetics over the quencher concentration range studied. A plot of the data for the series of hydrogen-atom donors is shown in Figure 4.32. The quenching rate constants, listed in Table 4.11, were obtained from the slopes of linear-least-squares fits to the data.

Results for  $^3\text{Pt}_2^*$  have been interpreted as supporting a hydrogen-abstraction pathway in which the H atom transfers via a linear Pt-H-E transition state with negligible charge transfer.<sup>56</sup> The reactivity of  $^3\text{Pt}_2^*$  has been compared to that of the  $n\pi^*$  excited states of ketones with similar triplet energies.<sup>2</sup> For  $^3\text{Ir}_2^*$ , a similar interpretation is possible; however, the analogy to the  $n\pi^*$  excited states of ketones cannot be made in that the triplet energy for this complex is 30 kcal/mol compared to 60 kcal/mol for  $^3\text{Pt}_2^*$ .

If the photoreaction is truly an atom-transfer process, the observed rate in the absence of steric effects should track the homolytic C-H bond energies of the substrates. In a series of quenchers, the observed rate constant decreases with increasing C-H bond energy (Figure 4.33).<sup>27</sup> However, this trend is not general. For cyclohexene, whose

**Table 4.10.** Observed photochemical reactivity of hydrocarbons with  $\text{Ir}_2(\text{TMB})_4^{2+}$ .

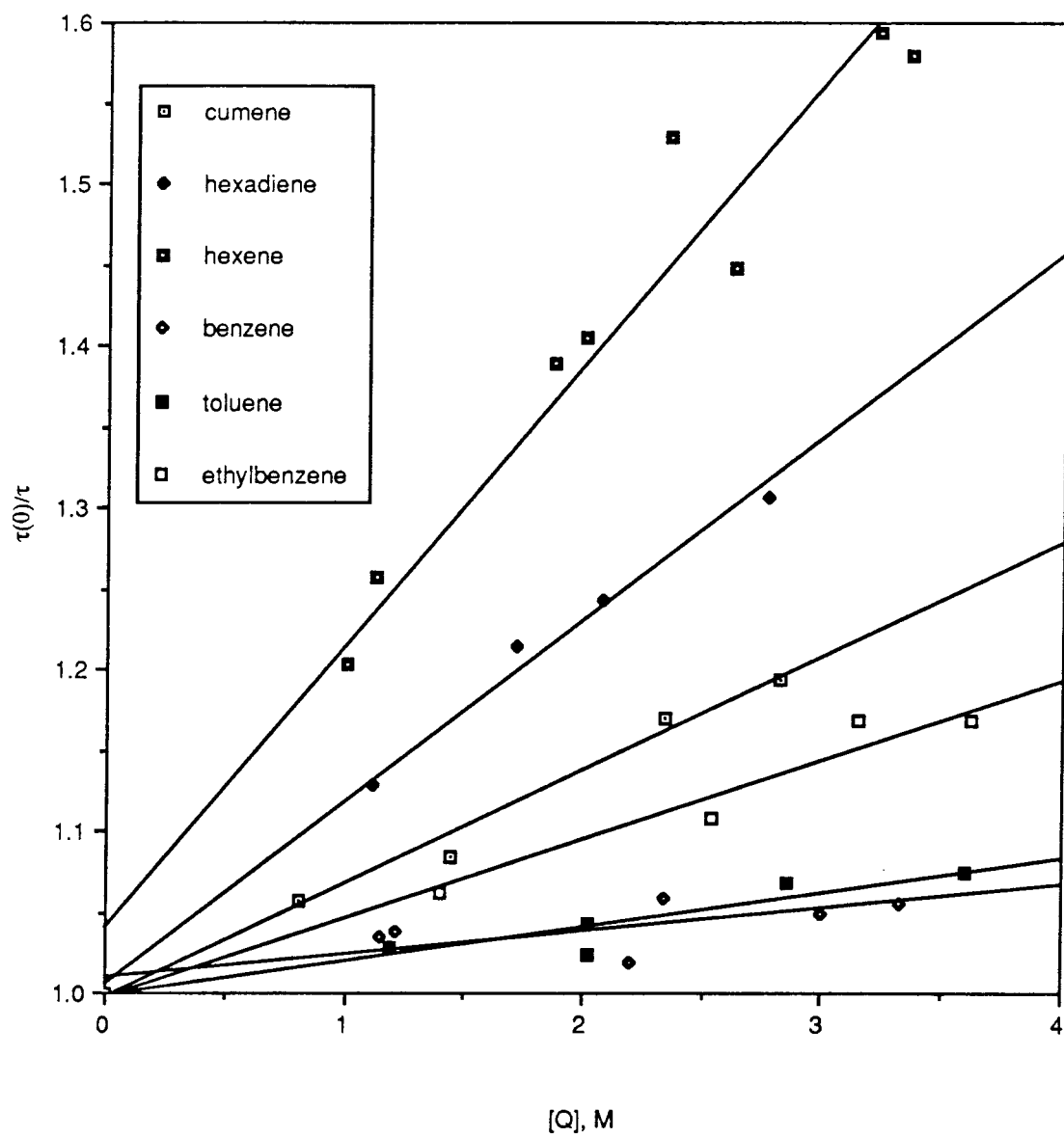
Substrate	D(C-H), kcal/mol <sup>a</sup>	Products
	77	$\text{Ir}_2\text{H}_2$ , 
	73	$\text{Ir}_2\text{H}_2$ , 
	82	$\text{Ir}_2\text{H}_2$ <sup>b</sup>
	88	NR <sup>c</sup>
	84.4	NR
$\left(\text{C}_6\text{H}_5\right)_3\text{CH}$	~75	NR

a. C-H bond dissociation energy, Reference 27.

b. Sluggish formation of  $\text{Ir}_2\text{H}_2$ .



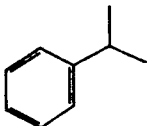
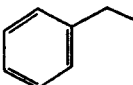
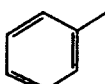


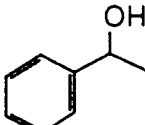
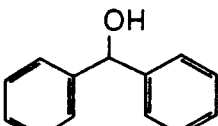
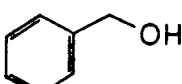
c. No reaction, no observed formation of  $\text{Ir}_2\text{H}_2$ .

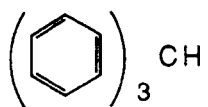
**Figure 4.32.** Stern-Volmer plot for hydrogen atom-transfer quenching of the  $^3(\text{d}\sigma^*\text{p}\sigma)$  excited state of  $\text{Ir}_2(\text{TMB})_4^{2+}$  for a series of hydrocarbons.





**Table 4.11.** Stern-Volmer quenching rate constants for  $\text{Ir}_2(\text{TMB})_4^{2+}$  and  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$ .

Substrate	BDE (kcal/mol) <sup>a</sup>	$^3\text{Ir}_2^*$ ( $\text{M}^{-1}\text{s}^{-1}$ )	$^3\text{Pt}_2^*$ ( $\text{M}^{-1}\text{s}^{-1}$ )
	82 <sup>b</sup>	$7.9 \times 10^5$ ( $k_H/k_D > 3$ )	$1.2 \times 10^6$
	73 <sup>b</sup>	$5.2 \times 10^5$	$8.2 \times 10^6$
	84.4 <sup>b</sup>	$3.1 \times 10^5$	$10^4$
	85.5 <sup>b</sup>	$2.3 \times 10^5$	
	88.0 <sup>b</sup>	$9.7 \times 10^4$	$10^4$
		$6.5 \times 10^4$	
	91.0 <sup>b</sup>		$5 \times 10^3$
	88		$1.9 \times 10^6$ ( $k_H/k_D = 4.6$ )
	82		$2 \times 10^5$
	83		$3 \times 10^6$
$\text{Et}_3\text{SiH}$	c	d	$2.0 \times 10^4$
$\text{Bu}_3\text{SnH}$	c	d	$1.2 \times 10^7$



~75

< 4 x 10<sup>5</sup>< 2 x 10<sup>4</sup>

a. C-H bond dissociation energy.

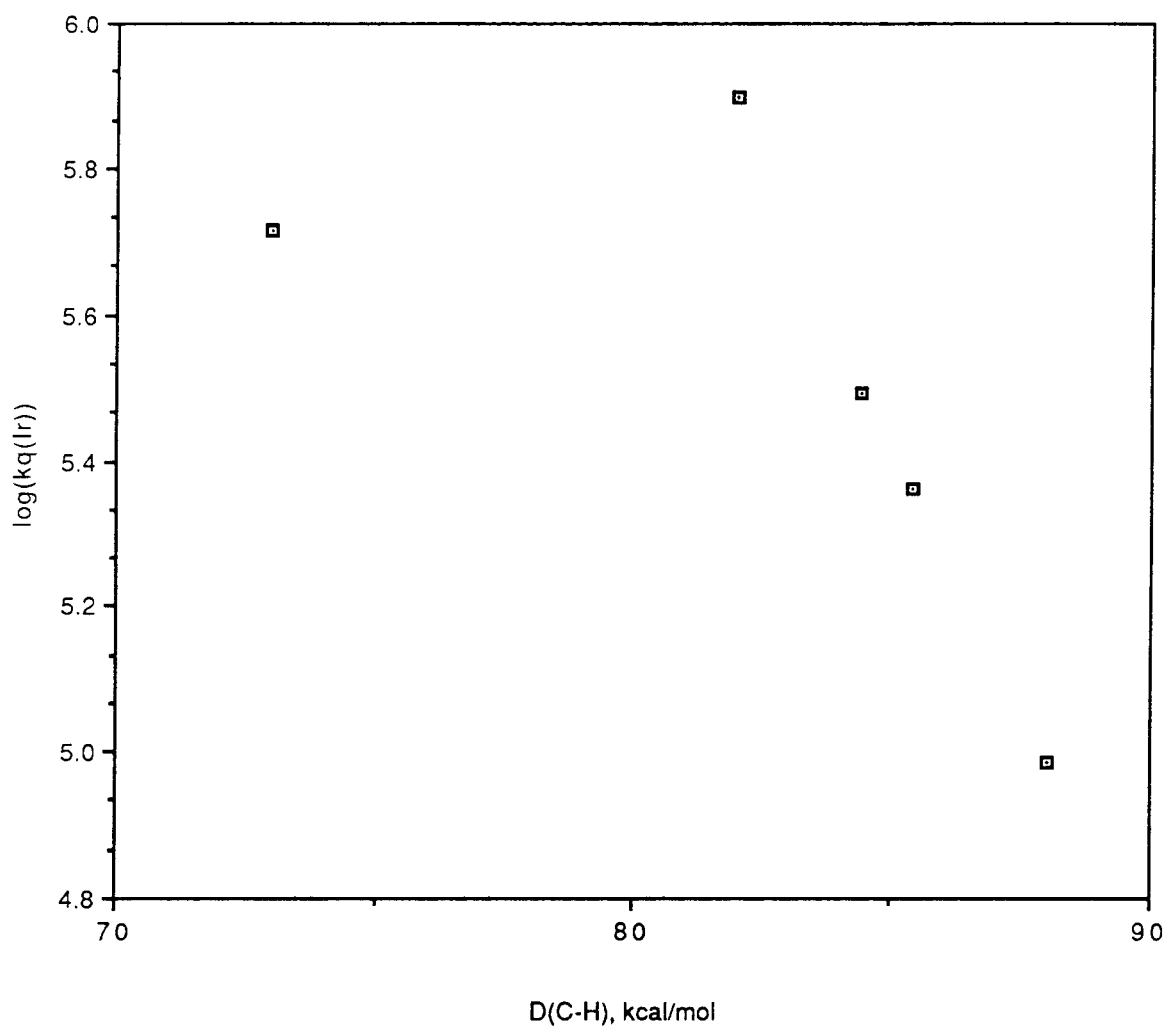
b. Reference 27.

c. The gas phase (298 K) E-H bond energies of (CH<sub>3</sub>)<sub>3</sub>EH species have been reported:

90 (Si) and 74 kcal/mol (Sn). Jackson, R.A. *J. Organomet. Chem.* **1979**, 166, 17-19.

d. Reacts upon mixing.

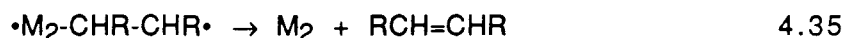
**Figure 4.33.** Plot of  $\log(k_q)$  versus the C-H bond dissociation energy ( $D(\text{C-H})$ ) of the hydrogen atom-transfer quenching of  $^3(\text{Ir}_2(\text{TMB})_4^{2+})^*$  by the hydrocarbon substrates.



homolytic C-H bond energy is ~10 kcal/mol greater than that of 1,4-cyclohexadiene, a larger rate constant is observed. However, steric arguments may be sufficient to reconcile the order of rates. Recent work has found the structure of 1,4-cyclohexadiene to be almost planar (Figure 4.34).<sup>93</sup> The structure of cyclohexene is a skewed boat (Figure 4.34).<sup>94</sup> For a linear M-H-C transition state, a much more unfavorable steric interaction is expected for the planar 1,4-cyclohexadiene. The angle defined by the M-H-C vector and the plane containing the allyl unit is much less for 1,4-cyclohexadiene than for cyclohexene.

The dihydride of  $\text{Ir}_2(\text{TMB})_4^{2+}$  ( $\text{Ir}_2\text{H}_2$ ) has been characterized as the primary photoproduct. For some substrates (toluene and cumene), very little  $\text{Ir}_2\text{H}_2$  product is observed. For cyclohexene,  $\text{Ir}_2\text{H}_2$  appears slowly; either efficient  $\text{Ir}_2\text{H}_2$  formation requires a rapid in-cage reaction of  $\text{Ir}_2\text{H}$  with the organic radical, or the factors that govern the second H atom-transfer reaction are unfavorable with either the substrates or the substrate radicals present.

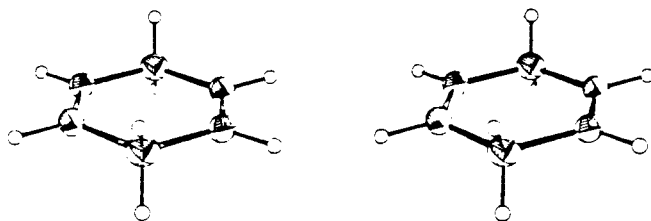
An alternate explanation for the variation in quenching rates and chemical conversions is an alternative quenching mechanism. Both energy-transfer quenching and electron-transfer quenching can be ruled out.<sup>95,96</sup> The energies required for these processes are too great. An alternate mechanism, suggested for the photoreaction of  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$  with simple alkenes, is an inner-sphere process that involves formation of a diradical intermediate.<sup>97</sup>



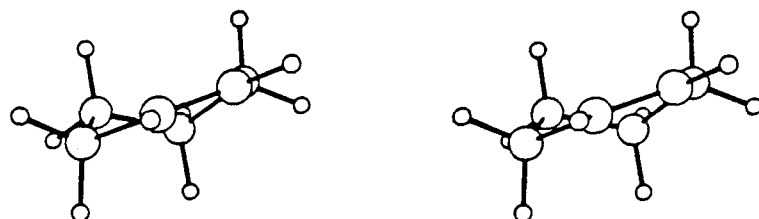
The second step in this mechanism involves the collapse of this intermediate back to the alkene and the ground-state metal complex. The rapid formation of  $\text{Ir}_2\text{H}_2$  with 1,4-

**Figure 4.34.** (a) Stereoview of the X-ray structure of 1,4-cyclohexadiene. (b) Structure of cyclohexene; (i) stereo representation, (ii) conventional symbolism.

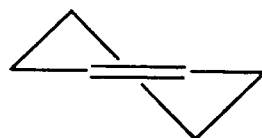
(a)



(b)



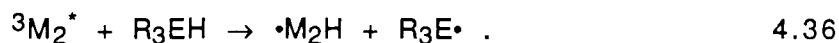
(i)



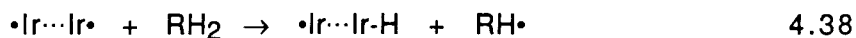
(ii)

cyclohexadiene and the large kinetic isotope effect observed for cyclohexene do not support such a mechanism.

From these data, the primary photoprocess from the  $^3(d\sigma^*p\sigma)$  excited state is hydrogen-atom transfer:



Observation of a large kinetic isotope effect is supportive of an atom-transfer process for  $^3Ir_2^*$ . The monohydride species for  $Ir_2(TMB)_4^{2+}$  has not been identified. The proposed pathway for formation of the observed products is



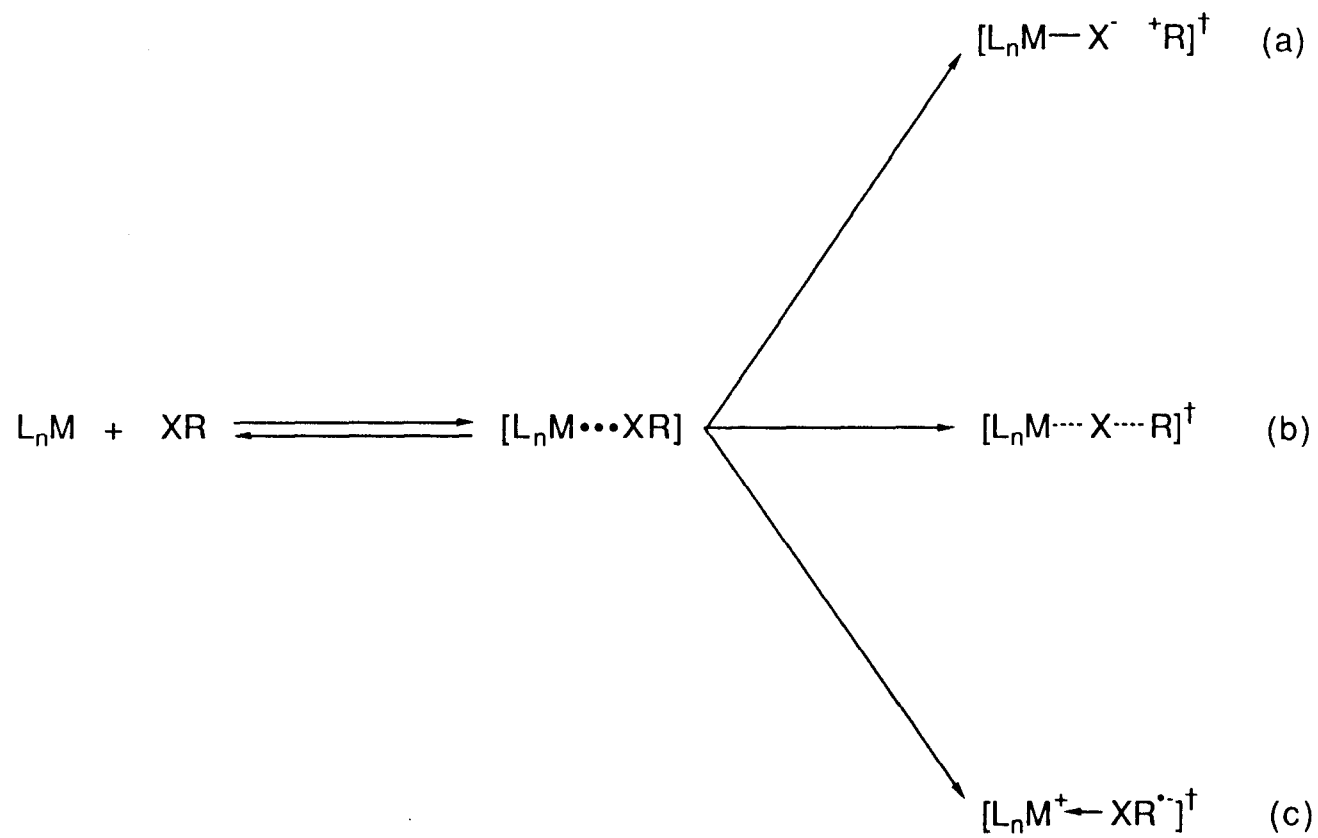
The  $^3(d\sigma^*p\sigma)$  excited state is written as  $\cdot Ir \cdots Ir \cdot$  to emphasize the diradical-like structure.

The mechanism of radical atom abstraction has been extensively studied for many years and numerous reviews have appeared.<sup>55</sup> Three limiting atom abstraction mechanisms are defined by the timing of the various steps: M-Y bond making, R-Y bond cleavage, and electron transfer from M to Y. The three limiting mechanisms are shown in Figure 4.35. The timing of the electron-transfer event with respect to R-Y bond cleavage and M-Y bond formation becomes earlier in going from a to c.

Halogen-atom abstraction from alkyl halides has been extensively studied by Brown and coworkers.<sup>98,99</sup> They showed that the rate constant for



**Figure 4.35.** Possible limiting mechanism of atom transfer. (a) Heterolytic R-X bond cleavage and M-X bond formation followed by electron transfer. (b) Synchronous electron transfer, R-X bond cleavage, and M-X bond formation. (c) Electron transfer followed by R-X<sup>-</sup> heterolytic cleavage and M-X bond formation.



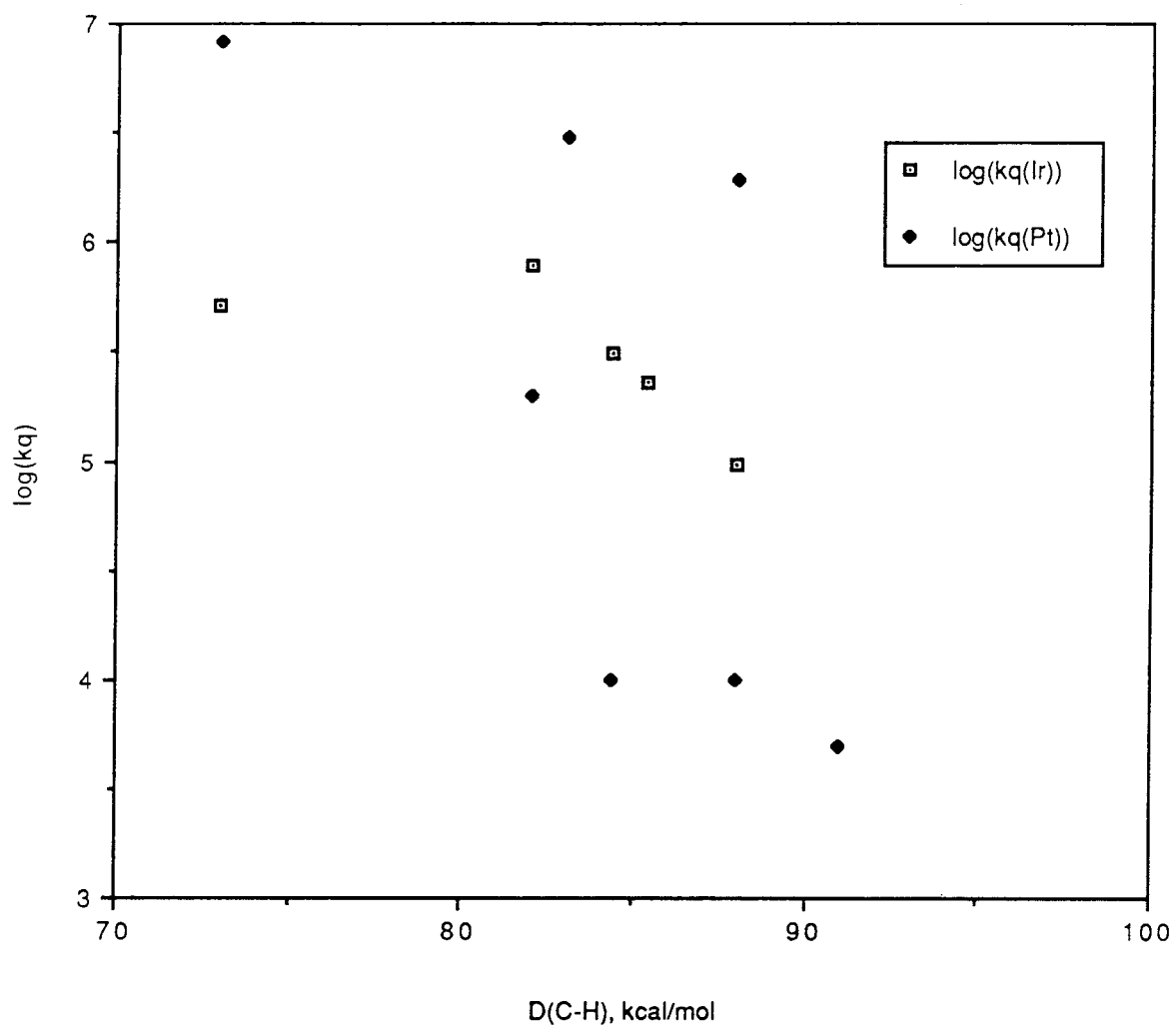


could be fit to the Marcus/Agmon-Levine equation, correlating the  $E_{1/2}$  value for the halides and the rate constants for atom abstraction. The correlation between the rate constants and the C-X bond dissociation energy was not very good. That a Marcus/Agmon-Levine correlation exists between  $E_{1/2}$  and the atom abstraction rates was interpreted as being consistent with some electron transfer, from the metal to the halogen-atom donor, in the transition state. Brown and coworkers were able to narrow down the possible pathways to b or c of Figure 4.35.

A limited correlation of the quenching rate constant with the C-H bond dissociation energy is observed for  $\text{Ir}_2(\text{TMB})_4^{2+}$  (Figure 4.33). Deviation from this correlation can be understood to be due to unfavorable steric interactions. If an atom-transfer mechanism with little to no charge development in the transition state is the preferred mechanism, correlation of the quenching rate constants for  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$  and the C-H bond dissociation energies is expected. The kinetic isotope effect observed for  $\text{Bu}_3\text{SnH}$  is interpreted to be in accord with hydrogen-atom transfer involving a linear Pt-H-Sn transition state with negligible charge transfer.<sup>56</sup>

A plot of the  $\log(k_q)$  versus the C-H bond dissociation energy  $D(\text{C-H})$  is given in Figure 4.36. While a limited correlation is observed for  $\text{Ir}_2(\text{TMB})_4^{2+}$ , a very poor correlation is found for  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$ . (A general downward drift in  $k_q$  is observed for increasing C-H bond energy.) This result may suggest an alternate quenching pathway for  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$ . With the greater redox energy, electron transfer in the transition state may be important for the reaction of  $^3\text{Pt}_2^*$ . An alternative explanation for the observed behavior is a variation in the steric factors for the series of hydrogen atom donors.

**Figure 4.36.** Plot of  $\log(k_q)$  versus the C-H bond dissociation energy ( $D(\text{C-H})$ ) of the hydrogen atom-transfer quenching of  $^3\text{M}_2^*$  by the hydrocarbon substrates,  $\text{Ir}_2(\text{TMB})_4^{2+}$  open squares and  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$  solid diamonds.



In comparing rate constants for  $^3\text{Ir}_2^*$  and  $^3\text{Pt}_2^*$ , some surprising trends emerge. For small substrates, such as 1,4-cyclohexadiene, the rate observed for  $^3\text{Pt}_2^*$  is greater. This is expected and is due to the more energetic excited state. Assuming a  $\text{M}_2\text{-H}$  dissociation energy of 60 kcal/mol,<sup>100</sup> for both Ir and Pt,<sup>101</sup> the driving force for hydrogen-atom transfer is estimated to be roughly  $(60 - D_{\text{EH}} + E_{\text{T}})$  47 kcal/mol for  $\text{Pt}_2$  and 17 kcal/mol for  $\text{Ir}_2$ . For cumene and possibly toluene, a larger rate is observed for  $^3\text{Ir}_2^*$ . These larger rates might be understood from the greater steric demands about the open axial site in  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$  in comparison to  $\text{Ir}_2(\text{TMB})_4^{2+}$ .

For reactions where most of the reorganization comes from the bonds being broken and formed, rather than from all the other coordinates, the free energy of activation is controlled by the work terms for the process ( $w_r$ ,  $w_p$ ), and the steric ( $S^r$ ,  $S^p$ ) and statistical factors ( $s^r$ ,  $s^p$ ), which describe the orientation of the reactants and products.<sup>102,103</sup>

$$\Delta G^* \cong W^r \quad (-\Delta G^{0'} \geq \lambda) \quad 4.43$$

$$\Delta G^* \cong \Delta G^{0'} + W^p \quad (\Delta G^{0'} \geq \lambda) \quad 4.44$$

$$W^r = w^r - RT \ln(S^r s^r)$$

$$W^p = w^p - RT \ln(S^p s^p)$$

$S^{r/p}$ ,  $s^{r/p}$  - steric and statistical factors that describe the orientation of the reactants and products, respectively.

$w^{r/p}$  - Coulombic and other work terms of the reactants and the products, respectively.

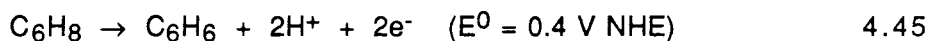
The poor correlation of the quenching rate of  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$  with the C-H bond dissociation energy can be understood as the result of variation in the steric factors of the reaction with variation of the substrate. Also, preassociation or favorable docking of the alcohol via hydrogen-bonding interactions with the  $\text{P}_2\text{O}_5\text{H}_2^{2-}$  ligand may enhance the

hydrogen-atom transfer process and may explain the increased quenching rate observed for alcohols in comparison to hydrocarbons with similar C-H bond dissociation energies. For  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$ , the greater steric demands about the open axial site and the possible substrate and/or solvent hydrogen-bonding interactions with the pyrophosphate ligand increase the sensitivity of the reaction to variations in the substrate. For  $\text{Ir}_2(\text{TMB})_4^{2+}$ , no preassociation or docking of the substrate is possible, and steric interactions alone discriminate between substrates.

The  $^3(\text{d}\sigma^*\text{p}\sigma)$  excited state of the  $\text{d}^8\text{-d}^8$  complexes has been shown to be involved in the photochemical hydrogen atom-transfer reaction. The atom-transfer reactivity of this state is attributed to the presence of a hole in the  $\text{d}\sigma^*$  orbital, analogous to the  $^3\text{n}\pi^*$  state of organic ketones. Interaction of the oxidizing hole with the electron pair of the C-H bond is the presumed pathway.

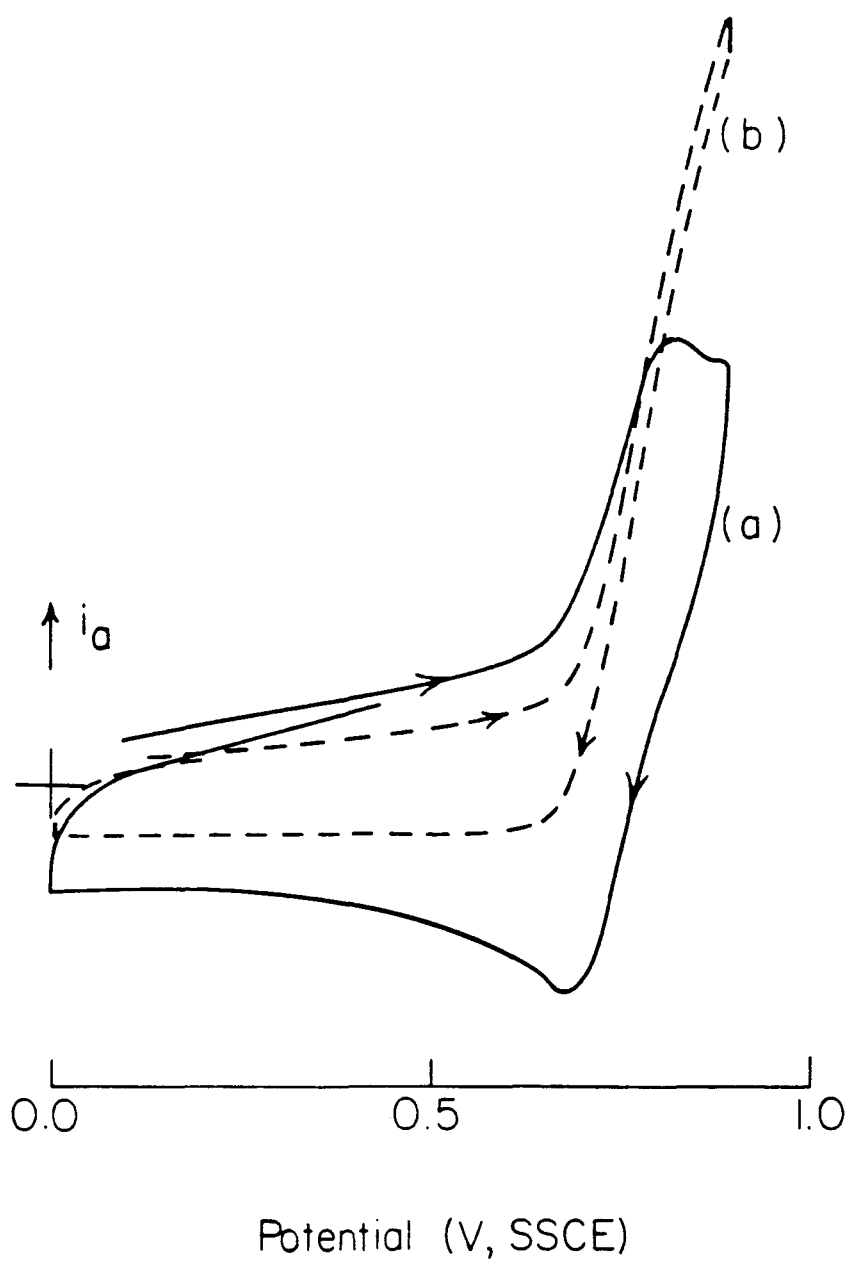
## Electrocatalysis

If production of an oxidizing hole in the  $\text{d}\sigma^*$  orbital is the important factor in the photochemical hydrogen-atom transfer reaction, then electrochemical generation of such a hole should produce a highly reactive intermediate that would mimic the initial step in the  $^3(\text{d}\sigma^*\text{p}\sigma)$  photoreaction. Several of the binuclear  $\text{d}^8$  complexes undergo reversible one-electron oxidations in noncoordinating solvents.<sup>43,104,105</sup> The complex  $\text{Rh}_2(\text{TMB})_4^{2+}$  possesses a quasi-reversible, one-electron oxidation at 0.74 V (SSCE). Electrochemical oxidation of  $\text{Rh}_2(\text{TMB})_4^{2+}$  in the presence of 1,4-cyclohexadiene exhibits an enhanced anodic current with loss of the cathodic wave, behavior indicative of an electrocatalytic process (Figure 4.37).<sup>106</sup> Bulk electrolysis of  $\text{Rh}_2(\text{TMB})_4^{2+}$  in an excess of 1,4-cyclohexadiene results in the formation of benzene and two protons (Equation 4.45).



**Figure 4.37.** Cyclic voltammograms of a dichloromethane solution (0.1 M TBAPF<sub>6</sub>) of (a) [Rh<sub>2</sub>(TMB)<sub>4</sub>](PF<sub>6</sub>)<sub>2</sub> (1 mM,  $i_a = 5 \mu\text{A}/\text{cm}$ ) with (b) a twentyfold excess of 1,4-cyclohexadiene ( $i_a = 10 \mu\text{A}/\text{cm}$ ).





The maximum number of turnovers observed is limited by the amount of substrate relative to solvent. The complex  $\text{Rh}_2(\text{TMB})_4^{2+}$  is slowly lost because of a competitive reaction with  $\text{CH}_2\text{Cl}_2$  to produce  $\text{Rh}_2(\text{TMB})_4\text{Cl}_2^{2+}$ .

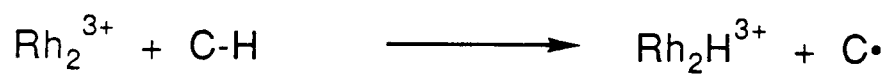
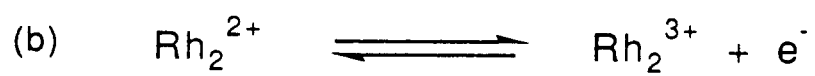
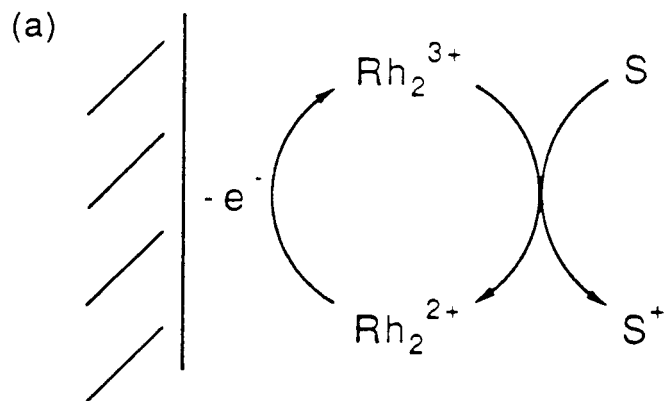
The pathway proposed to explain the observed oxidation of 1,4-cyclohexadiene is given in Figure 4.38. Oxidation  $\text{Rh}_2(\text{TMB})_4^{2+}$  generates the reactive  $d^8-d^7$  complex,  $\text{Rh}_2(\text{TMB})_4^{3+}$ . The metal complex abstracts a hydrogen atom from the organic substrate, generating a monohydride species and an organic radical. It is known that hydride complexes of  $\text{Rh}_2(\text{TMB})_4^{2+}$  are quite unstable.<sup>3</sup> (No reaction between  $\text{Rh}_2(\text{TMB})_4^{2+}$  and  $\text{HCl}$  is observed at room temperature, in the absence of air, suggesting that the equilibrium between the  $d^8-d^8$  and  $d^7-d^7$  adducts lies far to the side of the  $d^8-d^8$  complex.) Therefore, the monohydride rapidly loses  $\text{H}^+$ , regenerating  $\text{Rh}_2(\text{TMB})_4^{2+}$ . The cycle accounts for the generation of  $\text{H}^+$  (identified electrochemically at a platinum button electrode and by the decrease in pH) and the observed electrocatalytic behavior of  $\text{Rh}_2(\text{TMB})_4^{2+}$ .

The competitive reaction of  $\text{Rh}_2(\text{TMB})_4^{2+}$  with  $\text{CH}_2\text{Cl}_2$  can also be accounted for with a pathway similar to that for the hydrogen-atom abstraction. The  $d^8-d^7$  complex abstracts a halide atom, generating a monohalide species. The monohalide, rather than losing halide to regenerate  $\text{Rh}_2(\text{TMB})_4^{2+}$ , reacts to yield  $\text{Rh}_2(\text{TMB})_4\text{Cl}_2^{2+}$ .

This result demonstrates that highly reactive  $d^8-d^7$  species can be generated electrochemically, and that complexes with an oxidizing hole in the  $d\sigma^*$  orbital can participate in atom abstraction processes. This is the first direct example of chemistry for a  $d^8-d^7$  complex.

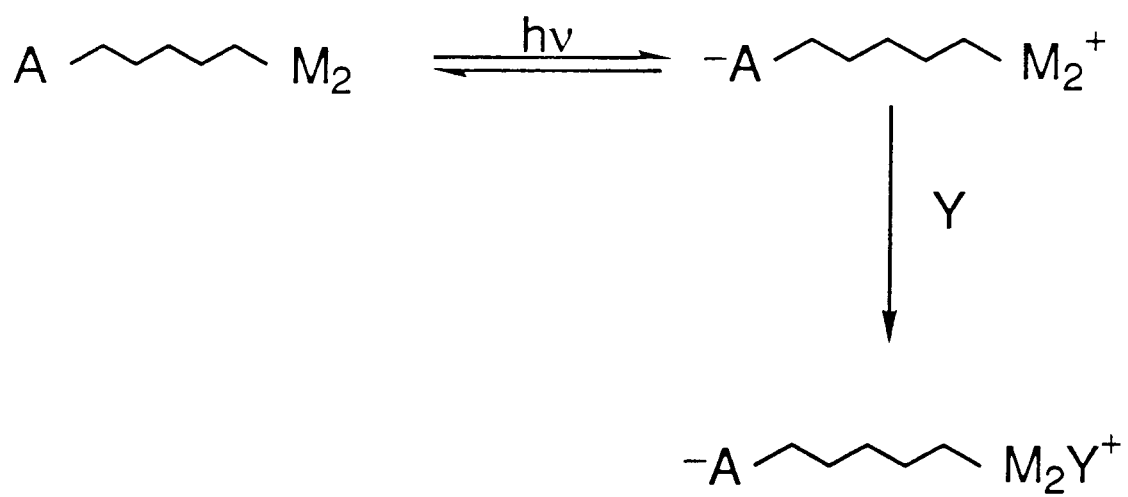
The importance of this result is twofold. First, it demonstrates that the  $d^8-d^7$  complex can be chemically intercepted. In constructing a system that can efficiently harvest light and store this energy in the form of a charge-separated state, some means of intercepting the charge-separated state is necessary. Systems using binuclear  $d^8$

**Figure 4.38.** (a) Pictorial representation of an electrocatalytic cycle which regenerates the  $\text{Rh}_2^{2+}$  complex. (b) Atom-transfer pathway for electrochemical oxidation of 1,4-cyclohexadiene.



complexes covalently linked to pyridinium acceptors have been shown to yield charge-separated species upon excitation.<sup>24</sup> Excited-state electron transfer generates the one-electron oxidized metal complex and the one-electron reduced pyridinium. Intercepting the  $d^8-d^7$  species of the charge-separated state would inhibit the back-electron-transfer step, creating a system where charge separation coupled with chemistry at the oxidized metal complex serves to harvest and store the photon energy (Figure 4.39). Second, the ability to generate reactive  $d^8-d^7$  complexes electrochemically should make possible reactions with the most inert hydrocarbons. While the  $^3(d\sigma^*p\sigma)$  excited state may possess a sufficient driving force for reaction with hydrocarbons with D(C-H) in excess of 100 kcal/mol, the expected rates for these reactions are orders of magnitude less than the lower limit for the observed rates for the  $^3(d\sigma^*p\sigma)$  excited state. For  $Pt_2(P_2O_5H_2)_4^{4-}$ , the lower limit of the observed rates is  $10^4 \text{ M}^{-1}\text{s}^{-1}$  for substrate with a D(C-H) of approximately 90 kcal/mol.<sup>2</sup> Electrochemically generated  $d^8-d^7$  complexes may overcome the kinetic limitation of the ns to  $\mu\text{s}$   $^3(d\sigma^*p\sigma)$  lifetime.

**Figure 4.39.** Pictorial representation of chemical step following charge separation to trap the charge-transfer state.



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**Chapter 5****Synthesis and Characterization of  $\text{Ir}_2(\text{TMB})_4\text{H}_2^{2+}$**



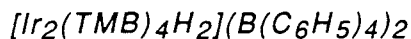
## Introduction

Hydrogen-atom transfer has been established as an important reaction pathway for the triplet  $d\sigma^*p\sigma$  excited state of binuclear  $d^8$  complexes.<sup>1-4</sup> Substrates that serve as hydrogen-atom donors include hydrocarbons (e.g., cyclohexene and 1,4-cyclohexadiene), alcohols with  $\alpha(\text{C-H})$  bonds, and triorganosilanes, -germanes, and -stannanes. Irradiation of the metal complex in the presence of these substrates produces complexes of the type  $\text{M}_2\text{H}_2$ . Earlier work has characterized the product formed with  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$  as the binuclear platinum(III) dihydride,  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{H}_2^{4-}$  ( $\text{Pt}_2\text{H}_2$ ).<sup>5</sup> In this chapter, the characterization of the dihydride of  $\text{Ir}_2(\text{TMB})_4^{2+}$  ( $\text{Ir}_2\text{H}_2$ ) is reported. In addition to NMR, UV-Vis, IR, and Raman spectra, the complex has been characterized crystallographically. The reactivity of  $\text{Ir}_2\text{H}_2$  is also reported.

## Experimental

### Synthesis

Standard Schlenk and high-vacuum techniques were used. Acetonitrile was used as received, freeze-pump-thaw degassed, and stored under vacuum over activated alumina. 1,4-Cyclohexadiene was distilled under argon from  $\text{NaBH}_4$ , freeze-pump-thaw degassed, and stored under vacuum. 9,10-Dihydroanthracene was recrystallized three times from absolute ethanol. Standard procedures were used to prepare  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$ .<sup>6</sup> The  $^1\text{H}$  NMR spectra were obtained on a 400 MHz JNM-GX400 FT NMR spectrometer. The IR spectra were measured on a Beckman IR 4240. Absorption spectra were recorded with a Hewlett-Packard 8450A spectrophotometer or a Shimadzu UV-260 spectrophotometer. Elemental analyses were obtained from Galbraith Laboratories, Inc.



A twentyfold excess of 1,4-cyclohexadiene was added to an acetonitrile solution of  $[\text{Ir}_2(\text{TMB})_4](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in an inert atmosphere. The solution was freeze-pump-thaw degassed three times and photolyzed for 2 hrs, using a 1000 W high-pressure Hg/Xe lamp equipped with a Corning cutoff filter ( $\lambda_{\text{ex}} > 550$  nm). The final solution was a clear, light orange. The solvent was removed under vacuum, being careful to shield the material from light. An off-white powder was obtained. The solid was recrystallized from acetonitrile/toluene. Calculated for  $\text{Ir}_2\text{C}_{95}\text{H}_{114}\text{N}_8\text{B}_2$ : C, 64.3; H, 6.5; N, 6.3. Found: C, 64.8, H, 6.8; N, 6.2.  $^1\text{H}$  NMR:  $\delta$  ( $\text{CD}_3\text{CN}$ , 20 °C): -10.6 (singlet, 1H, Ir-H), 1.4 (broad singlet singlet,  $\text{CH}_3$ ), 1.6 (broad singlet,  $\text{CH}_3$ ), 1.8 (broad singlet,  $\text{CH}_2$ ), 6.9 (triplet, 8H), 7.04 (triplet, 16H), 7.33 (multiplet, 16H). IR (Nujol mull, NaCl plates): 1940  $\text{cm}^{-1}$  (m),  $\nu(\text{Ir-H})$ ; 2160  $\text{cm}^{-1}$  (s),  $\nu(\text{N}\equiv\text{C})$ . UV-Vis ( $\text{CH}_3\text{CN}$ , 25 °C): 320 nm ( $\epsilon = 17000 \text{ M}^{-1}\text{cm}^{-1}$ )

A similar procedure was followed for the photolysis with 9,10-dihydroanthracene. The material obtained from the removal of the solvent contained the desired iridium complex, anthracene, and 9,10-dihydroanthracene. This material was washed with dry, degassed benzene to remove the organics. The remaining solid was dissolved in  $\text{CH}_3\text{CN}$  and filtered. The solvent was removed from the filtrate under vacuum, yielding a light blue-green powder.

## Reactivity

Acetonitrile and acetonitrile- $\text{d}_3$  were used as received, freeze-pump-thaw degassed, and stored under vacuum over activated alumina. HCl was freeze-pump-thaw degassed prior to addition to the reaction vessel. Styrene was used as received, freeze-pump-thaw degassed, vacuum-transferred to a graduate cylinder, and a known volume of styrene vacuum transferred to the reaction vessel.

### *Thermal Stability of $\text{Ir}_2(\text{TMB})_4\text{H}_2^{2+}$*

A  $\text{CH}_3\text{CN}$  solution of  $\text{Ir}_2\text{H}_2$  was prepared in a two-compartment cell consisting of a 10 ml bulb and a 1 cm pathlength square cuvette. The UV-Vis spectrum of the sample was measured. The cell was wrapped in foil to prevent exposure to light and immersed in an 80 °C oil bath. The UV-Vis spectrum was remeasured after 12 hrs. A 5% decrease in the intensity of the band at 320 nm was observed.

### *$\text{Ir}_2(\text{TMB})_4\text{H}_2^{2+} + \text{HCl}$*

An atmosphere of HCl was introduced to an  $\text{CH}_3\text{CN}$  solution of  $\text{Ir}_2\text{H}_2$  in a two-compartment cell. The UV-Vis spectrum of the sample was measured before and after the addition of HCl.

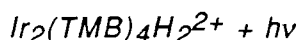
### *$\text{Ir}_2(\text{TMB})_4\text{H}_2^{2+} + \text{Styrene}$*

An  $\text{CH}_3\text{CN}$  solution of  $\text{Ir}_2\text{H}_2$  was prepared in a two-compartment cell. To this solution was added a large excess of styrene. A slow growth of  $\text{Ir}_2(\text{TMB})_4^{2+}$  was observed in the UV-Vis spectrum of the sample over a 24 hr period.

Two samples of  $\text{Ir}_2\text{H}_2$  were prepared in vacuum-adapted NMR tubes. To one of these solutions was added a thirtyfold excess of styrene. Both tubes were stored in the dark at room temperature. NMR spectra were recorded every 24 hrs over a 7-day period.

No change was observed for the solution that did not contain styrene. For the sample with styrene, resonances for  $\text{Ir}_2(\text{TMB})_4^{2+}$  slowly grew into the spectrum with loss of the signals for  $\text{Ir}_2(\text{TMB})_4\text{H}_2^{2+}$ . The sample that contained styrene was clear, dark-blue after 7 days. The NMR tube, which contained the styrene solution, was opened after 7 days, and the contents transferred to a round-bottom flask. The volatiles were vacuum-transferred to a new NMR tube. This solution was clear and colorless. A dark-blue material remained in the round-bottom flask. The NMR of the solvent, in addition

to strong styrene resonances, contained two weak signals attributable to the ethyl group of ethyl benzene ( $\delta$  (ppm): 1.2 (triplet, 3H, CH<sub>3</sub>); 2.62 (quartet, 2H, CH<sub>2</sub>)). These resonances appeared in the NMR spectrum of the reaction mixture as the NMR spectrum of Ir<sub>2</sub>(TMB)<sub>4</sub><sup>2+</sup> grew in.



A CH<sub>3</sub>CN solution of Ir<sub>2</sub>H<sub>2</sub>, prepared on a high-vacuum line in a two-compartment cell, was photolyzed using a 1000 W high-pressure Hg/Xe lamp equipped with Corning cut-off and band-pass filters. A set of filters was used, which yielded %T > 0.1% from 329 to 383 nm. The reaction was monitored by UV-Vis spectroscopy. Disappearance of the 320 nm band assigned to the  $d\sigma \rightarrow d\sigma^*$  transition of Ir<sub>2</sub>H<sub>2</sub> was observed with the growth of bands attributed to Ir<sub>2</sub>(TMB)<sub>4</sub><sup>2+</sup> ( $\lambda_{\text{max}}$  625 and 380 nm) and a new feature at 480 nm.

### Raman Spectroscopy

Raman and resonance Raman spectra were measured by Dr. Thomas Loehr at the Oregon Graduate Center.

### X-Ray Data Collection and Reduction for



Long, plate-like crystals were grown by slow evaporation of a toluene/acetonitrile solution; they appeared dichroic under a microscope with large, colorless faces and narrow, light-blue faces. A crystal was mounted in a capillary with a small amount of grease holding it in place, and was centered on a CAD-4 diffractometer equipped with graphite-monochromated MoK $\alpha$  radiation. Cell dimensions and an orientation matrix were obtained from the setting angles of 25 reflections with  $15^\circ < 2\theta < 25^\circ$ , and the intensity of reflections ( $\pm h, \pm k, l$ ) was measured out to  $2\theta$  of  $30^\circ$ . Absences

in the data for  $0k0$  with  $k$  odd identified the space group as either  $P2_1$  or  $P2_1/m$ . The coordinate of the iridium atoms, obtained from a Patterson map, suggested  $P2_1/m$ . The remaining nonhydrogen atoms were found by successive structure factor-Fourier calculations; however, the resulting structure appeared to be disordered in several regions.

Before attempting to derive an ordered model in the space group  $P2_1/m$ , an oscillation photograph was taken, the axis of rotation being the needle axis,  $[2,0,-1]$ .<sup>7</sup> This photograph showed weak, intermediate layer lines, which were interpreted in terms of a doubling of the  $c$  axis; none of these weak reflections had been noted by the CAD-4 diffractometer during the initial search. The crystal was remounted on a  $P2_1$  diffractometer again equipped with  $\text{MoK}\alpha$  radiation, and new cell dimensions were obtained. These were averaged with the cell dimensions determined by the CAD-4 (with the  $c$  axis doubled) to give the final values listed in Table 5.1. All reflections out to  $2\theta$  of  $30^\circ$  were measured twice ( $\pm h, \pm k, -l$  for  $0^\circ$  to  $20^\circ$ ,  $\pm h, \pm k, +l$  for  $20^\circ$  to  $30^\circ$ , and  $+h, +k, \pm l$  for  $h = 0, 1, 2$   $2\theta$  from  $30^\circ$  to  $40^\circ$ ). Absences in the data for  $0k0$ ,  $k=2n+1$  and  $h0l$ ,  $h+l=2n+1$  identified the space group as  $P2_1/c$ , #14. The weakness of reflections with  $l$  odd is due to the fact that the iridium atoms lie very close to  $y = 0.25$  and hence they plus many of the lighter atoms show a pseudotranslation of  $1/2$  along  $c$ .

Backgrounds were measured for each reflection at each end of scan; an average background as a function of  $2\theta$  was calculated and used to correct the measured counts. Absorption corrections were made, first on the basis of the correct value of  $\mu$  ( $38 \text{ cm}^{-1}$ ) and later, on the basis of a revised value ( $32.5 \text{ cm}^{-1}$ ), which led to improved agreement between symmetry-related measurements. With this revised value, the goodness-of-fit for merging the 9434 measured intensities into 4692 independent reflections was 0.95; 3915 of these reflections had  $F_o > 0$  and 2005 had  $F_o^2 > 3\sigma(F_o^2)$ . Variances of the

**Table 5.1.** Crystal data for  $[\text{Ir}_2(\text{C}_{10}\text{H}_{16}\text{N}_2)_4\text{H}_2](\text{B}(\text{C}_6\text{H}_5)_4)_2 \cdot \text{CH}_3\text{C}_6\text{H}_5$ .

Formula : $\text{Ir}_2\text{N}_8\text{C}_{95}\text{B}_2\text{H}_{114}$	Formula Weight: 1774.03
Crystal Color: Dichroic, colorless to light blue	Habit: flat plate
$a = 10.54(2) \text{ \AA}$	$\alpha = 90.0^\circ$
$b = 31.02(4) \text{ \AA}$	$\beta = 91.57(3)^\circ$
$c = 27.05(4) \text{ \AA}$	$\gamma = 90.0^\circ$
$V = 8840(13) \text{ \AA}^3$	$Z = 4$
$d_{\text{calc}} = 1.333(2) \text{ gcm}^{-3}$	
$\lambda = 0.71073 \text{ \AA}$	
Graphite monochromator	
Space group: $\text{P}2_1/\text{c}$ (#14)	Absences: $h0l, l = 2n+1; 0k0, k = 2n+1$
Crystal size: $0.2 \times 0.074 \times 0.8 \text{ mm}$	$\mu = 32.5 \text{ cm}^{-1}$
	Transimission coefficients: 0.53 - 0.79
Number of data collected: 9434	
Number of independent reflections: 4692	
Number with $F_o^2 > 0$ : 3915	
Number with $F_o^2 > 3\sigma(F_o^2)$ : 2005	
Goodness-of-fit for 4692 measured intensities: 1.16	

individual reflections were assigned, based on counting statistics plus an additional term,  $0.14|I|^2$ . Structure factors were taken from standard references.<sup>8</sup>

The earlier  $P2_1/m$  model was converted to an ordered model in  $P2_1/c$ . After several cycles of full-matrix least squares, minimizing  $\sum w(F_o^2 - F_c^2)^2$ , hydrogen atoms were introduced at calculated positions ( $C-H = 0.95 \text{ \AA}$ ). Each hydrogen was given an isotropic thermal parameter 10% greater than that of the carbon to which it was bonded. The hydrogens bound to iridium appeared on difference maps and were placed  $1.6 \text{ \AA}$  from the iridium atoms. No parameters of the hydrogen atoms bound to iridium were further adjusted, nor were those of the toluene ring, which appeared as diffuse maxima in a late difference map. In the final refinement cycles, 411 parameters were adjusted: the coordinates of 100 atoms, the anisotropic thermal parameters of the Ir atoms, the isotropic thermal parameters for the 98 C, N, and B atoms (excluding the toluene), and a scale factor. The final goodness-of-fit,  $[\sum w(F_o^2 - F_c^2)^2 / (n - p)]^{1/2}$ , was 1.16 for  $n = 4692$  and  $p = 411$ ; the R's were 0.118 for the 3915 reflections with  $F_o^2 > 0$  and 0.045 for the 2005 reflections with  $F_o^2 > 3\sigma(F_o^2)$ . Final parameters are given in Table 5.2, bond lengths and angles in Table 5.3, anisotropic coefficients for the Ir atoms in Table A3.1, hydrogen coordinates in Table A3.2 and observed and calculated F's in Table A3.3.

Table 5.2. Final Parameters (x,y,z and  $U_{eq} \times 10^4$ ).

Atom	x	y	z	$U_{eq}$ or $B$
IR1	45(.9)	2540(.5)	3276(.3)	298(4)
IR2	2126(.9)	2564(.4)	4010(.3)	321(4)
C1	1041(27)	2185(8)	2832(11)	4.4(8) *
N1	1679(21)	1998(6)	2560(8)	4.2(6) *
C2	2426(29)	1722(8)	2212(11)	3.9(7) *
C3	1959(31)	1263(9)	2278(12)	7.0(10) *
C4	2201(28)	1902(8)	1703(11)	6.5(9) *
C5	3767(28)	1757(8)	2425(10)	5.5(8) *
C6	4372(27)	2219(8)	2360(10)	5.5(8) *
C7	6060(28)	2041(8)	3037(11)	5.6(8) *
C8	5748(27)	2804(8)	2728(11)	5.8(8) *
C9	5154(25)	2375(9)	2835(10)	5.3(7) *
N2	4177(18)	2431(7)	3211(7)	4.0(5) *
C10	3395(21)	2523(10)	3504(8)	4.3(6) *
C11	712(25)	3059(8)	2980(10)	3.7(7) *
N3	1091(20)	3358(7)	2786(8)	4.0(6) *
C12	1861(28)	3716(8)	2531(11)	4.1(7) *
C13	2662(31)	3498(9)	2175(12)	6.6(9) *
C14	943(26)	4057(8)	2320(10)	5.1(8) *
C15	2691(25)	3900(8)	2975(10)	4.5(7) *
C16	1964(27)	4141(8)	3348(11)	5.4(8) *
C17	1451(30)	4253(8)	4248(12)	7.3(9) *
C18	3752(33)	4133(9)	4056(13)	7.2(10) *
C19	2270(33)	4042(9)	3897(12)	6.2(9) *
N4	2179(21)	3549(7)	3958(8)	3.9(6) *
C20	2108(30)	3184(9)	3990(12)	4.3(8) *
C21	-1073(30)	2858(8)	3713(11)	4.3(8) *
N5	-1749(25)	3060(7)	3938(9)	5.9(7) *
C22	-2584(33)	3295(10)	4303(13)	7.3(9) *
C23	-2512(31)	3763(9)	4129(13)	7.2(10) *
C24	-3944(35)	3099(10)	4246(13)	10.0(11) *
C25	-1953(29)	3200(9)	4790(12)	6.8(9) *
C26	-1930(31)	2755(9)	5017(12)	8.1(10) *
C27	-856(33)	2120(10)	5421(13)	9.2(11) *
C28	-245(34)	2862(10)	5661(14)	9.4(12) *
C29	-829(28)	2572(11)	5248(11)	7.0(8) *
N6	212(19)	2597(7)	4867(7)	5.0(5) *
C30	902(21)	2569(9)	4543(8)	4.7(7) *
C31	-494(24)	2024(8)	3553(9)	2.9(7) *
N7	-852(18)	1684(6)	3706(7)	2.9(5) *
C32	-1280(29)	1246(8)	3887(11)	4.8(7) *
C33	-2454(38)	1349(10)	4233(15)	9.6(12) *
C34	-1691(30)	986(9)	3459(12)	7.1(9) *
C35	-242(31)	1079(8)	4238(12)	6.9(9) *
C36	896(31)	939(8)	3927(11)	6.4(9) *



Atom	<i>x</i>	<i>y</i>	<i>z</i>	<i>U<sub>eq</sub></i> or <i>B</i>
C37	2340(30)	1006(9)	4747(13)	8.4(10)*
C38	3335(35)	938(9)	3933(13)	7.0(10)*
C39	2178(31)	1101(9)	4172(12)	6.2(9)*
N8	2215(21)	1589(7)	4132(8)	3.7(6)*
C40	2193(26)	1967(8)	4075(10)	2.4(7)*
B1	5694(38)	3954(11)	1047(14)	4.5(11)*
C111	6408(26)	3843(8)	1589(10)	2.9(6)*
C112	7483(34)	3555(9)	1634(13)	4.6(9)*
C113	8063(28)	3452(8)	2070(12)	4.6(8)*
C114	7623(30)	3657(9)	2508(11)	6.2(8)*
C115	6681(31)	3928(9)	2491(11)	5.8(8)*
C116	6001(28)	4023(8)	2034(13)	4.4(8)*
C121	6603(32)	4243(8)	717(12)	3.7(8)*
C122	7866(34)	4305(8)	784(11)	4.9(8)*
C123	8693(31)	4550(9)	470(12)	5.6(8)*
C124	8104(33)	4759(8)	79(12)	4.4(8)*
C125	6856(36)	4714(9)	-18(13)	5.0(9)*
C126	6154(31)	4475(8)	301(12)	3.9(8)*
C131	4358(26)	4240(8)	1134(9)	2.9(6)*
C132	4482(27)	4681(9)	1274(10)	5.5(8)*
C133	3281(32)	4917(9)	1354(11)	6.7(9)*
C134	2187(30)	4734(9)	1249(10)	5.8(8)*
C135	2098(28)	4319(9)	1080(10)	5.2(8)*
C136	3182(35)	4072(8)	1011(12)	4.7(8)*
C141	5330(27)	3495(8)	779(11)	3.9(8)*
C142	4743(27)	3159(9)	1031(10)	6.0(8)*
C143	4370(28)	2778(9)	803(12)	6.7(9)*
C144	4508(25)	2729(7)	320(11)	4.6(8)*
C145	5074(26)	3032(8)	68(10)	4.8(7)*
C146	5515(24)	3417(7)	273(10)	3.5(7)*
B2	6035(31)	1034(8)	1136(12)	1.5(8)*
C211	6778(28)	1462(8)	1358(10)	3.2(7)*
C212	7833(32)	1433(8)	1671(12)	4.5(9)*
C213	8508(26)	1788(8)	1908(9)	3.6(7)*
C214	8080(27)	2188(7)	1814(10)	3.7(7)*
C215	7025(26)	2236(7)	1514(9)	3.4(7)*
C216	6400(26)	1877(8)	1286(10)	4.4(8)*
C221	5911(24)	708(8)	1620(10)	3.1(7)*
C222	5828(29)	837(8)	2103(12)	4.8(9)*
C223	5575(27)	543(9)	2488(11)	6.0(8)*
C224	5358(28)	127(10)	2372(12)	6.5(9)*
C225	5403(26)	-21(8)	1905(11)	5.4(8)*
C226	5741(25)	264(8)	1548(10)	4.4(7)*
C231	6933(32)	788(8)	717(11)	3.9(8)*
C232	8270(30)	824(8)	648(11)	4.5(8)*

Atom	$x$	$y$	$z$	$U_{eq}$ or $B$
C233	8924(29)	582(9)	295(12)	5.3(8) *
C234	8311(37)	327(9)	-19(13)	6.3(10) *
C235	7023(38)	284(9)	10(13)	5.3(9) *
C236	6293(30)	497(8)	363(12)	3.6(8) *
C241	4666(27)	1131(7)	893(9)	2.7(6) *
C242	4655(30)	1415(8)	481(11)	4.3(8) *
C243	3409(32)	1502(8)	248(10)	5.3(8) *
C244	2296(30)	1336(9)	419(12)	6.3(8) *
C245	2303(32)	1064(9)	774(13)	6.5(9) *
C246	3494(31)	958(7)	1041(11)	3.1(7) *
CB1	1044	42	2523	4.0 *
CB2	1818	-12	2120	6.0 *
CB3	1335	59	1644	8.0 *
CB4	78	184	1571	10.0 *
CB5	-696	238	1974	8.0 *
CB6	-213	167	2450	6.0 *
CME	1579	-37	3051	12.0 *

$$^a U_{eq} = \frac{1}{3} \sum_i \sum_j [U_{ij}(a_i^* a_j^*)(\vec{a}_i \cdot \vec{a}_j)]$$

\*Isotropic displacement parameter,  $B$

Table 5.3. Bond lengths and angles.

Distance(Å)		Distance(Å)	
IR1 -IR2	2.920(2)	C35 -C36	1.55(4)
IR1 -C1	1.96(3)	C36 -C39	1.57(4)
IR1 -C11	1.94(3)	C37 -C39	1.59(4)
IR1 -C21	1.96(3)	C38 -C39	1.49(5)
IR1 -C31	1.86(3)	C39 -N8	1.52(4)
IR2 -C10	1.94(3)	N8 -C40	1.18(3)
IR2 -C20	1.92(3)	B1 -C111	1.67(5)
IR2 -C30	1.96(3)	B1 -C121	1.60(5)
IR2 -C40	1.86(3)	B1 -C131	1.68(4)
C1 -N1	1.17(3)	B1 -C141	1.64(5)
N1 -C2	1.51(4)	C111-C112	1.45(4)
C2 -C3	1.52(4)	C111-C116	1.40(4)
C2 -C4	1.50(4)	C112-C113	1.35(4)
C2 -C5	1.51(4)	C113-C114	1.43(4)
C5 -C6	1.58(4)	C114-C115	1.30(4)
C6 -C9	1.58(4)	C115-C116	1.44(4)
C7 -C9	1.50(4)	C121-C122	1.35(4)
C8 -C9	1.50(4)	C121-C126	1.41(4)
C9 -N2	1.48(3)	C122-C123	1.45(4)
N2 -C10	1.19(3)	C123-C124	1.37(4)
C11 -N3	1.14(3)	C124-C125	1.34(5)
N3 -C12	1.55(3)	C125-C126	1.37(5)
C12 -C13	1.46(4)	C131-C132	1.42(4)
C12 -C14	1.53(4)	C131-C136	1.38(4)
C12 -C15	1.57(4)	C132-C133	1.48(4)
C15 -C16	1.49(4)	C133-C134	1.31(4)
C16 -C19	1.54(4)	C134-C135	1.37(4)
C17 -C19	1.46(4)	C135-C136	1.39(4)
C18 -C19	1.63(5)	C141-C142	1.40(4)
C19 -N4	1.54(4)	C141-C146	1.41(4)
N4 -C20	1.14(4)	C142-C143	1.39(4)
C21 -N5	1.14(4)	C143-C144	1.33(4)
N5 -C22	1.53(4)	C144-C145	1.31(4)
C22 -C23	1.53(5)	C145-C146	1.39(4)
C22 -C24	1.56(5)	B2 -C211	1.65(4)
C22 -C25	1.49(5)	B2 -C221	1.66(4)
C25 -C26	1.51(4)	B2 -C231	1.68(4)
C26 -C29	1.42(4)	B2 -C241	1.60(4)
C27 -C29	1.48(5)	C211-C212	1.38(4)
C28 -C29	1.55(5)	C211-C216	1.36(4)
C29 -N6	1.53(4)	C212-C213	1.45(4)
N6 -C30	1.16(3)	C213-C214	1.34(4)
C31 -N7	1.20(3)	C214-C215	1.37(4)
N7 -C32	1.52(3)	C215-C216	1.43(4)
C32 -C33	1.60(5)	C221-C222	1.37(4)
C32 -C34	1.47(4)	C221-C226	1.40(4)
C32 -C35	1.52(4)	C222-C223	1.41(4)

Distance(Å)		Angle(°)	
C223-C224	1.35(4)	C11 -IR1 -C1	90.4(11)
C224-C225	1.35(4)	C21 -IR1 -C1	174.9(11)
C225-C226	1.36(4)	C31 -IR1 -C1	86.4(11)
C231-C232	1.43(4)	C21 -IR1 -C11	93.6(11)
C231-C236	1.47(4)	C31 -IR1 -C11	176.2(11)
C232-C233	1.41(4)	C31 -IR1 -C21	89.6(11)
C233-C234	1.32(5)	C20 -IR2 -C10	92.9(12)
C234-C235	1.37(5)	C30 -IR2 -C10	175.9(10)
C235-C236	1.41(5)	C40 -IR2 -C10	88.7(11)
C241-C242	1.42(4)	C30 -IR2 -C20	90.4(11)
C241-C246	1.42(4)	C40 -IR2 -C20	175.9(12)
C242-C243	1.47(4)	C40 -IR2 -C30	87.8(11)
C243-C244	1.37(4)	N1 -C1 -IR1	175.5(24)
C244-C245	1.28(4)	C2 -N1 -C1	174.9(25)
C245-C246	1.47(4)	C3 -C2 -N1	106.4(22)
CB1- CB2	1.390(0)	C4 -C2 -N1	106.7(22)
CB1- CB6	1.390(0)	C5 -C2 -N1	102.6(21)
CB1- CME	1.541(0)	C4 -C2 -C3	114.5(24)
CB2- CB3	1.390(0)	C5 -C2 -C3	109.0(23)
CB3- CB4	1.390(0)	C5 -C2 -C4	116.4(23)
CB4- CB5	1.390(0)	C6 -C5 -C2	113.6(22)
CB5- CB6	1.390(0)	C9 -C6 -C5	113.0(21)
		C7 -C9 -C6	113.3(22)
		C8 -C9 -C6	109.0(22)
		N2 -C9 -C6	103.7(20)
		C8 -C9 -C7	114.6(23)
		N2 -C9 -C7	106.2(21)
		N2 -C9 -C8	109.4(21)
		C10 -N2 -C9	172.8(23)
		N2 -C10 -IR2	169.7(22)
		N3 -C11 -IR1	177.0(23)
		C12 -N3 -C11	168.6(24)
		C13 -C12 -N3	106.2(22)
		C14 -C12 -N3	109.1(21)
		C15 -C12 -N3	101.9(20)
		C14 -C12 -C13	116.4(24)
		C15 -C12 -C13	110.6(23)
		C15 -C12 -C14	111.5(22)
		C16 -C15 -C12	114.5(22)
		C19 -C16 -C15	117.1(24)
		C17 -C19 -C16	115.3(26)
		C18 -C19 -C16	113.1(24)
		N4 -C19 -C16	106.8(23)
		C18 -C19 -C17	109.3(25)
		N4 -C19 -C17	109.6(24)
		N4 -C19 -C18	101.9(23)
		C20 -N4 -C19	178.3(27)

Angle(°)			Angle(°)		
N4	-C20	-IR2	174.8(26)	C114-C113-C112	117.8(27)
N5	-C21	-IR1	174.9(26)	C115-C114-C113	121.4(28)
C22	-N5	-C21	171.8(28)	C116-C115-C114	121.6(28)
C23	-C22	-N5	102.7(24)	C115-C116-C111	119.8(26)
C24	-C22	-N5	107.0(25)	C122-C121-B1	127.4(28)
C25	-C22	-N5	103.1(25)	C126-C121-B1	122.6(27)
C24	-C22	-C23	113.2(27)	C126-C121-C122	110.0(27)
C25	-C22	-C23	115.9(27)	C123-C122-C121	127.0(28)
C25	-C22	-C24	113.4(27)	C124-C123-C122	115.6(28)
C26	-C25	-C22	122.8(27)	C125-C124-C123	121.4(30)
C29	-C26	-C25	123.2(27)	C126-C125-C124	118.5(31)
C27	-C29	-C26	119.7(28)	C125-C126-C121	127.3(29)
C28	-C29	-C26	112.8(27)	C132-C131-B1	118.1(23)
N6	-C29	-C26	106.0(24)	C136-C131-B1	121.2(24)
C28	-C29	-C27	109.5(26)	C136-C131-C132	120.1(25)
N6	-C29	-C27	106.3(24)	C133-C132-C131	116.2(24)
N6	-C29	-C28	100.4(23)	C134-C133-C132	120.3(27)
C30	-N6	-C29	170.1(24)	C135-C134-C133	122.1(28)
N6	-C30	-IR2	175.6(22)	C136-C135-C134	120.9(27)
N7	-C31	-IR1	176.4(21)	C135-C136-C131	119.8(27)
C32	-N7	-C31	178.2(22)	C142-C141-B1	122.3(25)
C33	-C32	-N7	104.6(22)	C146-C141-B1	122.7(25)
C34	-C32	-N7	108.8(22)	C146-C141-C142	114.9(25)
C35	-C32	-N7	107.0(22)	C143-C142-C141	122.9(26)
C34	-C32	-C33	110.7(25)	C144-C143-C142	119.8(27)
C35	-C32	-C33	104.8(24)	C145-C144-C143	119.6(26)
C35	-C32	-C34	119.8(25)	C146-C145-C144	124.0(25)
C36	-C35	-C32	108.1(24)	C145-C146-C141	118.7(23)
C39	-C36	-C35	110.6(24)	C221-B2 -C211	104.5(21)
C37	-C39	-C36	114.9(25)	C231-B2 -C211	109.9(22)
C38	-C39	-C36	114.4(26)	C241-B2 -C211	114.3(22)
N8	-C39	-C36	108.2(23)	C231-B2 -C221	108.3(21)
C38	-C39	-C37	107.2(25)	C241-B2 -C221	110.5(22)
N8	-C39	-C37	104.7(23)	C241-B2 -C231	109.2(22)
N8	-C39	-C38	106.6(24)	C212-C211-B2	122.6(24)
C40	-N8	-C39	175.8(25)	C216-C211-B2	125.3(24)
N8	-C40	-IR2	177.6(23)	C216-C211-C212	112.0(25)
C121-B1	-C111		110.0(25)	C213-C212-C211	126.8(27)
C131-B1	-C111		110.1(24)	C214-C213-C212	117.3(25)
C141-B1	-C111		107.7(24)	C215-C214-C213	118.5(24)
C131-B1	-C121		107.4(25)	C216-C215-C214	122.0(23)
C141-B1	-C121		112.2(26)	C215-C216-C211	123.3(24)
C141-B1	-C131		109.4(24)	C222-C221-B2	125.6(24)
C112-C111-B1			122.4(24)	C226-C221-B2	120.1(22)
C116-C111-B1			122.1(24)	C226-C221-C222	114.0(24)
C116-C111-C112			115.4(25)	C223-C222-C221	122.2(27)
C113-C112-C111			123.8(28)	C224-C223-C222	118.6(27)

Angle( $^{\circ}$ )

C225 -C224-C223	122.4(28)
C226 -C225-C224	117.5(26)
C225 -C226-C221	125.0(25)
C232 -C231-B2	128.9(25)
C236 -C231-B2	117.5(24)
C236 -C231-C232	113.5(25)
C233 -C232-C231	123.3(26)
C234 -C233-C232	121.1(29)
C235 -C234-C233	119.5(32)
C236 -C235-C234	123.7(31)
C235 -C236-C231	118.8(27)
C242 -C241-B2	115.3(23)
C246 -C241-B2	126.7(23)
C246 -C241-C242	118.0(24)
C243 -C242-C241	116.0(25)
C244 -C243-C242	123.4(27)
C245 -C244-C243	120.9(30)
C246 -C245-C244	120.4(29)
C245 -C246-C241	121.1(25)
CB6 - CB1- CB2	120.0(0)
CME- CB1- CB2	120.0(0)
CME- CB1- CB6	120.0(0)
CB3 - CB2- CB1	120.0(0)
CB4 - CB3- CB2	120.0(0)
CB5 - CB4- CB3	120.0(0)
CB6 - CB5- CB4	120.0(0)
CB5 - CB6- CB1	120.0(0)

## Results and Discussion

### Crystal Structure

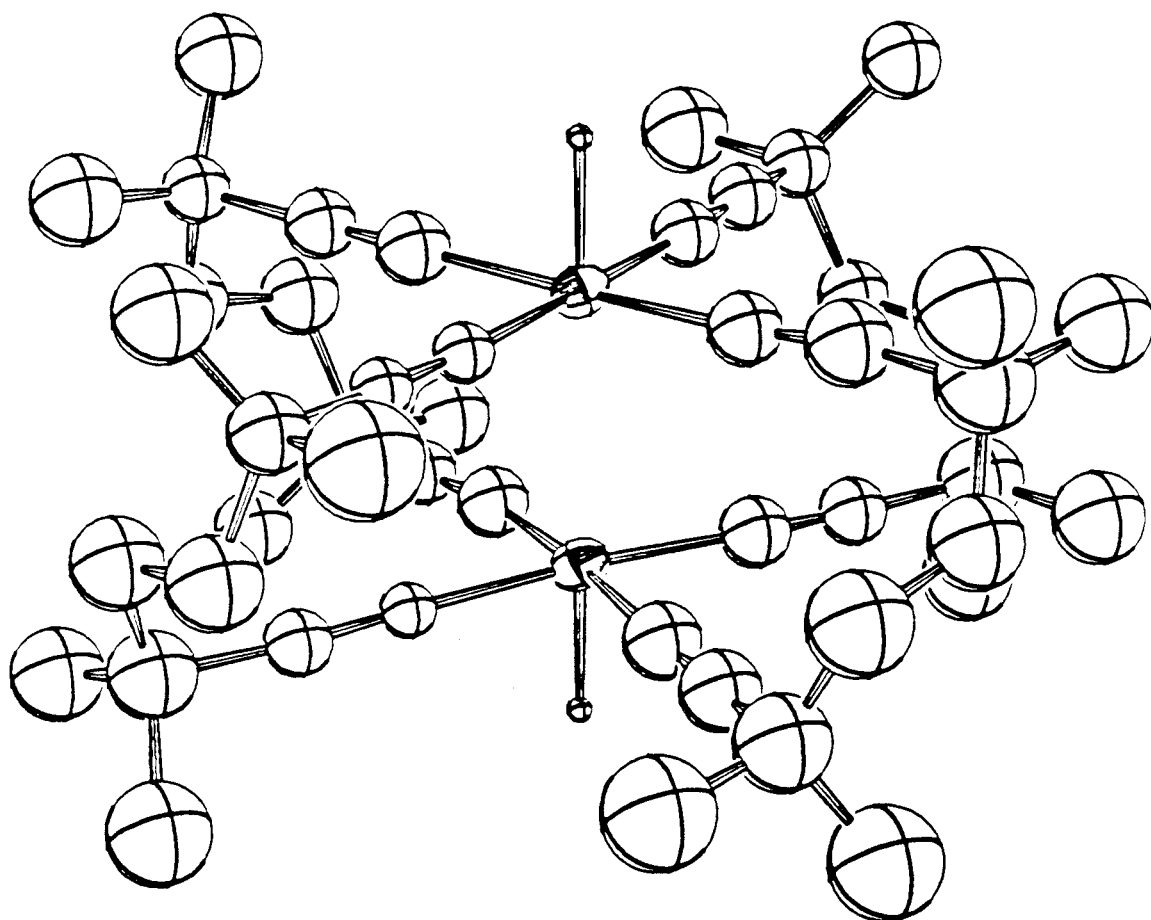
$\text{Ir}_2(\text{TMB})_4\text{H}_2^{2+}$  is the first example of a binuclear metal complex with a trans-dihydride (H-M-M-H) structure (Figure 5.1). The Ir-Ir separation is 2.920(2) Å, approximately 0.3 Å shorter than in the starting  $\text{d}^8$  dimer (3.199(1) Å),<sup>6</sup> indicating formation of an Ir-Ir bond. This Ir-Ir separation is ~0.1 Å longer than that in the analogous  $\text{Ir}_2(\text{TMB})_4\text{I}_2^{2+}$ ,<sup>9</sup> presumably a result of two transhydride effects. The hydrogen-atom positions were not refined; electron density was observed ~1.6 Å away from each Ir atom and assigned to the hydrogen positions.

A comparison of the geometric parameters for  $\text{Ir}_2(\text{TMB})_4\text{H}_2^{2+}$  ( $\text{Ir}_2\text{H}_2$ ),  $\text{Ir}_2(\text{TMB})_4^{2+}$  ( $\text{Ir}_2^{2+}$ ),<sup>6</sup>  $\text{Ir}_2(\text{TMB})_4\text{I}_2^{2+}$  ( $\text{Ir}_2\text{I}_2$ ),<sup>9</sup>  $\text{Rh}_2(\text{TMB})_4\text{Cl}_2^{2+}$  ( $\text{Rh}_2\text{Cl}_2$ ),<sup>9</sup> and  $\text{Rh}_2(\text{TMB})_4^{2+}$  ( $\text{Rh}_2$ )<sup>10</sup> is given in Table 5.4. Comparison of the structure for  $\text{Ir}_2\text{H}_2$  to the structure of  $\text{Ir}_2(\text{TMB})_4\text{Cl}_2^{2+}$  ( $\text{Ir}_2\text{Cl}_2$ ) is preferred over the comparison to  $\text{Ir}_2\text{I}_2$ . However, the structure of the dichloride complex is not known.

The structural parameters observed for  $\text{Ir}_2\text{H}_2$  compare well to those observed for  $\text{Ir}_2\text{I}_2$ . The geometry of the  $\text{Ir}(\text{CN})_4$  unit deviates only slightly from the idealized geometry (M-M-C, 90°; M-C-N, 180°) and compares well to those of  $\text{Ir}_2$  and  $\text{Ir}_2\text{I}_2$ . The two  $\text{Ir}(\text{CN})_4$  planes are nearly parallel to one another with an angle of approximately 3° between the vectors normal to the planes and are twisted with respect to each other (Figure 5.2). The torsion angle, N-M-M'-N', is 28°, identical to that observed for  $\text{Ir}_2$  and  $\text{Ir}_2\text{I}_2$ . A best-plane fit through the  $\text{Ir}(\text{CN})_4$  units shows a slight distortion about each metal center (Table 5.5), each experiencing a different distortion. Ir(1) experiences a tetrahedral distortion. The metal is displaced 0.05 Å toward the center of the dimer, away from the plane defined by the four isocyanide nitrogens. A pyramidal distortion is observed for Ir(2), displaced 0.11 Å away from the plane defined by the four isocyanide nitrogens, toward the center of the dimer.

**Figure 5.1.** View of the structure of  $\text{Ir}_2(\text{TMB})_4\text{H}_2^{2+}$ .





**Table 5.4.** Comparison of geometric parameters of Ir<sub>2</sub>H<sub>2</sub> to other binuclear d<sup>8</sup> and d<sup>7</sup> complexes.

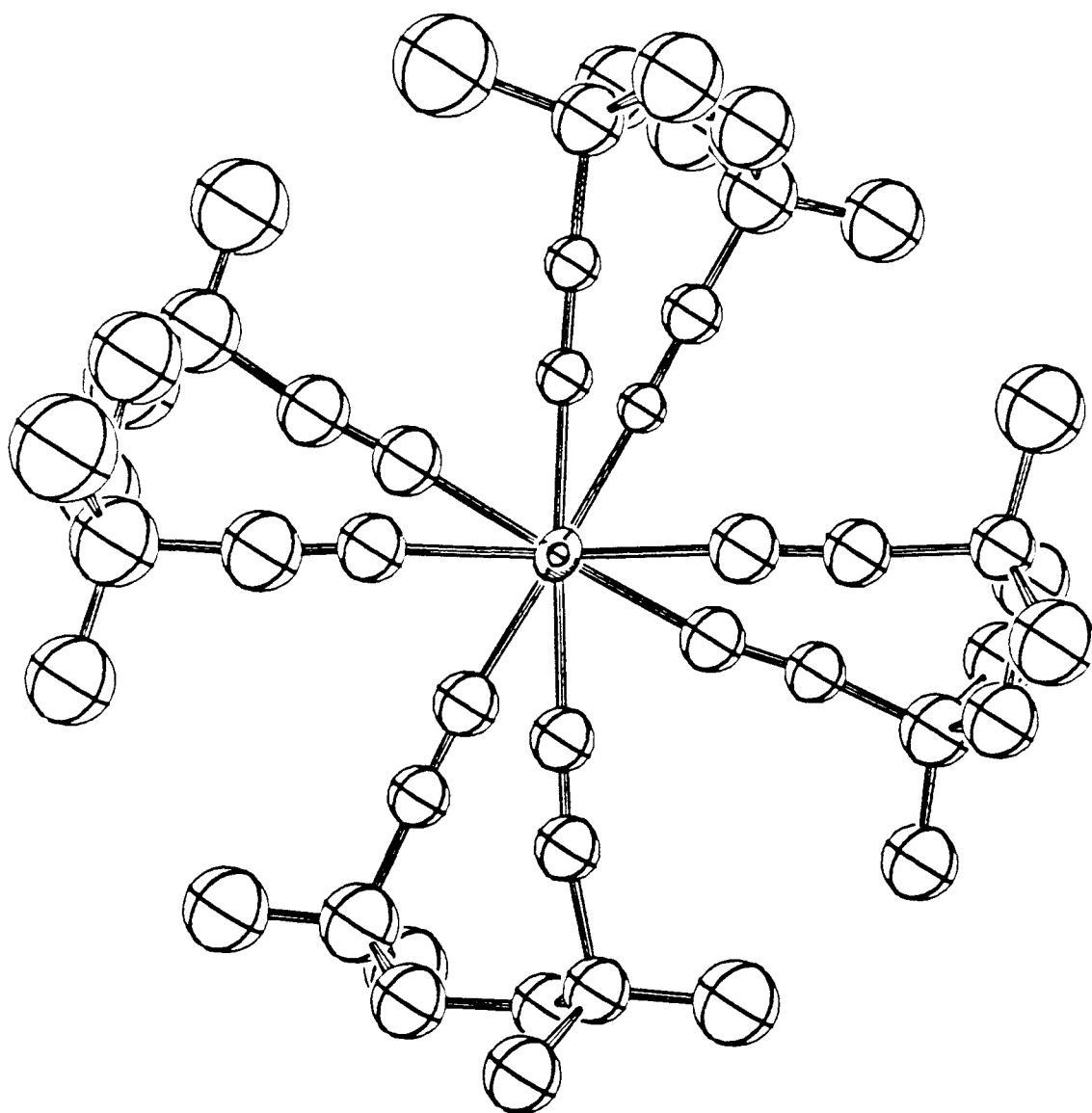
	Ir <sub>2</sub> H <sub>2</sub> <sup>a</sup>	Ir <sub>2</sub> I <sub>2</sub> <sup>b</sup>	Ir <sub>2</sub> <sup>a</sup>	Rh <sub>2</sub> Cl <sub>2</sub> <sup>b</sup>	Rh <sub>2</sub> <sup>c</sup>
M - M	2.920(2) Å	2.803(4) Å	3.199(1) Å	2.773(2) Å	3.262(1) Å
M - X	1.6 Å	2.712(7) Å		2.425(6) Å	
M - C	1.92(1) Å	1.95(1) Å	1.96(1) Å	1.95(1) Å	1.92-1.98 Å
C - N	1.16(1) Å	1.16(2) Å	1.12(1) Å	1.15(1) Å	1.16 Å
ω	28.2(3)°	31(2)°	28.1(1)°	33.1(1)°	31°
C <sub>β</sub> -C <sub>γ</sub> -C <sub>γ</sub> '-C <sub>β</sub> '	136(1)°	130(2)°	141(1)°	137(2)°	
N•••N'	3.43(1) Å		3.68(1) Å		
M - M - C	91.0(3)°	90.12(5)°	91.7(2)°	90.7(4)°	
M - C - N	175.2(1)°	173(1)°	177.12(6)°	176(1)°	175°
displacement of M from plane of the 4 N	Ir(1) 0.05 Å Ir(2) 0.11 Å toward center	0.15 Å toward center of dimer of dimer	0.08 Å toward center of dimer	0.15 Å toward center of dimer	
dihedral angle of M(CN) <sub>4</sub> planes	179°		177°		

a. This work.

b. Reference 9.

c. Reference 10.

**Figure 5.2.** View of the structure of  $\text{Ir}_2(\text{TMB})_4\text{H}_2^{2+}$  viewed down the Ir-Ir vector.



**Table 5.5.** Best plane fit of Ir(CN)<sub>4</sub> units.

Atom	Deviation, Å	Atom	Deviation, Å
Ir1	0.030	Ir2	-0.071
C1	-0.051	C10	-0.025
N1	-0.066	N2	0.023
C11	0.056	C20	-0.007
N3	0.033	N4	0.005
C21	-0.005	C30	-0.035
N5	-0.091	N6	0.030
C31	0.061	C40	-0.006
N7	0.034	N8	0.045

With the different metal-metal separations observed for the Ir complexes, some variation of geometric parameters is expected. No variation in the  $\text{Ir}(\text{CN})_4$  geometry is observed. A compression of the  $\text{Ir}(\text{CN})_4$  units toward one another is found (the  $\text{N}\cdots\text{N}'$  separation decreases 0.25 Å from  $\text{Ir}_2$  to  $\text{Ir}_2\text{H}_2$ ) in addition to a slightly increased pyramidal distortion of the  $\text{Ir}(\text{CN})_4$  moieties. With a compression of the  $\text{Ir}(\text{CN})_4$  units, an increase in the torsion angle ( $\text{N-M-M}'\text{-N}'$ ) might be expected. However, no variation in the torsion angle is observed through the series of complexes; rather, a decrease in the  $\text{C}_\beta\text{-C}_\gamma\text{-C}_\gamma'\text{-C}_\beta'$  dihedral angle is found. A decrease in the metal-metal distance (increased metal-metal bonding interaction) results in a compression or kinking of the four-carbon atom backbone of the TMB ligand. The other geometric parameters of the TMB ligand compare well through the series of complexes.

### Spectroscopy

The vibrational data for  $\text{Ir}_2(\text{TMB})_4\text{H}_2^{2+}$  are summarized in Tables 5.6 and 5.7. The IR spectrum in the region from  $1800\text{ cm}^{-1}$  to  $2400\text{ cm}^{-1}$  is shown in Figure 5.3. A moderate Ir-H absorption is observed at  $1940\text{ cm}^{-1}$ , in line with the terminal metal hydride structure. This resonance is  $100\text{ cm}^{-1}$  to higher energy of the Pt-H band for  $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{H}_2^{4,5}$  indicating a stronger metal hydrogen bond for iridium. The  $\text{C}\equiv\text{N}$  frequency for  $\text{Ir}_2\text{H}_2$  has shifted  $20\text{ cm}^{-1}$  to higher energy from that for  $\text{Ir}_2$  (a  $40\text{ cm}^{-1}$  shift to higher energy from the free ligand)<sup>11</sup> because of the increased  $\sigma$  donation from the isocyanide to the  $\text{Ir}(\text{II})$  center. This increased  $\sigma$  donation reduces the  $\sigma$  antibonding contribution of the carbon lone pair to the  $\text{C}\equiv\text{N}$  bond.

Raman spectra for  $\text{Ir}_2\text{H}_2$  are shown in Figure 5.4. Two low-energy bands are observed. The band at  $38\text{ cm}^{-1}$  is a plasma line. Excitation at 647.1 nm (on resonance for  $\text{Ir}_2$ , see Chapter 2) results in a large enhancement of the  $60\text{ cm}^{-1}$  band (Figure 5.4), identifying this resonance as the Ir-Ir frequency for  $\text{Ir}_2$ . Thus, a small amount of  $\text{Ir}_2$  is present in the  $\text{Ir}_2\text{H}_2$  material (hence the light-blue color of the crystals). The

Table 5.6. Infrared vibrational frequencies.

	$\nu(\text{M-H}), \text{cm}^{-1}$	$\nu(\text{N}\equiv\text{C}), \text{cm}^{-1}$
$\text{Ir}_2\text{H}_2$	1940	2160 a
$\text{Pt}_2\text{H}_2$ b	1840	
$\text{Ir}_2$		2140 a

a. Free ligand,  $2120 \text{ cm}^{-1}$ .

b. Reference 5.

**Table 5.7.** Raman vibrational frequencies.

	$\nu(\text{M-M}), \text{ cm}^{-1}$	$k, \text{ mdyne/\AA}$	$k, \text{ mdyne/\AA}^a$	$\text{M-M}, \text{ \AA}$	
				X-ray	W.R. <sup>a</sup>
$\text{Ir}_2\text{H}_2^b$	136	1.05 <sup>c</sup>	0.86	2.920(2)	2.85
$\text{Ir}_2\text{I}_2^d$	116	2.09 <sup>e</sup>	1.18	2.803(4)	2.55
$\text{Ir}_2\text{Cl}_2^d$	140	1.56 <sup>e</sup>			2.69
$\text{Rh}_2\text{Cl}_2^d$	155	1.141 <sup>e</sup>	1.27	2.773(2)	2.82

a. Calculated using Woodruff's relationship, Reference 13.

b. This work.

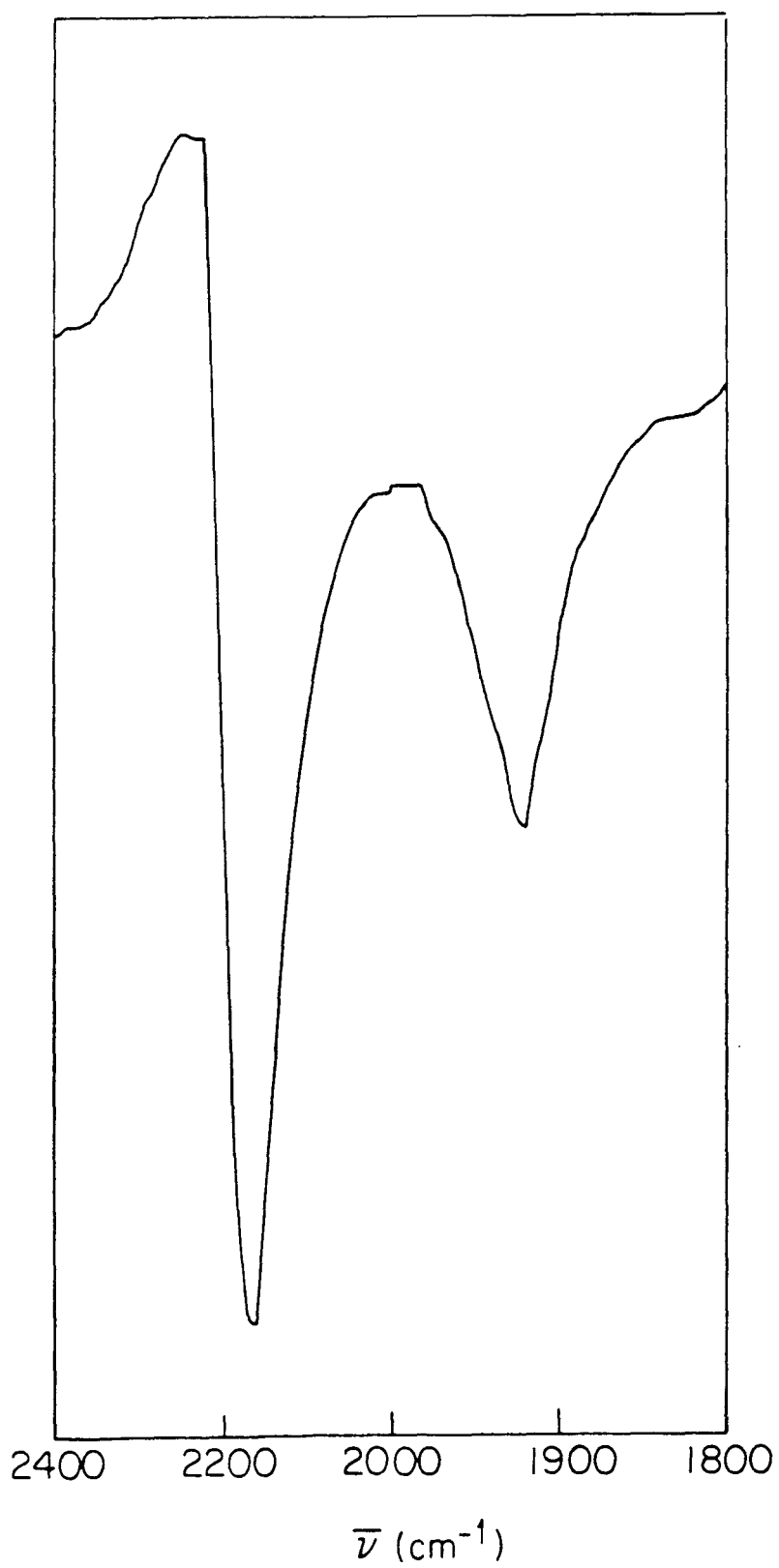
c. Harmonic approximation.

d. Reference 12.

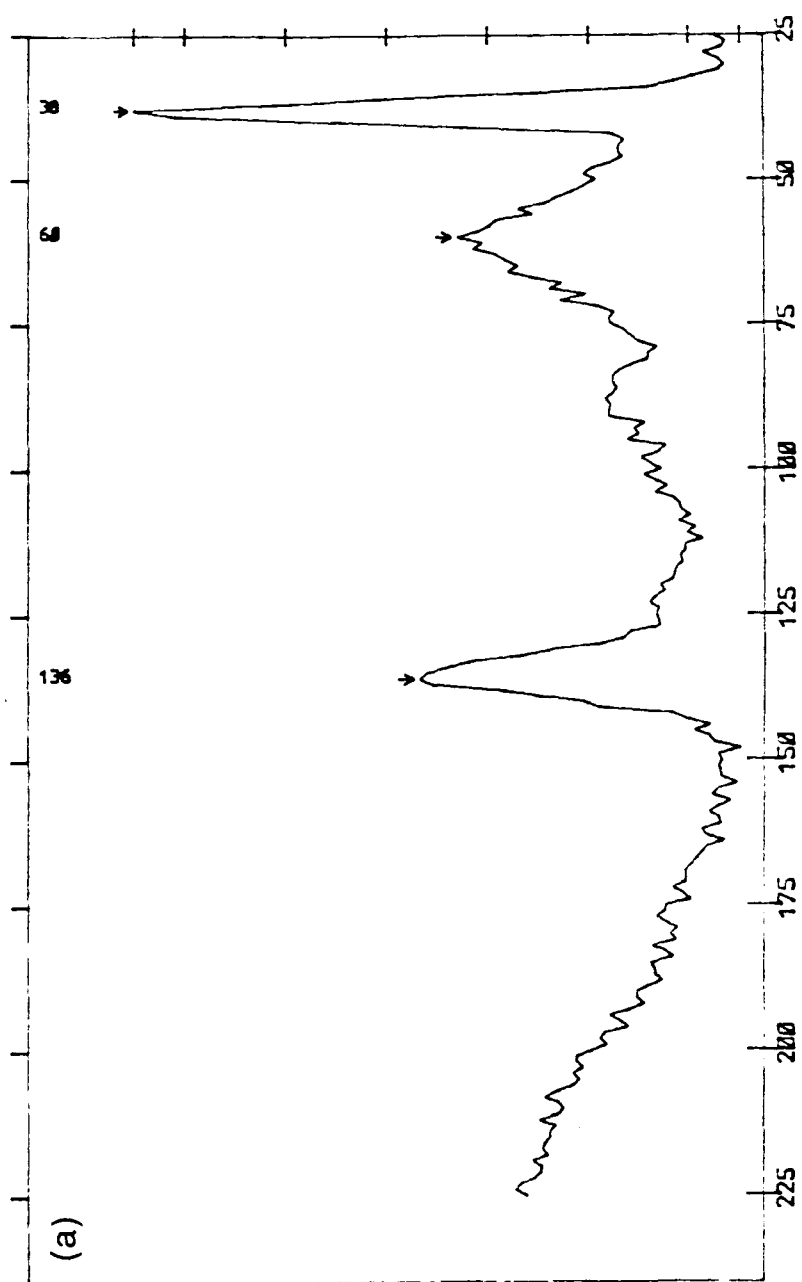
e. Determined from a force-field analysis, Reference 12.

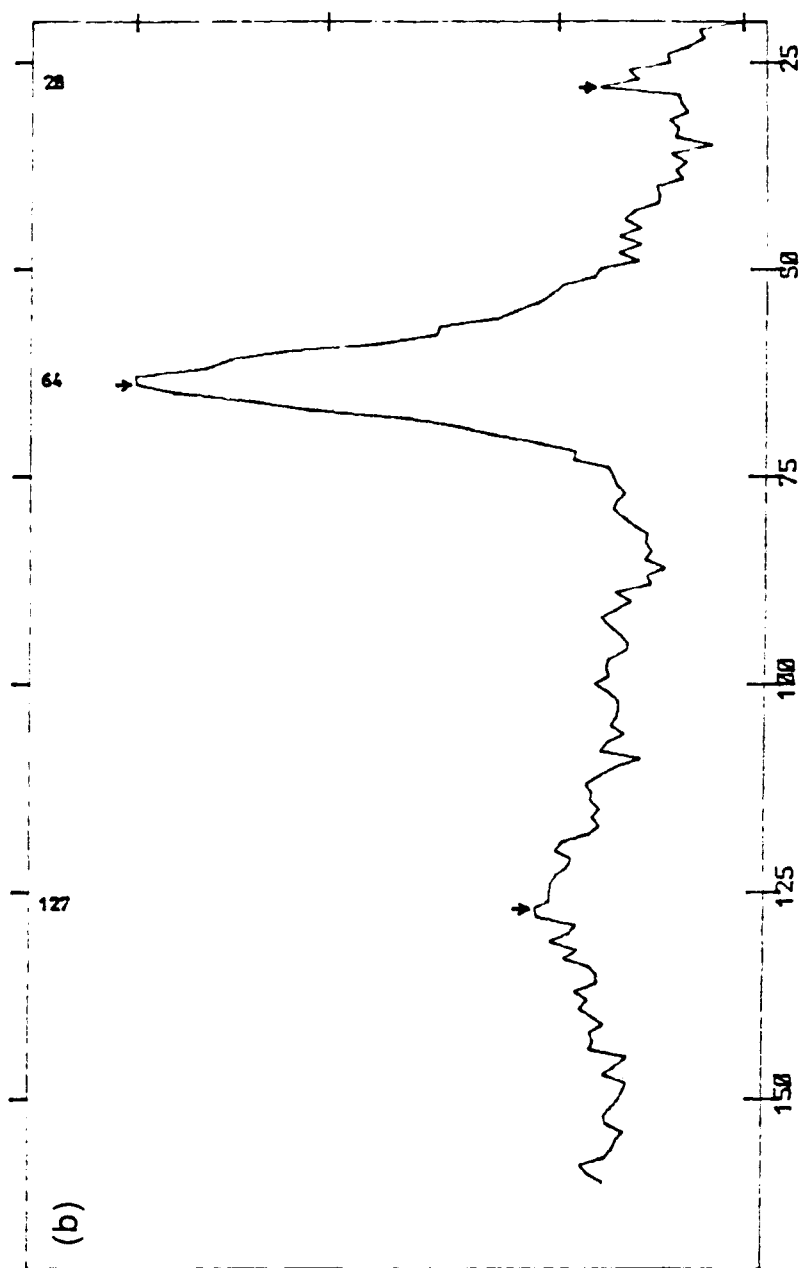


**Figure 5.3.** Infrared spectrum of  $[\text{Ir}_2(\text{TMB})_4\text{H}_2](\text{B}(\text{C}_6\text{H}_5)_4)_2 \cdot \text{CH}_3\text{C}_6\text{H}_5$ .



**Figure 5.4.** Raman spectra of  $[\text{Ir}_2(\text{TMB})_4\text{H}_2](\text{B}(\text{C}_6\text{H}_5)_4)_2 \cdot \text{CH}_3\text{C}_6\text{H}_5$ : (a) Solid, 90 K,  $\lambda_{\text{ex}} = 488 \text{ nm}$ ; (b) Solid, 90 K,  $\lambda_{\text{ex}} = 647.1 \text{ nm}$ .





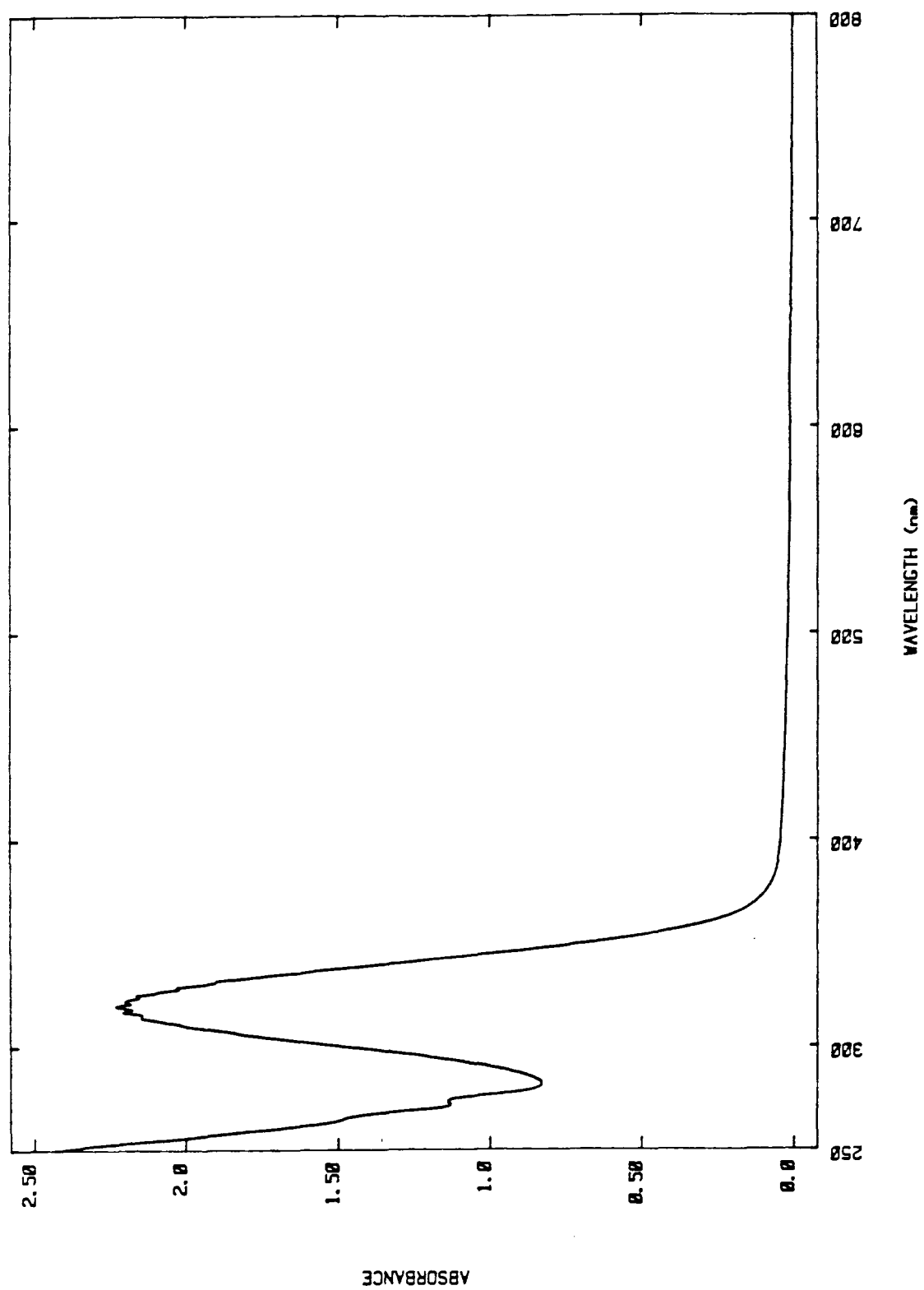
ratio of the concentration of  $\text{Ir}_2$  to  $\text{Ir}_2\text{H}_2$  (determined from optical absorption spectra) is 1/40. The band at  $136\text{ cm}^{-1}$  is the Ir-Ir stretching frequency of  $\text{Ir}_2\text{H}_2$ .

While the  $\text{Ir}_2\text{Cl}_2$  structure is not known, a comparison of the Ir-Ir interaction for  $\text{Ir}_2\text{Cl}_2$  to that of  $\text{Ir}_2\text{H}_2$  can be made using the Raman vibrational data.<sup>12</sup> A significantly larger Ir-Ir force constant is found for  $\text{Ir}_2\text{Cl}_2$ , indicating a stronger Ir-Ir interaction, and hence a smaller Ir-Ir separation. An estimate of the Ir-Ir distance for  $\text{Ir}_2\text{Cl}_2$  can be obtained using Woodruff's correlation (Table 5.7).<sup>13,14</sup> A significantly shorter Ir-Ir distance, in comparison to  $\text{Ir}_2\text{H}_2$ , is calculated for  $\text{Ir}_2\text{Cl}_2$ . The longer Ir-Ir distance for  $\text{Ir}_2\text{H}_2$  is a consequence of two transhydride effects.

The optical absorption spectrum of  $[\text{Ir}_2(\text{TMB})_4\text{H}_2](\text{B}(\text{C}_6\text{H}_5)_4)_2$  is shown in Figure 5.5. In analogy to other binuclear  $d^7$ - $d^7$  complexes, the broad feature ( $\lambda_{\text{max}}$  320 nm,  $\epsilon = 17000\text{ M}^{-1}\text{cm}^{-1}$ ) is assigned to the  $d\sigma \rightarrow d\sigma^*$  transition.<sup>12</sup> A weak absorption attributable to " $d\pi \rightarrow d\sigma^*$ " is not present in the spectrum. This band is presumably beneath the  $d\sigma \rightarrow d\sigma^*$  band. The large red shift observed for the  $d\sigma \rightarrow d\sigma^*$  transition from  $\text{Ir}_2\text{Cl}_2$  to  $\text{Ir}_2\text{H}_2$  is a result of the strong  $\sigma$ -donating character of the hydride ligand (a significantly greater charge-transfer contribution for  $\text{H}^-$  in comparison to  $\text{Cl}^-$ ).

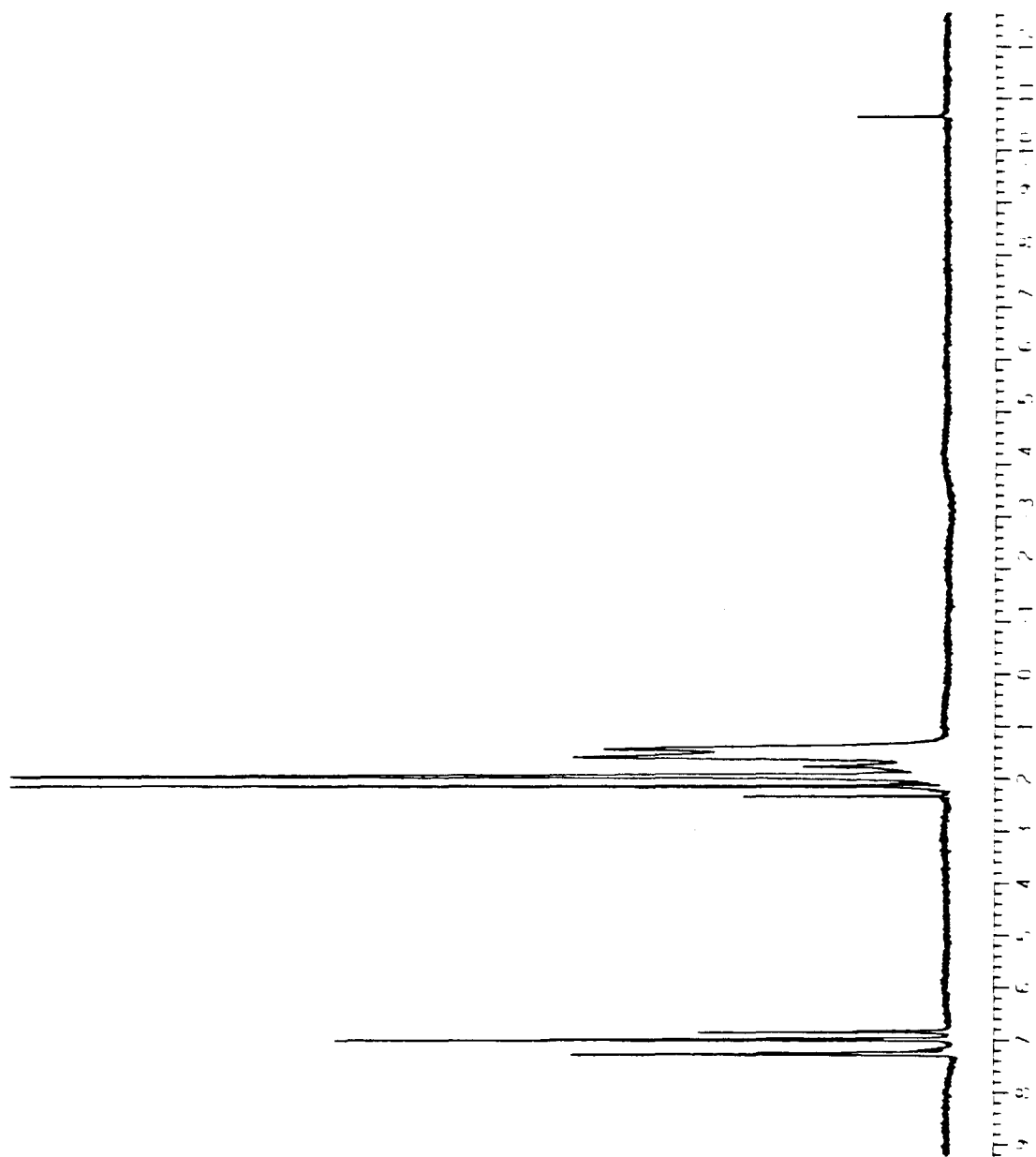
The  $^1\text{H}$  NMR spectrum of  $[\text{Ir}_2(\text{TMB})_4\text{H}_2](\text{B}(\text{C}_6\text{H}_5)_4)_2 \cdot \text{CH}_3\text{C}_6\text{H}_5$  is shown in Figure 5.6. The Ir-H resonance is a sharp singlet at -10.6 ppm. The resonances attributed to the TMB methyls and methylenes are assigned to the broad features at 1.4, 1.6 and 1.8 ppm. The  $\text{B}(\text{C}_6\text{H}_5)_4^-$  resonances are to lower field. One unique feature of this spectrum is the similarity of the TMB spectrum to that of the coalesced spectra for  $\text{Ir}_2(\text{TMB})_4^{2+}$  and  $\text{Ir}_2(\text{TMB})_4\text{I}_2^{2+}$ .<sup>15</sup> While  $\text{Ir}_2\text{H}_2$  is expected to be fluxional in solution at room temperature, the two types of methyl groups (one almost coplanar with the  $\text{Ir}(\text{CN})_4$  unit and the other projecting out over the  $\text{Ir}(\text{CN})_4$  plane, see Figure 5.1) are magnetically inequivalent and give rise to resonances at 1.4 and 1.6 ppm.

**Figure 5.5.** Electronic absorption spectrum of  $[\text{Ir}_2(\text{TMB})_4\text{H}_2](\text{B}(\text{C}_6\text{H}_5)_4)_2$  in  $\text{CH}_3\text{CN}$  at 25 °C.





**Figure 5.6.**  $^1\text{H}$  NMR spectrum of  $[\text{Ir}_2(\text{TMB})_4\text{H}_2](\text{B}(\text{C}_6\text{H}_5)_4)_2 \cdot \text{CH}_3\text{C}_6\text{H}_5$  in  $\text{CD}_3\text{CN}$  at  $20\text{ }^\circ\text{C}$ .



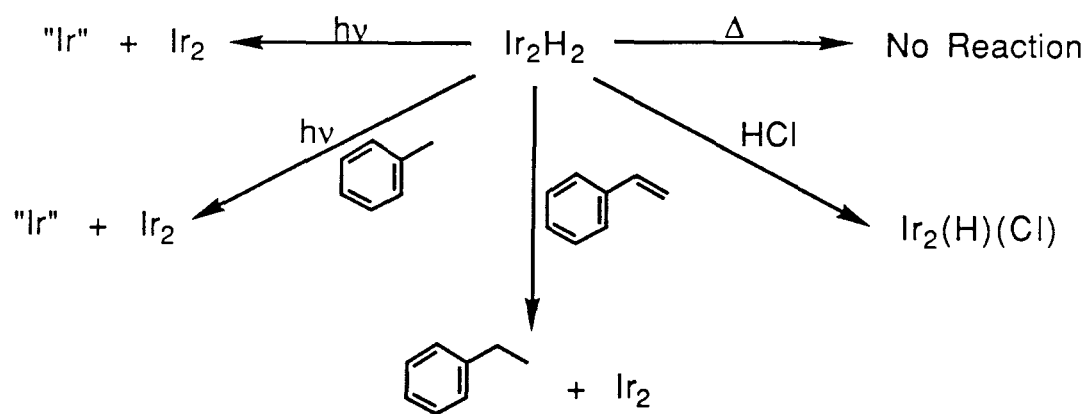
## Reactivity

An exciting aspect of atom-transfer reactivity is the possibility of photocatalytic processes. The metal dihydrides formed in the initial photochemical hydrogen atom-transfer reaction may be turned over to produce  $H_2$  and the  $d^8-d^8$  complex ( $M_2$ ) or reacted with a substrate to effect hydrogenation (and again production of  $M_2$ ).

Isolation and characterization of  $M_2H_2$  afford an opportunity to explore its reactivity (Scheme 5.1). Some of this work has already been reported for  $Pt_2(P_2O_5H_2)_4H_2^{4-}$  ( $Pt_2H_2$ ).<sup>1,5</sup>  $Ir_2H_2$  is found to be thermally quite stable. No significant decomposition of  $Ir_2H_2$  is observed at 80 °C in  $CH_3CN$  over 12 hrs. Slow thermal decomposition of  $Pt_2H_2$  is observed at room temperature. Both  $Pt_2H_2$  and  $Ir_2H_2$  react with HCl. Reaction of  $Ir_2H_2$  yields  $Ir_2(H)Cl$ . For  $Pt_2H_2$ , reaction with HCl and DCl generates  $H_2$  and HD, respectively.  $Pt_2H_2$  photochemically eliminates  $H_2$  to yield  $Pt_2$ . For  $Ir_2H_2$ , photolysis results in production of the starting  $Ir_2$  complex. This reaction is not clean, and a large amount of decomposition of the iridium material is observed.

A potentially more useful reaction has been found. Reaction of  $Ir_2H_2$  with styrene results in a slow, clean conversion to  $Ir_2$  and ethylbenzene. The reaction may be related to the hydrogenation of alkenes by  $Co(CN)_5^{3-}$ , which involves a coordinatively saturated metal hydride that can effect hydrogen atom transfer to give a stable 1-electron reduced transition metal ion.<sup>16</sup> A similar reaction (transhydrogenation of cyclohexene and pentene with isopropanol) is observed for  $Pt_2$ .<sup>17</sup> Thus, combination of the photochemical hydrogen atom-transfer reaction with the hydrogenation step completes a catalytic cycle in which the  $Ir_2$  complex functions as a photochemical two-hydrogen transfer reagent.

**Scheme 5.1.** Observed reactivity of  $[\text{Ir}_2(\text{TMB})_4\text{H}_2](\text{B}(\text{C}_6\text{H}_5)_4)_2$ . Iridium containing by-product not characterized, "Ir".



## References and Notes

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11. An increase in the N≡C stretching frequency from the free ligand to Ir<sub>2</sub> is counter to the expected shift, if  $\pi$ -backbonding were an important part of the metal-isocyanide interaction. Therefore, the  $\pi$ -backbonding contributions to the metal-isocyanide interaction are negligible with  $\sigma$  donation from the isocyanide to the iridium center being the dominant bonding interaction leading to an increase in the N≡C frequency.
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14. Poor agreement is found between the calculated and crystallographically determined Ir-Ir distances for  $\text{Ir}_2\text{I}_2$ . The significant mixing of the metal-halide vibration with the metal-metal vibration for  $\text{Ir}_2\text{I}_2$  makes force-field analysis of this system quite difficult. The force constant for the Ir-Ir bond is not well determined. For  $\text{Ir}_2\text{Cl}_2$  and  $\text{Ir}_2\text{Br}_2$ , the amount of mixing between the metal-halide mode and the metal-metal mode is much less, allowing force-field analyses to produce more reliable estimates of the Ir-Ir force constants. Therefore, better agreement between the calculated and the true Ir-Ir separations are expected. Good agreement is found between the calculated and crystallographically determined metal-metal distances for  $\text{Rh}_2\text{Cl}_2$ .

15.

	Chemical Shift <sup>a</sup>			$T_C, ^\circ\text{C}$	$\Delta G^\ddagger, \text{kcal/mol}$
	-CH <sub>2</sub> -	-CH <sub>3</sub>	-CH <sub>3</sub>		
	(fast)	(fast)	(slow)		
$\text{Ir}_2(\text{TMB})_4^{2+}$ <sup>b</sup>	1.72	1.46	1.33, 1.58	15	14.6
$\text{Ir}_2(\text{TMB})_4\text{I}_2^{2+}$ <sup>b</sup>	2.11	1.54	1.33, 1.63	-14	13.0

a. Given in  $\delta$  (ppm) vs. TMS. "Fast" and "slow" refer to values determined in the fast-exchange and stopped-exchange limits, respectively.

b. In  $\text{CD}_3\text{CN}$ , see reference 9.

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**Chapter 6*****Ab Initio* Calculations of  $\text{Re}_2\text{Cl}_8^{2-}$**



### Introduction

The discovery of multiply bonded metal-metal dimers has provoked numerous studies aimed toward understanding the electronic structure and spectral properties of these systems. To begin to understand the nature of the quadruple metal-metal bond, one can analyze the d orbital interaction for a simple diatomic system  $M_2$ . Figure 6.1 shows the five nonzero overlaps between pairs of d orbitals when two metal atoms approach each other.

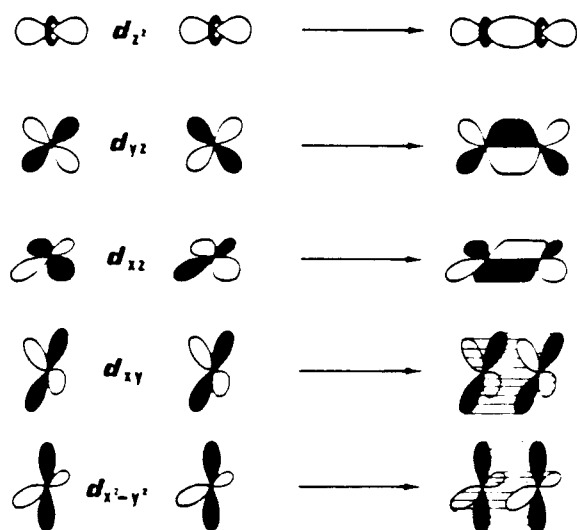


Figure 6.1. d-d overlaps between two metal atoms.

The positive overlap of the  $d_{z^2}$ ,  $d_{xz}$ , and  $d_{yz}$  orbitals gives rise to the  $\sigma$  and  $\pi$  bonds, with the  $\delta$  bonds being formed by the overlap of the  $d_{xy}$  and  $d_{x^2-y^2}$  orbitals. The relative magnitudes of the orbital overlap are  $\delta \ll \pi < \sigma$ . Using the simple Hückel concept that the molecular orbital (MO) energy should be proportional to the orbital overlap, the following orbital energy ordering is predicted.

$$\sigma < \pi \ll \delta < \delta^* \ll \pi^* < \sigma^*$$

In order to explain the nonbridging eclipsed conformation and the observed diamagnetism of  $\text{Re}_2\text{Cl}_8^{2-}$  Cotton proposed, based on the above MO scheme, that the ground-state configuration was  $\sigma^2\pi^4\delta^2$ . Because the  $\delta_{x^2-y^2}$  orbitals of the metal are involved in the metal-ligand  $\sigma$  bonds, the  $\delta_{x^2-y^2}$  and  $\delta_{x^2-y^2}^*$  orbitals were assumed to be high in energy and unimportant in the description of the metal-metal bonds (see Figure 6.2). Using this MO scheme, estimates of the quadruple bond and  $\delta$  bond strengths were made, as well as assignment of the optical absorption spectrum. However, the MO interpretation has been shown to be incorrect.<sup>1</sup>

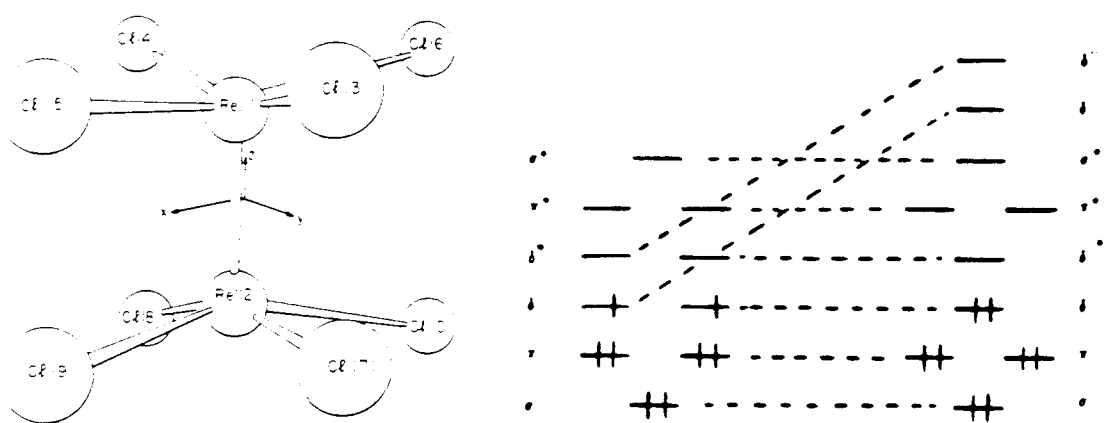


Figure 6.2. d orbitals for  $\text{M}_2$  and  $\text{M}_2\text{X}_8$ .

In the MO formalism, a single configuration is used to describe the electronic structure of a molecule. For a singly bonded system such as  $\text{H}_2$ , the MO wave function is

$$\Psi_{\text{MO}} = [\phi_{\sigma}(1)\phi_{\sigma}(2)] (\alpha\beta - \beta\alpha). \quad 6.1$$

In terms of atomic orbitals, the MO wave function has an equal covalent and ionic contribution.

$$\phi_{\sigma} = 1 + r \quad 6.2$$

$$\Psi_{\text{MO}} = [1 + r][1 + r](\alpha\beta - \beta\alpha) \quad 6.3$$

$$\Psi_{\text{MO}} = [(1r + r1) + (11 + rr)](\alpha\beta - \beta\alpha) \quad 6.4$$

$$\Psi_{\text{cov}} = 1r + r1 \quad \Psi_{\text{ion}} = 11 + rr \quad 6.5$$

The MO wave function for the ground state of  $\text{H}_2$ , while being bound at the equilibrium bond distance, is more than 2 eV above the exact energy. As the molecule is dissociated, the ground-state wave function should converge toward the covalent limit of two hydrogen atoms,  $1r + r1$ . At  $R = \infty$ , the MO wave function leads to an energy that is more than 8 eV above the correct dissociation limit. Since the MO wave function must have equal ionic and covalent contributions for all internuclear separations, it cannot correctly describe the long  $R$  region of the ground-state potential surface. This is a region of low orbital overlap that is best described by a covalent wave function.

For  $\text{Re}_2$ , with a metal-metal separation of 2.24 Å, the orbital overlaps are  $\sigma = 0.63$ ,  $\pi = 0.53$ , and  $\delta = 0.17$ . Because of the low  $\delta$  orbital overlap, the wave function for the  $\delta$  bond is required to be largely covalent. Therefore, an MO wave function will be an extremely inadequate description of the  $\delta$  bond.

In order to correct for the deficiencies of the MO description, a wave function can be constructed using the exact wave functions at  $R = \infty$  (i.e., the valence bond (VB) wave function). The VB wave function for  $\text{H}_2$  is

$$\Psi_{\text{VB}} = [1r + r1](\alpha\beta - \beta\alpha). \quad 6.6$$

This wave function is covalent for all internuclear distances and, while being only slightly better than the MO wave function at short  $R$ , correctly describes the long  $R$  region of the potential surface. Although dissociating to the correct limit, the VB wave function at  $R_e$  is ~1.5 eV above the exact energy of  $\text{H}_2$ .

Expressing the VB wave function for H<sub>2</sub> in terms of the MO wave functions, one sees that the correct description of the long R region requires a wave function that is made up of more than a single configuration.

$$\Psi_{\text{MO}} = [1r + r1](\alpha\beta - \beta\alpha) \quad 6.7$$

$$\varphi_{\sigma} = 1 + r \quad \varphi_{\sigma}^* = 1 - r \quad 6.8$$

$$\varphi_{\sigma}\varphi_{\sigma} = (11 + rr) + (1r + r1) \quad 6.9$$

$$\varphi_{\sigma}^*\varphi_{\sigma}^* = (11 + rr) - (1r + r1) \quad 6.10$$

$$\Psi_{\text{VB}} = [\varphi_{\sigma}\varphi_{\sigma} - \varphi_{\sigma}^*\varphi_{\sigma}^*](\alpha\beta - \beta\alpha) \quad 6.11$$

For the metal dimers, the VB wave function is a multiconfiguration description that correctly describes the weak coupling in the metal-metal bond.

A generalization to the VB description can be made by allowing the MO wave functions to optimize their contributions to the VB wave function. This generalization yields the generalized valence bond (GVB) wave function.

$$\Psi_{\text{GVB}} = [C_1(\varphi_{\sigma}\varphi_{\sigma}) - C_2(\varphi_{\sigma}^*\varphi_{\sigma}^*)](\alpha\beta - \beta\alpha) \quad 6.12$$

A multiconfiguration description is necessary to represent adequately the ground state wave function for the quadruply bonded metal dimers. This level of theory, while providing the correct qualitative description of the quadruple bond, gives a very poor quantitative description of the metal dimer. Using an extensive *ab initio* wave function for Mo<sub>2</sub>Cl<sub>8</sub><sup>4-</sup> and solving self-consistently for the <sup>1</sup>A<sub>2u</sub> excited state ( $\delta \rightarrow \delta^*$  transition), an excitation energy ~1.2 eV higher than the experimental assignment was obtained.<sup>2</sup> Such results have led some theoreticians to conclude that the earlier spectral assignments are incorrect.<sup>2</sup> Others, arguing that the errors are due to an inadequate description of the electron correlation, have corrected the calculated transition energies

by subtracting from these values the difference in the theoretical estimate and experimental assignment for one well-characterized transition.<sup>3</sup>

Recently, it has been shown from studies on Cr<sub>2</sub> and Mo<sub>2</sub> that the large correlation error in the *ab initio* wave function results from a difficulty in describing the atomic electron affinities.<sup>4,5</sup> There exists an atomic correlation energy associate with a doubly occupied atomic orbital that is normally accounted for by doing double excitation out of a bond pair. For the hextuply bonded metal dimers, the number of configurations necessary to describe this atomic correlation is impractical. Reducing the size of the calculation to a manageable level leads to an energy for the doubly occupied atomic orbitals that is too high. However, by making a simple modification to the GVB methods, accurate ground-state properties could be determined.<sup>5</sup>

For the quadruply bonded metal dimers, an accurate description of the doubly occupied atomic orbitals, while being important in determining the ground-state properties, is imperative in the calculation of excitation energies. The importance of the intrapair correlation energy in the excited state is evident from the VB wave function for the <sup>1</sup>A<sub>2u</sub> state.

$$\Psi(^1A_{2u}) = [\varphi_{\delta}\varphi_{\delta} + \varphi_{\delta}^*\varphi_{\delta}^*](\alpha\beta - \beta\alpha) \quad 6.13$$

$$\varphi_{\delta} = l + r \quad \varphi_{\delta}^* = l - r \quad 6.14$$

$$\varphi_{\delta}\varphi_{\delta}^* = (l + r)(l - r) \quad 6.15$$

$$\varphi_{\delta}^*\varphi_{\delta} = (l - r)(l + r) \quad 6.16$$

$$\Psi(^1A_{2u}) = [l\ l - r\ r](\alpha\beta - \beta\alpha) \quad 6.17$$

The <sup>1</sup>( $\delta\delta^*$ ) excited state is described by an ionic wave function, two electrons localized on the same atom. The correlation error associated with having two electrons in the same localized orbital is large, and to describe accurately the <sup>1</sup>( $\delta\delta^*$ ) state, it is necessary to include in the wave function the correct amount of correlation interaction

to account for this correlation error. The  $^1(\delta \rightarrow \delta^*)$  excitation energy is expected to have an error of 1-2 eV if a wave function is used that does not account for the intrapair correlation energy. A similar error is expected for any state that cannot be described by singly occupied orbitals on the metal centers.

For the *ab initio* calculations to date,<sup>1-3</sup> the level of correlation in the wave functions has not been sufficient to account for the intrapair correlation energy. Again, the size of the calculation necessary for an accurate description is impractical. However, the correlation neglected in the *ab initio* wave function can be accounted for by applying the modified generalized valence-bond (MGVB) method.

In addition to the *ab initio* calculations, many studies have been carried out using the extended Huckel and  $X\alpha$  scattered wave methods.<sup>1</sup> However, these methods are limited in their ability to describe accurately the multiply bonded metal dimers because of their MO formalisms and lack of correlation.

In order to determine accurate ground- and excited-state properties for quadruply bonded metal dimers, we have begun an *ab initio* study in which the ground and excited state wave functions are solved for self-consistently using the MGVB method. A detailed discussion of the results obtained from the *ab initio* study of the excited states for quadruply bonded metal dimers is presented elsewhere. In this chapter, the results of *ab initio* calculations of  $\text{Re}_2\text{Cl}_8^{2-}$  are reported.

### Computational Details

Hartree-Fock (HF), generalized valence bond (GVB), configuration interaction (CI), multiconfiguration self-consistent field (MCSCF), and modified generalized valence bond (MGVB) calculations were carried out on all metal dimers of interest.

The general form of the ground-state HF wave function for the quadruply bonded metal dimers is

$$\Psi_{\text{HF}} = A \Psi_{\text{core}} [\varphi_{a1g}(1)\alpha(1)] [\varphi_{a1g}(2)\beta(2)] [\varphi_{egx}(3)\alpha(3)] [\varphi_{egx}(4)\beta(4)] \\ [\varphi_{egy}(5)\alpha(5)] [\varphi_{egy}(6)\beta(6)] [\varphi_{b2g}(7)\alpha(7)] [\varphi_{b2g}(8)\beta(8)] , \quad 6.18$$

where  $A$  is the antisymmetrizer and  $\Psi_{\text{core}}$  represents all the ligand and nonvalence metal electron pairs. The overall symmetry of the systems studied is  $D_{4h}$ , and the orbitals represented above are those that are equivalent to the  $\sigma$ ,  $\pi$ , and  $\delta$  metal-metal bonding orbitals.

$$\varphi_{a1g} \approx \sigma \quad 6.19$$

$$\varphi_{eg} \approx \pi \quad 6.20$$

$$\varphi_{b2g} \approx \delta \quad 6.21$$

In the GVB wave function, unlike the HF wave function where an electron pair is described by a pair of identical orbitals  $\varphi_i(1)\varphi_i(2)$ , each electron in a bond pair is allowed to have a different orbital  $\varphi_i(1)\varphi_i'(2) + \varphi_i(2)\varphi_i'(1)$ , and every orbital is optimized self-consistently. For the metal dimers, the  $^1A_{1g}$  ground-state GVB-PP wave function has the form

$$\Psi_{\text{GVB}} = A \Psi_{\text{core}} [\varphi_{\sigma}(1)\varphi_{\sigma}'(2) + \varphi_{\sigma}'(1)\varphi_{\sigma}(2)] [\varphi_{\pi x}(3)\varphi_{\pi x}'(4) + \varphi_{\pi x}'(3)\varphi_{\pi x}(4)]$$

$$\left[ \varphi_{\pi y}(5) \varphi_{\pi y'}(6) + \varphi_{\pi y'}(5) \varphi_{\pi y}(6) \right] \left[ \varphi_{\delta}(7) \varphi_{\delta'}(8) + \varphi_{\delta'}(7) \varphi_{\delta}(8) \right] \chi \quad 6.22$$

$$\langle \varphi_i | \varphi_j \rangle = \delta_{ij}, \quad \langle \varphi_i | \varphi_i' \rangle = S_i, \quad 6.23$$

where, again, A is the antisymmetrizer,  $\psi_{\text{core}}$  represents the closed-shell HF core and  $\chi$  represents the spin function. This wave function describes the left-right static correlation in the two electron bonds. Where one electron is localized on the right, the other electron is more likely to be found on the left, thus relieving some of the electron-electron repulsion.

$$\varphi = I + \lambda r \quad 6.24$$

$$\varphi' = r + \lambda I \quad 6.25$$

In solving the GVB wave function, the GVB orbitals are cast in their orthogonal, natural orbital (NO) representation  $c_1 \varphi_{i1}(1) \varphi_{i1}(2) - c_2 \varphi_{i2}(1) \varphi_{i2}(2)$  (plus and minus combination of the one-electron orbitals), where  $\langle \varphi_{i1} | \varphi_{i2} \rangle = 0$  and  $c_1^2 + c_2^2 = 1$ .

For the quadruple metal-metal bond, the NO's would be bonding and antibonding orbitals of the appropriate bond pairs.

$$\Psi_{\text{GVB}}^{\text{NO}} = A \psi_{\text{core}} \left[ c_1 (\varphi_{\sigma})^2 - c_2 (\varphi_{\sigma}^*)^2 \right] \left[ c_3 (\varphi_{\pi x})^2 - c_4 (\varphi_{\pi x}^*)^2 \right] \left[ c_5 (\varphi_{\pi y})^2 - c_6 (\varphi_{\pi y}^*)^2 \right] \left[ c_7 (\varphi_{\delta})^2 - c_8 (\varphi_{\delta}^*)^2 \right] \chi \quad 6.26$$

The designation given this wave function is GVB-PP(4/8), four bonds correlated with eight natural orbitals.



In the GVB-PP wave function, only a single spin-coupling (or valence bond structure) is allowed and the wave function is designated PP for perfect pairing. In the NO representation, this spin restriction leads only to configurations where the occupation for any given bond pair is {2,0} or {0,2}.

The GVB-PP wave function often provides an excellent description of saturated organic systems; however, for transition metal systems, singly occupied d orbitals on the same metal prefer to be high-spin coupled, and it is necessary to optimize the spin coupling of the GVB wave function.

Originally, the full GVB wave function was calculated using the form

$$A\phi_a\phi_b\cdots\phi_z\chi, \quad 6.27$$

where the orbitals  $\phi_a\cdots\phi_z$  are all allowed to overlap and where the spin function  $\chi$  is optimized.<sup>6</sup> Such calculations are too tedious for studies of large molecules, and a series of studies led to an abbreviated approach in which the GVB-PP wave function of (6.26) is generalized to allow the {1,1} configuration in each bond pair. The orbitals and CI coefficients of the resulting wave function are solved for using the GVB3 program.<sup>4,7</sup> This wave function does not correspond exactly to the full GVB wave function; however, it does allow a completely general description of the spin coupling and an unambiguous description of the two electrons in each bond pair, just as in GVB. The resulting wave function is denoted GVB-RCI in order to indicate the restricted coupling used.

In nontransition metal systems, the full GVB wave function can often be well approximated by using the GVB-PP orbital in the GVB-RCI wave function. However, for transition metal system, the orbitals in the GVB-RCI wave function must be optimized self-consistently. Calculations in which all the GVB occupied orbitals are used and general occupations are allowed are denoted as GVB-CI. This allows a more complete description of the resonance and the correlation effects.

Recently, it has been shown from studies on  $\text{Cr}_2$  and  $\text{Mo}_2$  that quite significant correlation errors result from difficulties in describing the atomic electron affinities.<sup>4,8</sup> With this insight, a simple modification to the GVB method was made that corrects certain Coulombic integrals so as to provide the correct atomic electron affinities.<sup>4,5</sup> This method, MGVB, when applied to the metal dimers gave accurate bond lengths and bond energies.

### Modified Generalized Valence Bond Method

The MGVB method was developed by Goodgame and Goddard to account for the electron correlation missing in a GVB wave function.<sup>4,5</sup> Because of the importance of this method to the present investigation, a brief discussion is necessary.

In the long  $R$  limit, the homonuclear GVB wave function is covalent and has the form

$$\Psi_{\text{GVB}} = [\varphi_l(1)\varphi_r(2) + \varphi_r(1)\varphi_l(2)](\alpha\beta - \beta\alpha), \quad 6.28$$

where  $\varphi_l$  and  $\varphi_r$  represent the one-electron orbitals localized on the left and right centers, respectively. As the internuclear distance is decreased, the one-electron orbitals delocalize and the GVB wave function has the form

$$\Psi_{\text{GVB}} = [\varphi_a(1)\varphi_b(2) + \varphi_b(1)\varphi_a(2)](\alpha\beta - \beta\alpha), \quad 6.29$$

where the GVB orbitals can be thought of as consisting of optimized atomic orbitals (ignoring normalization),

$$\varphi_a = \varphi_l + \lambda\varphi_r \quad \text{and} \quad \varphi_b = \varphi_r + \lambda\varphi_l. \quad 6.30$$

Expanding the GVB wave function in terms of these optimized atomic functions,

$$\Psi_{\text{GVB}} = (\varphi_l \varphi_r + \varphi_r \varphi_l) + \lambda (\varphi_l \varphi_l + \varphi_r \varphi_r) , \quad 6.31$$

leads to a superposition of an ionic term (two electrons localized on the same center) and a covalent term (two electrons localized on different centers).

$$\Psi_{\text{GVB}} = C_c \Psi_{\text{cov}} + C_i \Psi_{\text{ion}} \quad 6.32$$

Goodgame and Goddard postulated that the GVB wave function leads to an excellent description of the covalent terms, but the ionic terms, involving as they do doubly occupied orbitals, should be too high by  $\sim 1.3$  eV.<sup>5</sup> They argued that the usual CI approach, in which one does double excitations out of each bond pair, is specifically correcting for the error in the ionic terms. They tested this hypothesis by carrying out calculations in which the one-center Coulomb integrals were modified so to have the effect of properly correcting the ionic terms at  $R = \infty$ .

There are several ways in which corrections for the error in the ionic terms can be determined. The simplest, used by Goodgame and Goddard in their first paper and referred to as MGVB/1,<sup>5</sup> corrects the limits so that the wave function  $ll$  or  $rr$  at  $R = \infty$  has the correct energy relative to  $lr$  or  $rl$ , leading to

$$CC = IP(\text{Exp} - \text{Calc}) - EA(\text{Exp} - \text{Calc}) . \quad 6.33$$

For the hextuply bonded metal dimers, the total error in the HF energy was

$$\delta E(s^0 d^5) + \delta E(s^2 d^5) - 2\delta E(s^1 d^5) = \delta IP_s - \delta EA_s \quad 6.34$$

$$\delta E(s^1 d^4) + \delta E(s^1 d^6) - 2\delta E(s^1 d^5) = \delta IP_d - \delta EA_d \quad 6.35$$

These correction values were then used to modify the s-s and d-d self-Coulomb terms while all other integrals were evaluated exactly.

In a new approach, MGVB/2, various one-center, one- and two-electron integrals are adjusted to fit the IP's, EA's and atomic-excitation energies. In MGVB/2, an energy expression based on the Condon and Shortley formalism is derived for various atom states. While the values of the one- and two-electron integrals in these energy expressions are dependent on the electron configuration of the atom, it has been shown that the correlation error expressed in terms of these integrals is relatively constant for a given ionic series and shows only minor fluctuation with changing ionicism. In other words, the correlation error in the HF energy appears to be constant.

MGVB/2 is an extension of the ideas developed in the earlier method. The corrections to the one- and two-electron Condon and Shortley parameters, obtained by fitting the difference between the calculated and experimental term energies with the parametric equations, are added to the appropriate integrals, while evaluating all integrals exactly. Now, instead of correcting only for the correlation error associated with placing two electrons in the same orbital, a spherical correction is applied that accounts not only for the correlation between electrons in the same orbital but also for the correlation of electrons in different orbitals.

### **Method of Correction**

In applying the MGVB corrections, the following procedure is used:

- (1) symmetric orthogonalization of the localized GVB orbitals;
- (2) evaluate all integrals in terms of the localized orbitals and apply the appropriate corrections to the one-center integrals;
- (3) transform these orbitals back to the GVB natural orbitals for use in a CI or an MCSCF calculation;

- (4) recalculate the usual GVB-RCI wave function, using the modified integrals;
- (5) in calculating the CI wave function of (4), the corrected integrals are used only for bond pairs in which bonding or antibonding GVB natural orbitals are doubly occupied, since only in this case is there static ionic character in the wave function.

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## Bond Energy and Other Properties of the Re-Re Quadruple Bond

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**Abstract:** Using generalized valence bond (GVB) methods designed for obtaining accurate bond energies, we predict an Re-Re quadruple bond strength of  $85 \pm 5$  kcal/mol for  $\text{Re}_2\text{Cl}_8^{2-}$ . This is much less than early estimates of 370 kcal/mol and somewhat lower than estimates (124 to 150 kcal/mol) based on Birge-Sponer extrapolation but is in reasonable agreement with a recent thermochemical study ( $97 \pm 12$  kcal/mol). We obtain a rotational barrier of 3.0 kcal/mol and a singlet-triplet excitation energy of  $3100\text{ cm}^{-1}$ , and we conclude that the intrinsic strength of the  $\delta$  bond is  $6 \pm 3$  kcal/mol.

Since their discovery in 1965, quadruply bonded metal dimers have provoked numerous theoretical and experimental studies. A particularly controversial issue has been the strength of the quadruple bond and, in particular, the contribution of the  $\delta$  bond to the observed structure of the unbridged dimers.<sup>1-4</sup> We here report the results of ab initio calculations of  $\text{Re}_2\text{Cl}_8^{2-}$  designed to provide accurate bond energies and torsion barriers as well as accurate shapes for the potential curves. These studies use the generalized valence bond (GVB) approach in which electron correlations are included for all eight electrons available for the

quadruple bond, while solving self-consistently for all orbitals.<sup>5,6</sup>

We use the modified-GVB (M-GVB) approach of Goodgame and Goddard.<sup>7</sup> They pointed out that ab initio descriptions of multiple bonds in transition metals lead to substantial errors in the bond energy due to an inadequate treatment of the electron correlations in the ionic part of the wave function describing the bond.

$$\psi^{\text{GVB}} = \psi_{\text{cov}} + \lambda \psi_{\text{ionic}} = [\phi_1(1)\phi_r(2) + \phi_r(1)\phi_1(2)] + \lambda[\phi_1(1)\phi_1(2) + \phi_r(1)\phi_r(2)]$$

In GVB, electron correlation in the covalent part of the wave

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**Table II.** Comparison of Calculated Spectroscopic Constants for the  $^1A_{1g}$  and  $^3A_{2g}$  States at the M-GVB-RCI Level

	$^1A_{1g}$	$^3A_{2g}$ <sup>a</sup>	diff
bond length (Å)	2.26	2.30	+0.04
Re-Re force constant (mdyn/Å)	4.65	4.67	+0.02
vibrational frequency (cm <sup>-1</sup> )	293	293	0.0
bond energy (kcal/mol)	85.0	76.5	-8.5

<sup>a</sup>Optimization of eclipsed rotomer neglecting relaxation of ReCl<sub>4</sub> fragments from their ground-state structure.

neutralized by counter charges (e.g., K<sup>+</sup>), and in calculating bond energies it is necessary to include the effects of these counter charges (the Coulomb energy for two charges at 2.24 Å is 150 kcal/mol!). In order to ameliorate this problem, fractional charges (leading to a net charge of 0 on each fragment) are placed at positions extrapolated from the crystallographic positions of the counterions. The charges are placed so as to maintain the overall symmetry of the system and to be invariant under a 45° rotation of the ReCl<sub>4</sub><sup>-</sup> units.<sup>10</sup> The rotational invariance allows for a comparison of the properties of the eclipsed and staggered geometries. The results for the Re-Re bond distance optimization are shown in Figure 3 and in Tables I and II.<sup>11</sup>

The bond energy for the Re-Re quadruple bond calculated at the M-GVB level is 85.0 kcal/mol, the first ab initio estimate of the energy of a quadruple metal-metal bond. Similar calculations on Mo<sub>2</sub> and Cr<sub>2</sub> lead to bond energies within 5 kcal/mol of experiment, and we believe that similar accuracy can be expected in the Re-Re quadruple bond studies, leading to  $D_e$  (Re<sup>4</sup>-Re) = 85 ± 5 kcal/mol. A more conservative estimate of the uncertainty, perhaps, ±10 kcal/mol, would be warranted given the lack of comparable theoretical and experimental studies on other systems. Experimental values for the bond energy of the quadruple bond have been quite difficult to obtain. Early estimates have ranged as high as 370 kcal/mol.<sup>1</sup> Birge-Sponer extrapolations using the harmonic stretching frequency and the anharmonicity constant determined from resonance Raman measurements suggest an Re-Re bond energy of 152 ± 19 kcal/mol for Re<sub>2</sub>Cl<sub>8</sub><sup>2-</sup> and 139 ± 24 kcal/mol for Re<sub>2</sub>Br<sub>8</sub><sup>2-</sup>.<sup>12</sup> Such Birge-Sponer extrapolations for multiple bonds are fraught with peril and could easily lead to errors of 50–70 kcal/mol.<sup>13</sup> A more reliable thermochemical study places the Re-Re bond energy for Cs<sub>2</sub>Re<sub>2</sub>Br<sub>8</sub> at 97 ± 12 kcal/mol,<sup>14</sup> this estimate depending on empirically based assumptions in order to estimate the Re-Br bond energy. Given the various experimental uncertainties, we believe that the theoretical value of 85 kcal/mol is the best current estimate of the bond energy. Early extended Hückel calculations placed the quadruple bond in Re<sub>2</sub>Cl<sub>8</sub><sup>2-</sup> at 370 kcal/mol.<sup>1</sup> Recent theoretical work with the Hartree-Fock-Slater transition-state method has estimated the Re-Re triple bond energy in Re<sub>2</sub>Cl<sub>4</sub>·(PH<sub>3</sub>)<sub>4</sub> to be 134 kcal/mol.<sup>15</sup> Generally, these latter methods

lead to an overestimate of the bond energy by 1 to 2 eV,<sup>16</sup> suggesting  $D_e$  ≈ 90 to 110 kcal/mol, in line with the GVB results.

Given a total bond energy of 85 kcal/mol, the question is how strong is the δ-bond? One way to establish this is by the singlet-triplet gap at the eclipsed geometry (calculated at 3100 cm<sup>-1</sup> = 0.38 eV = 8.8 kcal/mol), suggesting a δ bond strength of 4.5–9.0 kcal/mol (this estimate ignores spin-coupling with the other orbitals of the bond). A second approach is to use the rotational barrier at the ground-state equilibrium bond distance. The energy difference between the eclipsed and staggered geometries involves opposing energy contributions. Rotating from eclipsed to staggered changes the overlap, i.e., the δ bond goes to zero so that bonding is lost. Simultaneously electrostatic and steric interactions between Cl's bonded to opposite Re's are relieved. With this competitive effect in the total energy, one can only obtain bounds to the pure δ contribution to the bonding. The calculations give a direct barrier of 3.0 kcal/mol,<sup>17a</sup> a lower bound on the strength of the δ bond. The calculated barrier for the triplet state (favoring staggered) of 3.0 kcal/mol<sup>17b</sup> should be an upper limit on the steric barrier, suggesting an upper limit on the δ bond of 6.0 kcal/mol. From the above analysis, a reasonable estimate of the δ bond strength is 6 ± 3 kcal/mol, which is weak compared with σ and π bonds but sufficiently large to explain the ubiquity of the eclipsed geometry for d<sup>4</sup>-d<sup>4</sup> metal dimers.

There are not experimental estimates of the δ bond strength; however, from dynamic NMR studies of meso-substituted molybdenum porphyrin dimers, the activation energy for rotation about the quadruple metal-metal bond has been estimated to be 10.1 ± 0.5 kcal/mol.<sup>18</sup> This is larger than the value calculated for Re<sub>2</sub>Cl<sub>8</sub><sup>2-</sup>, perhaps because steric interactions (which reduce the barrier) are smaller in the porphyrin. From measured δ → δ\* splittings in odd-electron complexes [e.g., Mo<sub>2</sub>(SO<sub>4</sub>)<sub>4</sub><sup>3-</sup> and Tc<sub>2</sub>Cl<sub>8</sub><sup>3-</sup>], Trogler and Gray have estimated the δ bond energy to be 9–10 kcal/mol.<sup>19</sup>

Early self-consistent-field (SCF) calculations found the ground state of the staggered conformation to be a triplet, about 60 kcal/mol lower than the singlet ground state of the eclipsed geometry. Addition of CI reduced the difference between eclipsed and staggered geometries to 4 kcal/mol, yet still favored the triplet state of the staggered conformation.<sup>20</sup> Other recent ab initio calculations find the ground state for the eclipsed and staggered geometries to be a singlet with essentially no barrier between the two rotomers.<sup>21</sup>

The M-GVB calculations lead to an Re-Re bond distance of 2.26 Å, in excellent agreement with the experimental value of 2.24 Å. This is consistent with the results on Mo<sub>2</sub> where M-GVB leads to an error of -0.01 Å.<sup>7</sup> We calculate a metal-metal stretching force constant of 4.6 mdyn/Å, and using a valence force field,<sup>22</sup> we obtain a metal-metal stretching frequency of 293 cm<sup>-1</sup>.<sup>23</sup> This can be compared with the vibrational frequency for (n-Bu<sub>4</sub>N)<sub>2</sub>[Re<sub>2</sub>Cl<sub>8</sub>] (determined from resonance Raman) of 275 cm<sup>-1</sup>.<sup>24</sup>

Summarizing, our calculations (i) provide the first prediction based on ab initio studies of an eclipsed ground state for a

(10) An average metal-ion distance of 4.3 Å was maintained (determined from the X-ray diffraction data). Charges of +0.125 were placed in the same plane as defined by the four Cl's of each ReCl<sub>4</sub> unit. The charges were placed at 22.5° off the Re-Cl vector. The validity of this approach is supported by the small effect on excitation energies and rotational barriers where use of the counterions (i) increases the <sup>3</sup>(δ-δ\*) excitation energy by 10 cm<sup>-1</sup> for the eclipsed conformation and less than 10 cm<sup>-1</sup> for the staggered conformation and (ii) decreases the rotational barrier by 0.5 kcal/mol.

(11) The calculations lead to a slight barrier in the potential curves as the neutral ReCl<sub>4</sub> fragments are brought together. This barrier is 0.45 kcal/mol (at 4.18 Å) for GVB-RCI and 0.13 kcal/mol (at 4.92 Å) for M-GVB. This occurs because the nonplanar ReCl<sub>4</sub> fragments are kept at a fixed nonplanar geometry as the fragments are separated, leading to a small dipole moment and hence a slight repulsive interaction at large R.

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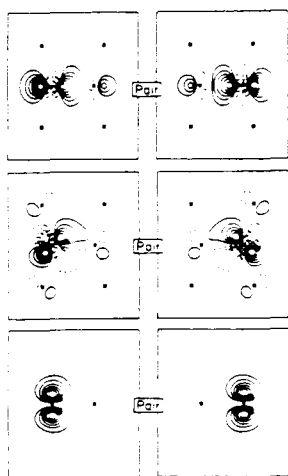


Figure 1. GVB orbitals involved in the  $\sigma$ ,  $\pi$ , and  $\delta$  bonds of  $\text{Re}_2\text{Cl}_4^{2-}$  (reading top to bottom). Spacing between contours is 0.05 au. Negative contours are denoted by dashed lines; asterisks denote position of nuclei. The  $\sigma$  and  $\pi$  orbitals are plotted in the  $xz$  plane; the  $\delta$  orbitals are plotted in a plane rotated  $45^\circ$  from the  $xz$  plane.

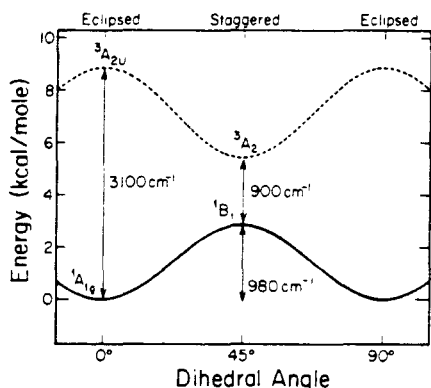


Figure 2. Energy profile for the ground state and  $^3(\delta\delta^*)$  excited state of  $\text{Re}_2\text{Cl}_4^{2-}$  as a function of dihedral angle. The  $^3(\delta\delta^*)$  excitation energy was calculated for  $\phi$  of  $0^\circ$  and  $45^\circ$ . The plot assumes a periodic energy function with the energies at  $\phi$  of  $0^\circ$  and  $45^\circ$  as the minimum and maximum of the function for the ground state and the maximum and minimum of the function for the  $^3(\delta\delta^*)$  excited state.

function is well described, but the ionic terms are forced to use doubly occupied orbitals. Because of the spatial compactness of the d orbitals, there are substantial electron correlations in such doubly occupied orbitals, leading to large errors in certain atomic electron affinities and ionization potentials. In M-GVB, this correlation error of the atom is built into the atomic Coulomb integrals, but calculations are otherwise as in normal GVB.<sup>7</sup> The electron correlations imbedded in M-GVB would have been included in normal GVB wave functions by suitably high-level excitations (CI) and hence with M-GVB we cannot go beyond GVB-level calculations (since some electron correlation effects would be double corrected). M-GVB has been successfully applied in studying the sextuple bonds of  $\text{Mo}_2$  and  $\text{Cr}_2$ , where an accurate description of bond energies ( $D_0$ ) and bond distances ( $R_e$ ) was obtained.<sup>7</sup> Thus, for  $\text{Mo}_2$ ,  $D_0^{\text{calcd}} = 90.9$  kcal/mol,  $D_0^{\text{exptl}} = 97 \pm 5$  kcal/mol,  $R_e^{\text{calcd}} = 1.92$  Å,  $R_e^{\text{exptl}} = 1.93$  Å; and for  $\text{Cr}_2$ ,  $D_0^{\text{calcd}} = 42.9$  kcal/mol,  $D_0^{\text{exptl}} = 46 \pm 7$  kcal/mol,  $R_e^{\text{calcd}} = 1.61$  Å,  $R_e^{\text{exptl}} = 1.68$  Å. The current studies involve exactly the same approach as in these earlier studies except that we now utilize effective core

Table I. Comparison of Calculated and Experimental Spectroscopic Constants

	GVB-PP <sup>a</sup>	GVB-RCI <sup>b</sup>	M-GVB-RCI	exptl
bond length (Å)	2.37	2.36	2.26	2.24 <sup>c</sup>
Re-Re force constant (mdyn/Å)	2.75	2.48	4.64	
vibrational frequency <sup>d</sup> ( $\text{cm}^{-1}$ )	247	240	293	275 <sup>e</sup>
bond energy (kcal/mol)	g	25.8	85.0	$(97 \pm 12)^f$
rotational barrier (kcal/mol)	-1.6	0.1	2.8	

<sup>a</sup> For a description of GVB-PP wave functions, see ref 6. <sup>b</sup> CI wave function similar to that described in Table Ia of Moss, B. J.; Goddard, W. A., III *J. Chem. Phys.* 1975, 63, 3523-3531. <sup>c</sup> Based on X-ray diffraction study of  $\text{K}_4[\text{Re}_2\text{Cl}_4] \cdot 2\text{H}_2\text{O}$ ; ref 9. <sup>d</sup> The calculated Re-Re bond length and force constants were used with an earlier determined valence force field in the vibrational analysis; ref 22. <sup>e</sup> Based on solid-state resonance Raman study of  $(n\text{-Bu}_4\text{N})_2[\text{Re}_2\text{Cl}_4]$ ; ref 24. <sup>f</sup> Based on thermochemical study of  $\text{Cs}_2\text{Re}_2\text{Br}_8$ ; ref 14. <sup>g</sup> Keeping a constant spin coupling, GVB-PP leads to a dissociation error of 55.5 kcal/mol.

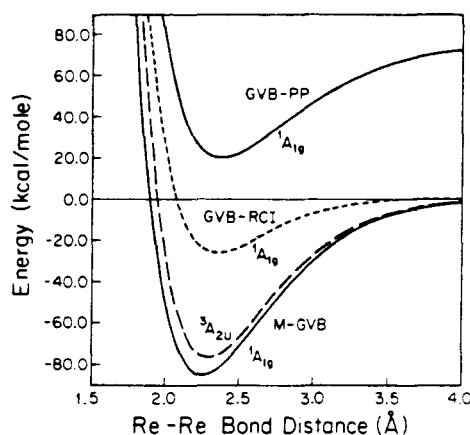


Figure 3. Calculated potential curves for  $\text{Re}_2\text{Cl}_4^{2-}$ .

potentials (ECP) so that only the 15 outermost electrons (5s,5p,6s,5d) of each Re are treated explicitly; however, the potentials describe the effects of core electrons including the dominant relativistic effects.<sup>8</sup>

The GVB orbitals of the singlet ground state ( $^1A_{1g}$ ) are shown in Figure 1 where we see that four electrons in  $d\sigma$ ,  $d\pi_x$ ,  $d\pi_y$ , and  $d\delta_z$  orbitals on each Re are spin-paired to form the quadruple bond. Uncoupling the spins of the two electrons in  $\delta$  orbitals leads to the lowest triplet state ( $^3A_{2u}$ ). Rotating one  $\text{ReCl}_4$  group about the bond by  $45^\circ$  leads (see Figure 2) to essentially identical bonding in the  $\sigma$ ,  $\pi_x$ , and  $\pi_y$  orbitals, but the  $\delta$  bond is broken with the result that the singlet and triplet states (denoted as  $^1B_1$  and  $^3A_2$ ) are nearly degenerate. Some indications of the relative strengths of these bonds are given by the overlaps in the GVB orbitals,  $S_\sigma = 0.68$ ,  $S_\pi = 0.54$ ,  $S_\delta = 0.12$  for the eclipsed configurations, but  $S_\sigma = 0.68$ ,  $S_\pi = 0.53$ , and  $S_\delta = 0.0$  for the staggered. As the bond is stretched, the overlaps of the GVB orbitals decrease continuously to zero, leading to a smooth description of bond dissociation (molecular orbital based schemes often lead to the wrong dissociation limit, complicating predictions of bond energies).

In these calculations the geometry of each  $(\text{ReCl}_4)^-$  fragment was fixed as that in the crystal,<sup>9</sup> and the Re-Re bond distance was optimized. An important issue here for bond energies concerns the charges. In crystals and solution, the charge of  $\text{Re}_2\text{Cl}_4^{2-}$  is

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quadruply bonded metal dimer, (ii) predict an Re—Re quadruple bond energy of  $85 \pm 5$  kcal/mol—the first direct estimate of the strength of this prototypical quadruple bond, and (iii) suggest that the  $\delta$  bond energy is  $6 \pm 3$  kcal/mol.

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## **Appendix 1**

Table A1.1. Hydrogen Coordinates  $\times 10^4$ .

Atom	$x$	$y$	$z$	$B$
H3A	-2119	180	-103	5.8
H3B	-1996	521	555	5.8
H3C	-2813	523	-303	5.8
H4A	-1871	319	-1353	4.8
H4B	-2547	673	-1669	4.8
H4C	-1595	752	-1522	4.8
H8A	-1515	2369	356	5.7
H8B	-1812	2047	-363	5.7
H8C	-2378	2131	112	5.7
H9A	-986	2147	1873	5.0
H9B	-1790	1874	1707	5.0
H9C	-864	1685	2121	5.0
H13A	2745	325	-970	8.1
H13B	2647	621	-1696	8.1
H13C	2912	792	-797	8.1
H14A	1224	39	-1778	8.3
H14B	441	339	-2060	8.3
H14C	1011	326	-2547	8.3
H18A	332	2219	-2078	5.1
H18B	1084	1936	-2034	5.1
H18C	194	1926	-2820	5.1
H19A	-883	1960	-1866	7.2
H19B	-1164	1625	-2559	7.2
H19C	-941	1503	-1637	7.2
H23A	4104	207	3565	8.0
H23B	4925	484	3896	8.0
H23C	4016	678	3652	8.0
H24A	4288	-85	2344	6.9
H24B	4411	193	1686	6.9
H24C	5165	140	2578	6.9
H28A	4675	2095	1755	7.1
H28B	4872	1791	2495	7.1
H28C	5436	1784	2001	7.1
H29A	3955	1841	224	7.6
H29B	4623	1493	348	7.6
H29C	3644	1388	21	7.6
H33A	936	88	3735	5.0
H33B	1801	332	4133	5.0
H33C	1127	372	4502	5.0
H34A	-455	459	2873	6.4
H34B	-337	780	3565	6.4
H34C	-446	925	2685	6.4
H38A	3641	1845	4263	7.0
H38B	3673	1527	4934	7.0
H38C	3637	1377	4083	7.0
H39A	2252	2193	4120	7.7
H39B	1424	1931	3908	7.7
H39C	2177	1920	4807	7.7

Atom	$x$	$y$	$z$	$B$
H5A	-2417	1285	-686	4.2
H5B	-1439	1357	-444	4.2
H6A	-1306	1133	975	3.9
H6B	-2234	1319	625	3.9
H15A	1380	1128	-2268	5.3
H15B	1628	1296	-1371	5.3
H16A	165	987	-1679	5.7
H16B	10	1116	-2581	5.7
H25A	5176	971	2871	5.3
H25B	4275	1168	2655	5.3
H26A	4004	882	1194	6.0
H26B	4985	988	1540	6.0
H35A	997	1184	4224	4.8
H35B	897	1327	3347	4.8
H36A	2330	948	3835	4.8
H36B	2400	1092	4701	4.8
H42	8254	1384	3283	5.5
H43	7720	1077	1909	6.1
H44	7081	433	1659	6.1
H45	6805	98	2738	7.3
H46	7212	436	4039	5.1
H52	8697	1716	5999	6.1
H53	8315	2421	6078	7.7
H54	7344	2746	4833	7.3
H55	6701	2400	3560	6.9
H56	7119	1694	3487	5.2
H62	6513	1072	4514	5.7
H63	5826	872	5378	7.5
H64	6648	487	6607	6.0
H65	8137	412	7117	6.0
H66	8855	689	6305	4.7
H72	9714	1659	5405	5.9
H73	11257	1468	6042	7.0
H74	11618	785	6216	6.2
H75	10631	253	5743	6.6
H76	9085	430	5053	4.5
H82	9803	2838	2336	4.6
H83	10994	2430	2387	8.4
H84	12280	2730	2494	6.4
H85	12412	3442	2612	5.8
H86	11293	3862	2722	4.5
H92	8885	3490	3639	5.2
H93	7692	3067	3471	5.3
H94	6970	2667	2254	5.1
H95	7306	2774	1107	5.2
H96	8478	3215	1239	4.5
H102	10825	3439	4138	4.8
H103	11213	3758	5507	6.1

Atom	$x$	$y$	$z$	$B$
H104	10802	4415	5560	6.9
H105	9846	4803	4381	9.0
H106	9392	4470	2984	5.8
H112	10233	4000	1325	3.8
H113	9682	4530	340	5.0
H114	8422	4880	126	4.3
H115	7631	4681	892	4.7
H116	8137	4128	1849	3.4

**Table A1.2.** Anisotropic Thermal Parameters  $\times 10^4$ .

Atom	$U_{11}$	$U_{22}$	$U_{33}$	$U_{12}$	$U_{13}$	$U_{23}$
IR1	374(5)	410(5)	229(5)	58(4)	128(4)	-8(4)
IR2	384(5)	388(5)	236(5)	28(4)	150(4)	18(4)
C3	584(126)	472(140)	1055(173)	-10(105)	437(125)	136(124)
C4	656(130)	571(140)	435(112)	116(110)	46(96)	-26(103)
C8	698(126)	444(130)	983(164)	308(103)	319(119)	124(117)
C9	926(148)	466(129)	659(132)	-9(108)	438(116)	-139(104)
C13	1439(165)	939(199)	904(180)	346(150)	880(145)	6(154)
C14	2151(217)	687(149)	601(148)	242(150)	609(148)	134(126)
C18	628(133)	638(146)	458(131)	-137(112)	182(110)	116(112)
C19	562(119)	1413(193)	745(169)	138(125)	398(116)	555(149)
C23	693(143)	1553(214)	551(133)	-30(138)	-92(110)	118(132)
C24	604(135)	664(148)	1007(172)	332(113)	119(125)	-33(126)
C28	534(126)	896(163)	1462(182)	-139(116)	359(126)	80(140)
C29	1175(175)	1143(177)	932(165)	157(144)	772(144)	212(141)
C33	837(136)	745(129)	489(131)	-58(107)	340(110)	32(110)
C34	694(122)	1118(189)	784(146)	-72(120)	505(109)	58(135)
C38	711(124)	1161(184)	611(149)	-301(123)	118(110)	102(139)
C39	1343(169)	861(164)	590(167)	160(133)	164(138)	-45(131)

Table A1.3. More Distances and Angles.

		Distance(Å)			Distance(Å)
N1	-C2	1.47(2)	C52	-C53	1.41(3)
C2	-C3	1.51(3)	C53	-C54	1.37(3)
C2	-C4	1.52(2)	C54	-C55	1.40(3)
C2	-C5	1.53(2)	C55	-C56	1.40(3)
C5	-C6	1.49(2)	C61	-C62	1.38(3)
C6	-C7	1.53(2)	C61	-C66	1.41(2)
C7	-C8	1.53(3)	C62	-C63	1.40(3)
C7	-C9	1.56(2)	C63	-C64	1.39(3)
C7	-N2	1.48(2)	C64	-C65	1.37(3)
N3	-C12	1.50(2)	C65	-C66	1.42(3)
C12	-C13	1.46(3)	C71	-C72	1.41(3)
C12	-C14	1.55(3)	C71	-C76	1.38(2)
C12	-C15	1.54(3)	C72	-C73	1.43(3)
C15	-C16	1.48(3)	C73	-C74	1.31(3)
C16	-C17	1.51(3)	C74	-C75	1.34(3)
C17	-C18	1.50(3)	C75	-C76	1.41(3)
C17	-C19	1.50(3)	B2	-C81	1.64(3)
C17	-N4	1.47(2)	B2	-C91	1.62(3)
N5	-C22	1.51(2)	B2	-C101	1.63(3)
C22	-C23	1.56(3)	B2	-C111	1.65(2)
C22	-C24	1.51(3)	C81	-C82	1.35(2)
C22	-C25	1.50(3)	C81	-C86	1.38(2)
C25	-C26	1.55(3)	C82	-C83	1.41(3)
C26	-C27	1.48(3)	C83	-C84	1.35(3)
C27	-C28	1.55(3)	C84	-C85	1.34(3)
C27	-C29	1.54(3)	C85	-C86	1.43(3)
C27	-N6	1.48(2)	C91	-C92	1.40(2)
N7	-C32	1.48(2)	C91	-C96	1.39(2)
C32	-C33	1.51(2)	C92	-C93	1.39(3)
C32	-C34	1.47(3)	C93	-C94	1.37(3)
C32	-C35	1.51(2)	C94	-C95	1.36(3)
C35	-C36	1.50(3)	C95	-C96	1.38(3)
C36	-C37	1.53(3)	C101	-C102	1.41(2)
C37	-C38	1.50(3)	C101	-C106	1.36(3)
C38	-C39	2.53(3)	C102	-C103	1.43(3)
C37	-C39	1.50(3)	C103	-C104	1.33(3)
C37	-N8	1.50(3)	C104	-C105	1.36(3)
B1	-C41	1.63(3)	C105	-C106	1.47(3)
B1	-C51	1.59(3)	C111	-C112	1.41(2)
B1	-C61	1.64(3)	C111	-C116	1.39(2)
B1	-C71	1.65(3)	C112	-C113	1.39(3)
C41	-C42	1.41(3)	C113	-C114	1.34(3)
C41	-C46	1.38(3)	C114	-C115	1.38(2)
C42	-C43	1.42(3)	C115	-C116	1.40(2)
C43	-C44	1.34(3)	N9	-C200	1.26(5)
C44	-C45	1.39(3)	C200	-C201	1.33(5)
C45	-C46	1.41(3)	N10	-C203	1.22(5)
C51	-C52	1.41(3)	C203	-C204	1.25(6)
C51	-C56	1.40(3)			



Angle(°)			Angle(°)		
C1 -IR1 -C11	92.2	(6)	C23 -C22 -C25	108.4	(15)
C1 -IR1 -C21	175.5	(7)	C24 -C22 -C25	116.3	(16)
C1 -IR1 -C31	85.9	(6)	C22 -C25 -C26	115.5	(15)
C11 -IR1 -C21	90.1	(7)	C25 -C26 -C27	116.4	(16)
C11 -IR1 -C31	176.8	(6)	C26 -C27 -C28	112.7	(15)
C21 -IR1 -C31	91.6	(7)	C26 -C27 -C29	110.6	(16)
C10 -IR2 -C20	87.5	(7)	C28 -C27 -C29	110.2	(15)
C10 -IR2 -C30	175.4	(7)	C26 -C27 -N6	108.1	(15)
C10 -IR2 -C40	88.5	(7)	C28 -C27 -N6	108.0	(14)
C20 -IR2 -C30	92.4	(7)	C29 -C27 -N6	107.0	(15)
C20 -IR2 -C40	175.3	(7)	C27 -N6 -C30	172.4	(16)
C30 -IR2 -C40	91.3	(7)	C31 -N7 -C32	173.0	(16)
C1 -N1 -C2	169.5	(17)	N7 -C32 -C33	107.4	(13)
N1 -C2 -C3	108.9	(14)	N7 -C32 -C34	108.6	(14)
N1 -C2 -C4	106.6	(13)	C33 -C32 -C34	111.1	(14)
C3 -C2 -C4	108.6	(14)	N7 -C32 -C35	107.3	(13)
N1 -C2 -C5	106.4	(13)	C33 -C32 -C35	112.6	(14)
C3 -C2 -C5	112.4	(14)	C34 -C32 -C35	109.6	(14)
C4 -C2 -C5	113.8	(14)	C32 -C35 -C36	117.1	(15)
C2 -C5 -C6	116.8	(14)	C35 -C36 -C37	117.1	(15)
C5 -C6 -C7	115.4	(14)	C36 -C37 -C38	111.8	(16)
C6 -C7 -C8	113.4	(14)	C36 -C37 -C39	110.8	(16)
C6 -C7 -C9	109.9	(13)	C38 -C37 -C39	115.1	(17)
C8 -C7 -C9	108.4	(14)	C36 -C37 -N8	106.7	(15)
C6 -C7 -N2	108.9	(13)	C38 -C37 -N8	105.6	(15)
C8 -C7 -N2	108.5	(13)	C39 -C37 -N8	106.3	(15)
C9 -C7 -N2	107.6	(13)	C37 -N8 -C40	172.5	(18)
C7 -N2 -C10	172.5	(16)	C41 -B1 -C51	112.1	(14)
C11 -N3 -C12	174.5	(17)	C41 -B1 -C61	110.0	(14)
N3 -C12 -C13	109.0	(16)	C41 -B1 -C71	105.5	(14)
N3 -C12 -C14	106.7	(15)	C51 -B1 -C61	105.1	(14)
C13 -C12 -C14	110.2	(17)	C51 -B1 -C71	113.5	(14)
N3 -C12 -C15	107.2	(15)	C61 -B1 -C71	110.8	(14)
C13 -C12 -C15	113.8	(16)	B1 -C41 -C42	121.1	(15)
C14 -C12 -C15	109.7	(16)	B1 -C41 -C46	124.7	(15)
C12 -C15 -C16	118.0	(16)	C46 -C41 -C42	113.9	(16)
C15 -C16 -C17	116.2	(16)	C41 -C42 -C43	122.5	(17)
C16 -C17 -C18	115.1	(16)	C42 -C43 -C44	120.3	(18)
C16 -C17 -C19	112.5	(16)	C43 -C44 -C45	119.9	(19)
C18 -C17 -C19	110.6	(16)	C44 -C45 -C46	118.5	(19)
C16 -C17 -N4	106.4	(15)	C45 -C46 -C41	124.8	(17)
C18 -C17 -N4	107.1	(15)	B1 -C51 -C52	119.6	(15)
C19 -C17 -N4	104.3	(15)	B1 -C51 -C56	126.2	(15)
C17 -N4 -C20	170.2	(18)	C56 -C51 -C52	113.9	(16)
C21 -N5 -C22	176.6	(17)	C51 -C52 -C53	125.0	(18)
N5 -C22 -C23	105.1	(15)	C52 -C53 -C54	117.2	(20)
N5 -C22 -C24	108.6	(15)	C53 -C54 -C55	121.5	(20)
C23 -C22 -C24	110.4	(16)	C54 -C55 -C56	118.6	(19)
N5 -C22 -C25	107.5	(15)	C55 -C56 -C51	123.7	(17)

Angle(°)				Angle(°)			
B1	-C61	-C62	119.1(15)	C111-C112-C113	121.5(16)		
B1	-C61	-C66	125.0(15)	C112-C113-C114	123.1(17)		
C66	-C61	-C62	115.9(16)	C113-C114-C115	118.1(16)		
C61	-C62	-C63	124.3(18)	C114-C115-C116	119.2(16)		
C62	-C63	-C64	116.5(19)	C115-C116-C111	124.6(15)		
C63	-C64	-C65	122.6(19)	N9 -C200-C201	172.6(37)		
C64	-C65	-C66	118.3(18)	N10 -C203-C204	174.3(46)		
C65	-C66	-C61	122.0(16)				
B1	-C71	-C72	125.4(15)				
B1	-C71	-C76	118.8(15)				
C76	-C71	-C72	115.8(15)				
C71	-C72	-C73	120.8(17)				
C72	-C73	-C74	118.3(19)				
C73	-C74	-C75	125.2(20)				
C74	-C75	-C76	116.8(18)				
C75	-C76	-C71	123.0(17)				
C81	-B2	-C91	108.8(13)				
C81	-B2	-C101	108.6(13)				
C81	-B2	-C111	112.0(13)				
C91	-B2	-C101	110.0(13)				
C91	-B2	-C111	108.1(13)				
C101-B2	-C111		109.4(13)				
B2	-C81	-C82	124.3(15)				
B2	-C81	-C86	120.7(14)				
C86	-C81	-C82	114.9(15)				
C81	-C82	-C83	121.2(17)				
C82	-C83	-C84	123.4(20)				
C83	-C84	-C85	117.2(19)				
C84	-C85	-C86	119.7(18)				
C85	-C86	-C81	123.3(16)				
B2	-C91	-C92	124.2(15)				
B2	-C91	-C96	122.3(15)				
C96	-C91	-C92	113.5(15)				
C91	-C92	-C93	121.3(16)				
C92	-C93	-C94	123.5(18)				
C93	-C94	-C95	116.0(17)				
C94	-C95	-C96	121.4(17)				
C95	-C96	-C91	124.1(16)				
B2	-C101-C102		119.7(15)				
B2	-C101-C106		121.9(15)				
C106-C101-C102			117.9(16)				
C101-C102-C103			120.1(16)				
C102-C103-C104			118.4(18)				
C103-C104-C105			126.1(21)				
C104-C105-C106			114.5(20)				
C105-C106-C101			122.9(18)				
B2	-C111-C112		126.4(14)				
B2	-C111-C116		120.4(14)				
C116-C111-C112			113.2(14)				

**Table A1.4.** Structure Factors.

The columns contain, in order  $k$ ,  $10F_o$ ,  $10F_c$ , and  $10[(F_o^2 - F_c^2)/\sigma(F_o^2)]$ .

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				3	851	884	-9	10	675	708	-7	10	1412	1355	17
-14	k	4		4	652	709	-14					11	336	432	-15
				5	1351	1253	29	-13	k	9		12	1150	1131	5
0	1621	1556	19	6	369	394	-3					13	588	531	11
1	292	304	-1	7	-84	138	-5	0	-257	248	-30	14	397	391	0
2	335	239	11	8	595	592	0	1	830	784	12				
3	707	668	9	9	538	549	-2	2	777	776	0	-12	k	4	
4	734	710	5	10	520	560	-8	3	905	1023	-34				
								4	618	701	-20	0	660	581	20
-14	k	5		-13	k	4		5	763	780	-4	1	1666	1695	-9
1	660	700	-9	1	-233	130	-17	6	1517	1541	-7	2	1038	997	13
2	766	728	9	2	849	864	-4	7	803	937	-36	3	858	881	-7
3	1134	1070	18	3	1145	1075	21	8	493	359	21	4	335	318	2
4	1320	1296	7	4	1139	1097	12	9	660	665	0	5	357	413	-9
5	357	331	3	5	778	776	0					6	1585	1553	10
6	486	531	-8	6	861	815	12	-13	k	10		7	1418	1402	2
				7	942	839	30					8	957	861	27
-14	k	6		8	307	236	8	1	943	900	12	9	342	355	-1
0	472	500	-5	9	973	1018	-13	2	1693	1693	0	10	220	34	11
1	1010	1021	-3	10	1118	1068	14	3	572	596	-5	11	2019	1963	16
2	1061	1046	4	11	245	163	6	4	-276	178	-24	12	514	497	3
3	641	680	-9					5	323	436	-17	13	501	498	0
4	458	449	1	-13	k	5		6	390	348	5	14	408	430	-3
5	404	244	20					7	1251	1246	1				
6	1452	1387	18	0	555	515	8	8	593	614	-4	-12	k	5	
				1	1538	1545	-2								
-14	k	7		2	1211	1160	15	-13	k	11		1	1578	1586	
1	1195	1226	-9	3	851	925	-21	0	1509	1533	-7	2	1034	1072	-12
2	1746	1622	35	4	389	530	-28	1	954	1029	-22	3	245	140	10
3	677	641	8	5	424	412	2	2	790	761	7	4	1358	1318	13
4	714	621	20	6	1346	1343	0	3	281	239	4	5	537	410	25
5	933	905	7	7	947	977	-9	4	206	196	0	6	421	405	3
6	-218	117	-13	8	1174	1201	-8	5	1395	1372	6	7	1002	973	9
				9	285	209	7					8	1521	1488	10
-14	k	8		10	619	523	19	-12	k	1		9	1968	1978	-3
0	584	609	-5	11	1937	1870	19					10	1007	935	21
1	269	367	-12					1	1129	1119	2	11	139	82	2
2	720	709	2	-13	k	6		2	465	596	-29	12	348	235	14
3	728	745	-4	1	1228	1277	-15	3	529	646	-28	13	1333	1277	16
4	-108	30	-2	2	1058	1066	-2	4	1236	1209	8	14	1510	1475	10
5	1524	1548	-7	3	741	754	-3	5	135	203	-5	15	936	883	13
				4	868	873	-1	6	508	543	-7				
-14	k	9		5	655	651	0	7	436	464	-5	-12	k	6	
1	1334	1354	-5	6	438	293	21	8	977	964	3	0	2687	2757	-20
2	855	836	5	7	1733	1693	12	9	1489	1512	-7	1	400	217	27
3	581	555	5	8	1281	1299	-5	10	231	268	-4	2	425	284	23
				9	1424	1509	-26	11	283	248	3	3	711	796	-25
-13	k	1		10	1230	1165	18	12	159	7	5	4	1055	1027	8
0	610	651	-9	11	307	138	15					5	1585	1634	-16
1	697	700	0					-12	k	2		6	177	26	7
2	833	800	8	0	1808	1848	-12	0	973	979	-1	7	87	158	-4
3	243	308	-7	1	245	47	13	1	-72	185	-10	8	1175	1161	4
4	-237	140	-17	2	546	510	7	2	271	325	-7	9	809	755	14
5	238	566	-56	3	703	769	-17	3	1213	1153	18	10	903	882	6
6	964	988	-6	4	447	483	-6	4	332	368	-5	11	164	64	5
				5	1324	1333	-2	5	1690	1626	-20	12	1268	1320	-16
-13	k	2		6	-95	18	-2	6	928	967	-11	13	1471	1400	21
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2	377	290	17				14	967	1004	-16	5	290	9	58	
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7 k 4			7 k 8			18 931 943			-3	1 426 489			-11		
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4	1114	1156	-15	5	714	777	-15	2	1249	1316	-25	7 977 912			19
5	542	551	-2	6	908	940	-8	3	1462	1454	2				
6	731	845	-39	7	988	1005	-4	4	1829	1885	-20	9 k 0			
7	829	726	32					5	-155	162	-16				
8	599	587	3	8 k 0				6	1215	1277	-23	1	1441	1373	34
9	1869	1878	-3					7	302	249	8	2	2994	2994	0
10	1312	1258	18	0	2081	2177	-43	8	330	311	3	3	1124	1120	2
11	205	2	11	1	852	840	6	9	2291	2242	15	4	244	57	27
12	1588	1587	0	2	669	657	6	10	1126	1097	9	5	1206	1232	-13
13	1223	1184	17	3	101	52	4	11	624	611	3	6	610	686	-35
14	493	432	11	4	333	438	-41	12	1774	1645	41	7	1395	1331	31
15	935	920	4	5	1545	1568	-12	13	1506	1416	28	8	969	1009	-21
16	-234	11	-13	6	321	365	-15	14	866	929	-21	9	474	357	37
17	159	266	-10	7	997	1044	-26	15	738	669	17	10	1703	1663	18
18	891	847	11	8	1847	1877	-14	16	194	24	8	11	233	249	-3
19	1345	1209	38	9	251	120	23	17	302	260	5	12	1384	1471	-43
7 k 5												13 -158 118 -16			
												14 779 809 -12			
												15 1138 1057 33			
0	2935	3078	-40	11	1183	1229	-24	8 k 4			16 567 557 3				
1	106	14	3	12	1085	1187	-54				17 596 456 40				
2	-313	79	-33	13	905	949	-21	0	3092	3145	-14	18 319 261 10			
3	1025	1052	-9	14	1068	1115	-23	1	408	320	16	19 268 103 18			
4	1082	1151	-24	15	357	422	-19	2	96	223	-11				
5	1536	1559	-8	16	595	584	4	3	840	880	-12	9 k 1			
6	615	676	-17	17	1498	1476	9	4	1008	1096	-31				
7	-261	164	-28	18	1048	977	29	5	1183	1104	26	0	949	891	20
8	1676	1610	21	19	424	295	28	6	146	270	-14	1	1653	1696	-16
9	573	579	-1	20	339	250	16	7	599	640	-11	2	804	835	-11
10	1525	1422	33	21	914	916	0	8	1736	1657	25	3	2410	2405	1
11	213	38	10	8 k 1				9	619	659	-10	4	1125	1080	16
12	1855	1845	3	1	1960	2083	-46	10	1107	957	44	5	1182	1200	-6
13	1215	1085	38	2	2743	2831	-27	11	-196	105	-12	6	1590	1640	-18
14	625	552	16	3	751	762	-4	12	1332	1280	16	7	878	858	7
15	-129	316	-26	4	322	230	17	13	757	645	27	8	289	251	6
16	195	270	-7	5	633	641	-2	14	268	409	-20	9	811	799	4
17	692	698	-1	6	1042	1077	-13	15	293	298	0	10	539	419	28
7 k 6												11 2410 2320 26			
												12 1280 1243 12			
												13 1106 1133 -8			
1	1956	1975	-6	7	2182	2163	6	8 k 5			14 575 687 -28				
2	1153	1162	-3	8	1642	1561	29	1	1140	1241	-35	15	315	134	18
3	-193	77	-12	9	1615	1493	43	2	1393	1533	-49	16	1458	1456	0
4	749	709	11	10	2651	2557	27	3	-98	354	-34	17	392	455	-11
5	483	511	-6	11	501	433	16	4	764	901	-41	18	664	548	24
6	150	302	-17	12	1571	1525	16	5	253	262	-1	9 k 2			
7	1112	1065	14	13	374	234	22	6	-162	145	-12				
8	909	855	16	14	493	445	11	7	1188	1131	17	1	1616	1707	-33
9	2098	2147	-14	15	1498	1402	32	8	1311	1290	6	2	1166	1265	-37
10	840	742	25	16	565	492	16	9	1464	1463	0	3	1040	987	18
11	266	89	13	17	1410	1414	-1	10	955	893	17	4	2433	2459	-7
12	-135	95	-6	18	742	677	16	11	588	583	1	5	334	267	11
13	1158	1196	-11	19	596	566	6	12	521	545	-4	6	1291	1384	-34
14	1132	1073	16	20	1823	1659	47	13	1014	917	26	7	258	339	-13
15	774	636	30	8 k 2				14	1270	1236	9	8	718	593	34
7 k 7												9 1965 1941 7			
												10 1181 1225 -15			
0	813	837	-6	0	1015	1012	1	8 k 6			11 337 314 3				
1	1800	1781	6	1	1327	1344	-6				12 1452 1532 -26				
2	1029	956	21	2	1126	1154	-10	0	259	362	-14	13	1452	1400	16
3	851	761	23	3	1772	1856	-31	1	1109	1150	-12	14	1061	1062	0
4	183	228	-4	4	927	962	-13	2	1216	1173	13	15	298	217	9
5	633	639	-1	5	944	987	-16	3	662	698	-8	16	454	355	16
6	1139	1128	3	6	1476	1440	13	4	-286	37	-20	17	565	644	-17
7	1497	1450	14	7	1566	1477	31	5	230	281	-5	9 k 3			
8	793	740	13	8	258	191	8	6	1584	1536	14				
9	-189	116	-11	9	693	657	12	7	1329	1258	21	0	2243	2332	-28
10	707	630	17	10	619	459	40	8	619	669	-11	1	829	889	-19
11	1955	1939	4	11	2961	2929	8	9	262	304	-5	2	375	333	7
12	484	360	19	12	1085	890	61	10	76	8	1	3	251	433	-34
												11 2097 2109 -3			
												8 k 7			

## Iridium - TMB Dimer

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4	752	746	1	4	812	849	-18	5	602	683	-20	10	742	841	-27
5	1700	1752	-18	5	1169	1087	39	6	61	8	0	11	1048	1040	2
6	203	39	10	6	1419	1441	-10	7	1209	1203	1	12	345	336	1
7	686	661	6	7	853	774	35	8	1180	1207	-8				
8	1325	1296	9	8	655	679	-10	9	589	581	1	11	k	2	
9	600	559	9	9	374	343	7	10	1249	1258	-2				
10	1652	1637	4	10	892	921	-13	11	353	304	6	1	762	773	-2
11	632	510	26	11	1721	1648	30	12	869	904	-4	2	1415	1460	-14
12	1096	1032	18	12	1548	1508	17					3	1166	1222	-17
13	1296	1290	1	13	1313	1222	38	10	k	4		4	644	623	5
14	1184	1184	0	14	-186	18	-12					5	763	755	2
15	318	425	-16	15	349	361	-2	0	636	570	14	6	215	330	-14
	9	k	4	16	825	802	8	1	974	964	2	7	1394	1455	-19
				17	194	137	5	2	904	917	-3	8	1468	1462	1
								3	904	985	-23	9	719	797	-20
1	1255	1245	3	10	k	1		4	675	661	3	10	1560	1547	3
2	1576	1568	2					5	-207	77	-11				
3	228	198	3	1	1654	1755	-36	6	1427	1451	-7	11	k	3	
4	833	789	12	2	487	541	-13	7	1039	1072	-9				
5	599	625	-6	3	1092	1091	0	8	111	38	2	0	718	625	21
6	-138	95	-7	4	1927	1958	-10	9	221	135	6	1	1072	1142	-21
7	995	916	22	5	-213	38	-14					2	977	1008	-9
8	1322	1265	17	6	799	804	-1	10	k	5		3	1123	1172	-14
9	884	851	10	7	84	147	-4					4	735	765	-7
10	1427	1449	-7	8	845	850	-1	1	701	786	-21	5	580	666	-19
11	270	326	-7	9	2133	2097	10	2	469	479	1	6	1299	1266	9
12	766	784	-4	10	588	514	16	3	1253	1245	2	7	584	574	2
13	444	327	17	11	319	244	9	4	803	748	13				
	9	k	5	12	1360	1290	21					12	k	0	
				13	1394	1392	0	11	k	0		0	1177	1129	20
0	200	65	8	15	394	99	29	1	1288	1322	-16	1	964	949	6
1	1161	1135	8	16	999	949	14	2	216	261	-7	2	748	759	-3
2	767	765	0					3	659	719	-25	3	-229	39	-18
3	821	894	-20	10	k	2		4	1574	1592	-8	4	-262	21	-24
4	256	347	-12					5	-161	16	-9	5	1050	1052	0
5	103	268	-14	0	1963	2087	-41	6	542	409	38	6	744	636	36
6	1328	1311	5	1	1051	1031	6	7	306	342	-7	7	875	895	-7
7	807	767	10	2	831	887	-19	8	1020	1017	1	8	1060	987	28
8	429	571	-28	3	-123	124	-8	9	1834	1750	33	9	258	216	6
9	-266	53	-17	4	328	389	-10	10	162	6	8	10	538	496	11
10	-181	41	-7	5	1580	1591	-3	11	306	209	15				
	9	k	6	6	678	656	5	12	376	333	9	12	k	1	
				7	986	840	43	13	988	942	18				
1	743	775	-8	8	1867	1868	0	14	1230	1198	12	1	275	339	-8
2	762	720	10	9	496	468	5					2	1449	1429	6
3	1436	1433	0	10	1722	1694	8	11	k	1		3	1140	1105	10
4	651	610	8	11	782	522	57					4	-102	303	-23
5	372	416	-6	12	1042	1049	-2	0	1718	1750	-10	5	682	692	-2
6	680	669	2	13	1285	1339	-16	1	1017	1057	-12	6	321	370	-7
	10	k	0	14	1127	1256	-39	2	526	626	-24	7	1216	1103	31
								3	-170	232	-21				
								4	57	137	-3	12	k	2	
0	1627	1660	-15	1	1013	1092	-25	5	1602	1582	6	0	870	803	17
1	1249	1288	-19	2	1711	1683	8	6	714	660	13	1	729	733	0
2	686	686	0	3	258	363	-15	7	1051	987	18	2	910	900	2
3	1735	1735	0	4	786	741	13	8	1312	1403	-29				
								9	450	345	17				



## **Appendix 2**

C PROGRAM DIFFAN

```

C
  OPTIONS/G_FLOATING
  PARAMETER (LW=1000,LIW=LW/8+2)
  IMPLICIT DOUBLE PRECISION (A-H,O-Z)
  INTEGER IW(LIW)
  DIMENSION W(LW)
  EXTERNAL F,D01AJF
  COMMON/CONST/AK,C,KOUNT
  READ(5,101) IZA,IZB,X1,SI,T,RA,RB,OBSK,DA,DB,N
101 FORMAT(2I4,F4.2,E10.4,F6.2,2F4.2,E9.3,2E10.4,I5)
  TK = 273.15 + T
  AB = RA + RB
  IF (X1.EQ.0) X1 = AB
  DIEL = EXP(3.69031 - 0.00430493*T)
  ETA = EXP(885.636/TK - 4.03662)
  AID = SI/(DIEL*TK)
  AK = 50.2900016*SQRT(AID)
  TB = IZA*IZB*167102.4/(DIEL*TK)
  BB = TB/X1
  CHK = ABS(IZA*IZB)
  IF (CHK.GT.0) GO TO 1
  RESULT = 1.0/X1
  GO TO 2
1 DC = EXP(AK*X1)/(1.0 + AK*X1)
  C = TB*DC
  EPSABS = 0.0
  EPSREL = 1.0E-5
  A = X1
  B = N*X1
  KOUNT = 0
  IFAIL = -1
  CALL D01AJF(F,A,B,EPSABS,EPSREL,RESULT,ABSERR,W,LW,IW,LIW,IFAIL)
  IF (IFAIL.NE.0) THEN
    WRITE(6,1001) IFAIL
1001  FORMAT(' D01AJF FAILED. IFAIL = ',I2)
    STOP
  END IF
2 FF = 1.0/(RESULT*X1)
  PDC = 7.3246393*TK/(ETA*10.0**8)
  DAA = PDC/RA
  IF (DAA.GT.0) GO TO 3
  DA = DAA
3 DBB = PDC/RB
  IF (DBB.GT.0) GO TO 4
  DB = DBB
4 DK = 75.675251*AB*FF*(DA + DB)*10.0**12
  DHF = EXP(BB/(1.0 + X1*AK))
  BDK = 3.0*(DA + DB)*FF*DHF*10.0**16/AB/AB
  EQK00 = 0.0025225083*AB*AB*AB
  EQK = EQK00/DHF
  EQK0 = EQK00/EXP(BB)
  ACTK = BDK/(DK/OBSK - 1.0)
  ACT1 = 1.0/(1.0/OBSK - 1.0/DK)
  WRITE(6,103) IZA,IZB,T,SI,X1,RA,RB,OBSK,DA,DB
103 FORMAT(5X,'CHARGE ON ION A = ',I4,'/5X'CHARGE ON ION B = ',I4,
  *'/5X','TEMPERATURE (DEG. CELCIUS) = ',F6.2,'/5X','IONIC STRENGTH
  *(MOLES/LITER) = ',E10.4,'/5X','A-B SEPERATION DISTANE (A) = ',F5.2,
  *'/5X','RADIUS OF ION A (A) = ',F5.2,'/5X','RADIUS OF ION B (A) = ',
  *F5.2,'/5X','OBSERVED RATE CONSTANT (1/Msec) = ',E9.3,'/5X','DIFFUSION
  *CONSTANT FOR ION A (CM2/SEC) = ',E10.4,'/5X','DIFFUSION CONSTANT
  *FOR ION B (CM2/SEC) = ',E10.4)
  WRITE(6,104)
104 FORMAT(/3X,'THE CALCULATED PARAMETERS:')

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WRITE(6,105)DIEL,ETA,BB,AK,FF,DHF,DAA,DBB
105 FORMAT(5X,'DIELECTRIC CONSTANT OF CH3CN = ',F8.3,/5X,'VISCOSITY OF
*CH3CN = ',F7.4,/5X,'DEBYE-HUCKEL PARAMETER, B , = ',E12.6,/5X,'DEBYE-
*HUCKEL PARAMETER, KAPPA (A-1), = ',E12.5,/5X,'DEBYE CHARGE-IONIC ATMOS
*PHERE PARAMETER, F, = ',E12.6,/5X,'EIGEN-FUOSS k-d IONIC
*PARAMETER = ',E12.6,/5X,'THE CALCULATED DIFFUSION COEFFICIENT OF
*ION A (CM2/SEC) = ',E10.4,/5X,'THE CALCULATED DIFFUSION COEFFICIENT
*OF ION B (CM2/SEC) = ',E10.4)
WRITE(6,106)
106 FORMAT(/3X,'THE CALCULATED RESULTS ARE:')
WRITE(6,107)DK,BDK,EQK,EQK0,ACTK,ACT1
107 FORMAT(5X,'DIFFUSION CONTROLLED RATE, (1/Msec) = ',E10.4,/5X,'DISSOCIA
*TION RATE (1/SEC) = ',E10.4,/5X,'EQUILIBRIUM CONSTANT (MOLES/LITER)
*= ',E10.4,/5X,'EQUILIBRIUM CONSTANT AT ZERO IONIC STRENGTH (MOLES/LITER)
*= ',E10.4,/5X,'ACTIVATED 1ST ORDER RATE CONSTANT (1/SEC) = ',E10.4,
*/5X,'ACTIVATED 2ND ORDER RATE CONSTANT (1/Msec) = ',E10.4)
WRITE(6,108)RESULT,KOUNT,IFAIL,ABSERR,IW(1)
108 FORMAT(/5X,'RESULT = ',F10.4,5X,'KOUNT = ',I4,5X,'IFIAL = ',I4,
*/5X,'ABSERR = ',E10.4,5X,'IW(1) = ',I4)
END
REAL*8 FUNCTION F(X)
IMPLICIT DOUBLE PRECISION (A-H,O-Z)
COMMON/CONST/AK,C,KOUNT
INTEGER KOUNT
KOUNT = KOUNT + 1
F = EXP(C/(X*EXP(AK*X)))/(X**2)
RETURN
END

```

## Results for 4-cyano-N-methylpyridinium.

CHARGE ON ION A = 1  
 CHARGE ON ION B = 2  
 TEMPERATURE (DEG. CELCIUS) = 22.00  
 IONIC STRENGTH (MOLES/LITER) = 0.1000E+00  
 A-B SEPERATION DISTANE (A) = 11.19  
 RADIUS OF ION A (A) = 3.34  
 RADIUS OF ION B (A) = 7.85  
 OBSERVED RATE CONSTANT (1/Msec) = 0.450E+09  
 DIFFUSION CONSTANT FOR ION A (CM2/SEC) = 0.1824E-04  
 DIFFUSION CONSTANT FOR ION B (CM2/SEC) = 0.7760E-05  
 UPPER BOUND OF INTEGRATION,N\*X1, N = 100

## THE CALCULATED PARAMETERS:

DIELETRIC CONSTANT OF CH3CN = 36.438  
 VISCOSITY OF CH3CN = 0.3549  
 DEBYE-HUCKEL PARAMETER, B, = 0.277709E+01  
 DEBYE-HUCKEL PARAMETER, KAPPA (A-1), = 0.15335E+00  
 DEBYE CHARGE-IONIC ATMOSPHERE PARAMETER, F, = 0.750328E+00  
 EIGEN-FUOSS k-d IONIC PARAMETER = 0.278012E+01  
 THE CALCULATED DIFFUSION COEFFICIENT OF ION A (CM2/SEC) = 0.1824E-04  
 THE CALCULATED DIFFUSION COEFFICIENT OF ION B (CM2/SEC) = 0.7760E-05

## THE CALCULATED RESULTS ARE:

DIFFUSION CONTROLLED RATE,(1/Msec) = 0.1652E+11  
 DISSOCIATION RATE (1/SEC) = 0.1299E+11  
 EQUILIBRIUM CONSTANT (MOLES/LITER) = 0.1271E+01  
 EQUILIBRIUM CONSTANT AT ZERO IONIC STRENGTH (MOLES/LITER) = 0.2199E+00  
 ACTIVATED 1ST ORDER RATE CONSTANT (1/SEC) = 0.3639E+09  
 ACTIVATED 2ND ORDER RATE CONSTANT (1/Msec) = 0.4626E+09

RESULT = .119101994      KOUNT = 273      IFIAL = 0  
 ABSERR = 0.5720E-09      IW(1) = 28

Results for 4-carbomethoxy-N-methylpyridinium.

CHARGE ON ION A = 1  
 CHARGE ON ION B = 2  
 TEMPERATURE (DEG. CELCIUS) = 22.00  
 IONIC STRENGTH (MOLES/LITER) = 0.1000E+00  
 A-B SEPERATION DISTANE (A) = 11.72  
 RADIUS OF ION A (A) = 3.87  
 RADIUS OF ION B (A) = 7.85  
 OBSERVED RATE CONSTANT (1/Msec) = 0.330E+09  
 DIFFUSION CONSTANT FOR ION A (CM2/SEC) = 0.1574E-04  
 DIFFUSION CONSTANT FOR ION B (CM2/SEC) = 0.7760E-05  
 UPPER BOUND OF INTEGRATION,N\*X1, N = 100

THE CALCULATED PARAMETERS:

DIELECTRIC CONSTANT OF CH3CN = 36.438  
 VISCOSITY OF CH3CN = 0.3549  
 DEBYE-HUCKEL PARAMETER, B, = 0.265150E+01  
 DEBYE-HUCKEL PARAMETER, KAPPA (A-1), = 0.15335E+00  
 DEBYE CHARGE-IONIC ATMOSPHERE PARAMETER, F, = 0.773665E+00  
 EIGEN-FUOSS k-d IONIC PARAMETER = 0.258026E+01  
 THE CALCULATED DIFFUSION COEFFICIENT OF ION A (CM2/SEC) = 0.1574E-04  
 THE CALCULATED DIFFUSION COEFFICIENT OF ION B (CM2/SEC) = 0.7760E-05

THE CALCULATED RESULTS ARE:

DIFFUSION CONTROLLED RATE, (1/Msec) = 0.1613E+11  
 DISSOCIATION RATE (1/SEC) = 0.1025E+11  
 EQUILIBRIUM CONSTANT (MOLES/LITER) = 0.1574E+01  
 EQUILIBRIUM CONSTANT AT ZERO IONIC STRENGTH (MOLES/LITER) = 0.2865E+00  
 ACTIVATED 1ST ORDER RATE CONSTANT (1/SEC) = 0.2141E+09  
 ACTIVATED 2ND ORDER RATE CONSTANT (1/Msec) = 0.3369E+09

RESULT = .110285706      KOUNT = 273      IFIAL = 0  
 ABSERR = 0.4390E-09      IW(1) = 28

## Results for 3-carbomethoxy-N-benzylpyridinium.

CHARGE ON ION A = 1  
 CHARGE ON ION B = 2  
 TEMPERATURE (DEG. CELCIUS) = 22.00  
 IONIC STRENGTH (MOLES/LITER) = 0.1000E+00  
 A-B SEPERATION DISTANE (A) = 12.49  
 RADIUS OF ION A (A) = 4.64  
 RADIUS OF ION B (A) = 7.85  
 OBSERVED RATE CONSTANT (1/Msec) = 0.280E+08  
 DIFFUSION CONSTANT FOR ION A (CM2/SEC) = 0.1313E-04  
 DIFFUSION CONSTANT FOR ION B (CM2/SEC) = 0.7760E-05  
 UPPER BOUND OF INTEGRATION, N\*X1, N = 100

## THE CALCULATED PARAMETERS:

DIELETRIC CONSTANT OF CH3CN = 36.438  
 VISCOSITY OF CH3CN = 0.3549  
 DEBYE-HUCKEL PARAMETER, B, = 0.248804E+01  
 DEBYE-HUCKEL PARAMETER, KAPPA (A-1), = 0.15335E+00  
 DEBYE CHARGE-IONIC ATMOSPHERE PARAMETER, F, = 0.802966E+00  
 EIGEN-FUOSS k-d IONIC PARAMETER = 0.234788E+01  
 THE CALCULATED DIFFUSION COEFFICIENT OF ION A (CM2/SEC) = 0.1313E-04  
 THE CALCULATED DIFFUSION COEFFICIENT OF ION B (CM2/SEC) = 0.7760E-05

## THE CALCULATED RESULTS ARE:

DIFFUSION CONTROLLED RATE, (1/Msec) = 0.1585E+11  
 DISSOCIATION RATE (1/SEC) = 0.7573E+10  
 EQUILIBRIUM CONSTANT (MOLES/LITER) = 0.2094E+01  
 EQUILIBRIUM CONSTANT AT ZERO IONIC STRENGTH (MOLES/LITER) = 0.4083E+00  
 ACTIVATED 1ST ORDER RATE CONSTANT (1/SEC) = 0.1340E+08  
 ACTIVATED 2ND ORDER RATE CONSTANT (1/Msec) = 0.2805E+08

RESULT = .099710350      KOUNT = 273      IFIAL = 0  
 ABSERR = 0.3123E-09      IW(1) = 28

### **Appendix 3**

Table A3.1.  $U_{ij}$ 's, Ir atoms.

Atom	$U_{11}$	$U_{22}$	$U_{33}$	$U_{12}$	$U_{13}$	$U_{23}$
IR1	275(9)	371(9)	250(8)	-12(9)	42(7)	-19(8)
IR2	346(10)	366(9)	251(8)	-42(10)	21(7)	6(8)

The form of the displacement factor is:

$$\exp -2\pi^2(U_{11}h^2a^{*2} + U_{22}k^2b^{*2} + U_{33}\ell^2c^{*2} + 2U_{12}hka^*b^* + 2U_{13}h\ell a^*c^* + 2U_{23}k\ell b^*c^*)$$



Table A3.2. Hydrogen coordinates.

Atom	$x, y \text{ and } z \times 10^4$			
	$x$	$y$	$z$	$B$
H1	-1095	2527	2873	5.0
H2	3266	2577	4413	5.0
H3A	2411	1074	2073	6.0
H3B	2135	1180	2620	6.0
H3C	1087	1247	2213	6.0
H4A	2651	1741	1469	6.0
H4B	1319	1897	1621	6.0
H4C	2489	2195	1698	6.0
H5A	4288	1554	2262	6.0
H5B	3754	1694	2766	6.0
H6A	3704	2418	2290	6.0
H6B	4920	2209	2086	6.0
H7A	6701	1994	2801	6.0
H7B	6416	2133	3338	6.0
H7C	5600	1779	3079	6.0
H8A	6370	2766	2480	6.0
H8B	5102	2994	2596	6.0
H8C	6119	2921	3013	6.0
H13A	3152	3707	2011	6.0
H13B	3205	3301	2346	6.0
H13C	2139	3348	1942	6.0
H14A	1425	4277	2167	6.0
H14B	386	3928	2085	6.0
H14C	478	4179	2581	6.0
H15A	3296	4092	2841	6.0
H15B	3103	3669	3136	6.0
H16A	1106	4087	3287	6.0
H16B	2149	4440	3297	6.0
H17A	1553	4563	4208	6.0
H17B	635	4172	4202	6.0
H17C	1778	4185	4575	6.0
H18A	3871	4435	4025	6.0
H18B	3845	4045	4384	6.0
H18C	4247	3980	3842	6.0
H23A	-2994	3941	4336	6.0
H23B	-1635	3853	4154	6.0
H23C	-2798	3785	3799	6.0
H24A	-4483	3244	4470	6.0
H24B	-4246	3144	3918	6.0
H24C	-3906	2804	4322	6.0
H25A	-2383	3376	5026	6.0
H25B	-1112	3289	4768	6.0
H26A	-2160	2564	4757	6.0
H26B	-2546	2753	5262	6.0
H27A	-1507	2088	5654	6.0

Atom	$x$	$y$	$z$	$U_{eq}$ or $B$
H27B	-1035	1934	5145	6.0
H27C	-63	2049	5566	6.0
H28A	-837	2879	5921	6.0
H28B	521	2750	5769	6.0
H28C	-144	3146	5525	6.0
H33A	-2779	1088	4370	6.0
H33B	-2179	1531	4504	6.0
H33C	-3108	1491	4052	6.0
H34A	-1937	708	3573	6.0
H34B	-2355	1122	3290	6.0
H34C	-979	951	3248	6.0
H35A	-529	838	4420	6.0
H35B	36	1298	4461	6.0
H36A	823	1047	3609	6.0
H36B	935	628	3921	6.0
H37A	2337	696	4782	6.0
H37B	3116	1116	4857	6.0
H37C	1652	1124	4902	6.0
H38A	3349	625	3951	6.0
H38B	3330	1016	3594	6.0
H38C	4075	1043	4095	6.0
H112	7761	3434	1331	4.0
H113	8737	3251	2072	5.0
H114	8060	3590	2811	6.0
H115	6445	4055	2791	5.0
H116	5295	4211	2042	4.0
H122	8269	4182	1075	4.0
H123	9585	4570	534	5.0
H124	8589	4930	-139	6.0
H125	6451	4855	-290	5.0
H126	5271	4456	230	4.0
H132	5279	4820	1319	4.0
H133	3274	5205	1477	5.0
H134	1425	4896	1286	6.0
H135	1276	4202	1022	5.0
H136	3099	3790	875	4.0
H142	4567	3189	1378	4.0
H143	3976	2553	975	5.0
H144	4267	2468	152	6.0
H145	5186	2992	-278	5.0
H146	5926	3627	80	4.0
H212	8123	1149	1744	4.0
H213	9252	1736	2111	5.0
H214	8497	2429	1965	6.0
H215	6714	2522	1453	5.0

Atom	$x$	$y$	$z$	$U_{eq}$ or $B$
H216	5673	1934	1074	4.0
H222	5927	1133	2187	4.0
H223	5564	629	2824	5.0
H224	5155	-70	2628	6.0
H225	5243	-316	1826	5.0
H226	5864	155	1222	4.0
H232	8726	1031	839	4.0
H233	9826	610	296	5.0
H234	8807	169	-248	6.0
H235	6614	116	-245	5.0
H236	5405	440	381	4.0
H242	5423	1546	380	4.0
H243	3404	1666	-52	5.0
H244	1530	1440	275	6.0
H245	1521	941	865	5.0
H246	3463	761	1313	4.0

**Table A3.3.** Structure Factors.

The columns contain, in order  $k$ ,  $10F_o$ ,  $10F_c$ , and  $10[(F_o^2 - F_c^2)/\sigma(F_o^2)]$ .

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Page 1

-7 k 1	0 1680-1795 -7	3 -534 -213 -13	6 723 545 8
1 586 -125 11	1 2628 2709 -10	4 518 453 2	7 688 61 19
2 558 52 10	2 1876 1869 0	5 -609 40 -15	8 484 -221 7
3 -323 239 -5	3 2377-2352 3	6 1000 -783 15	-6 k 11
4 362 -255 2	-6 k 1	7 390 -381 0	1 -151 40 0
5 549 -417 4	1 166 250 -1	8 941 890 3	2 480 -147 7
6 -529 -160 -11	2 -146 192 -2	9 491 387 3	3 -151 -13 0
7 -245 148 -3	3 462 -288 5	10 509 -651 -6	4 606 -138 13
8 465 23 8	4 797 82 26	11 711 -309 17	5 -253 126 -3
9 728 -642 4	5 569 -202 12	12 1066 1044 1	6 -72 376 -5
-7 k 2	6 903 855 3	13 645 598 2	7 433 57 7
0 2035-1881 12	7 -79 330 -5	-6 k 6	-6 k 12
1 109 -368 -4	8 870 -929 -4	0 3447 3606 -18	0 -455 -34 -4
2 2017 2050 -3	9 600 -621 -1	1 -90 -802 -23	1 1456 1448 0
3 -559 145 -11	10 638 711 -4	2 3118-2972 22	2 -209 218 -3
4 1949-1954 0	11 -272 350 -8	3 596 487 4	3 1354-1235 10
5 271 -310 0	12 828 -999 -12	4 2624 2687 -10	4 739 -83 20
6 1567 1567 0	13 647 -57 17	5 654 -358 12	-5 k 1
7 127 133 0	14 1251 1265 -1	6 2188-2064 18	1 531 -416 4
8 1651-1512 14	-6 k 2	7 -215 460 -10	2 -398 127 -8
9 -288 82 -3	0 3242-3268 -3	8 2750 2835 -14	3 1319-1310 1
-7 k 3	1 1360 1396 -3	9 -82 -500 -10	4 628 -654 -1
1 714 -45 17	2 3386 3344 7	10 2408-2283 18	5 373 223 4
2 599 332 8	3 1415-1344 7	11 -509 310 -15	6 242 163 1
3 514 227 7	4 2616-2658 -6	12 2131 2371 -36	7 329 -481 -6
4 -489 7 -8	5 1716 1703 1	-6 k 7	8 765 -692 5
5 258 -240 0	6 2997 3001 0	1 425 66 6	9 605 345 12
6 363 -487 -3	7 1644-1648 0	2 -802 130 -24	10 1004 1058 -5
7 -58 433 -7	8 3297-3157 24	3 567 -522 1	11 342 -572 -10
8 912 783 8	9 1738 1706 4	4 -461 140 -9	12 1415-1435 -2
-7 k 4	10 2651 2549 16	5 432 437 0	13 721 607 7
0 1193 1274 -4	11 966-1270 -27	6 483 303 6	14 1310 1139 18
1 3027-2990 5	12 2226-2310 -12	7 791 -294 22	15 848 -866 -1
2 845-1005 -9	13 1193 1306 -11	8 718 -228 19	16 1393-1272 13
3 2496 2540 -5	14 2244 2267 -3	9 491 542 -2	17 683 718 -2
4 1087 925 10	-6 k 3	10 695 649 2	-5 k 2
5 2482-2375 14	1 -511 -169 -11	11 731 -693 2	0 1474-1352 10
6 909 -676 13	2 761 -453 14	-6 k 8	1 3690 3768 -14
7 2446 2357 12	3 583 657 -3	0 1014 619 14	2 1524 1319 24
8 579 557 0	4 339 251 2	1 2346 2174 22	3 3631-3634 0
-7 k 5	5 751 -607 8	2 483 -608 -5	4 1578-1587 -1
1 546 117 9	6 -234 -283 -5	3 1572-1486 9	5 3906 3829 14
2 -405 -46 -5	7 186 251 -1	4 387 26 5	6 1278 1202 9
3 -358 -455 -11	8 775 547 12	5 1614 1697 -10	7 3071-3070 0
4 651 383 9	9 894 -626 17	6 267 252 0	8 1840-1816 3
5 232 297 -1	10 -237 -263 -5	7 1742-1644 12	9 3386 3324 11
6 -535 -198 -11	11 1096 1051 3	8 454 19 8	10 1282 1223 6
7 523 -560 -1	12 -476 486 -19	9 1755 1768 -1	11 2632-2708 -13
-7 k 6	13 1238-1133 9	10 277 127 2	12 1454-1202 29
0 3070 3021 5	-6 k 4	11 1465-1423 4	13 2466 2384 14
1 1472 1577 -9	0 1445-1477 -2	-6 k 9	14 702 816 -7
2 2723-2820 -13	1 3700-3690 1	1 -339 -58 -3	15 2267-2317 -7
3 1259-1135 9	2 1625 1484 15	2 -237 151 -2	16 689 -914 -14
4 2882 2805 11	3 2810 2822 -1	3 260 201 1	17 2506 2466 6
5 1607 1489 11	4 1098-1250 -14	4 -185 -47 -1	-5 k 3
6 2782-2762 2	5 3125-3108 2	5 -343 -268 -7	1 310 114 3
-7 k 7	6 1403 1300 11	6 217 167 0	2 870 620 18
1 -345 110 -4	7 3269 3257 2	7 -469 99 -9	3 384 114 6
2 -306 151 -3	8 845-1039 -15	8 -693 -168 -21	4 642 524 6
3 587 240 9	9 3131-2978 25	9 289 -100 3	5 588 -412 9
4 -202 -204 -2	10 700 738 -2	10 475 284 5	6 -668 255 -26
5 622 591 1	11 3133 3129 0	-6 k 10	7 1135 1029 11
-7 k 8	12 1015-1089 -6	0 -856 -221 -17	8 -271 -219 -6
1 278 150 2	13 2609-2678 -10	1 -548 267 -12	9 377 -552 -8
2 589 -228 11	-6 k 5	2 578 193 10	10 364 168 5
	1 278 150 2	3 -538 112 -10	11 1177 1053 12
	2 589 -228 11	4 -347 28 -4	12 828 -581 16
		5 386 -3 5	13 1274-1267 0

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14	570	649	-4	15	94	-137	0	7	2599-2552	7					
15	1181	1189	0					8	1284	1252	3	-4	k	3	
16	-36	-473	-9		-5	k	8	9	2093	2173	-11				
								10	1538-1518		2	1	-316	-199	-6
	-5	k	4	0	708	-894	-7					2	-541	80	-15
				1	-340	-522	-15	-5	k	13		3	-263	229	-6
0	2415-2352	7		2	781	592	11					4	185	3	2
1	1024-1036	-1		3	362	323	1	1	-572	37	-12	5	547	474	4
2	3015-2986	5		4	492	-410	3	2	-525	77	-11	6	357	-115	7
3	1786-1760	3		5	-141	1	-1	3	-205	202	-3	7	530	-418	6
4	2508-2447	10		6	695	367	16	4	475	-15	9	8	214	-494	-12
5	1718-1632	12		7	273	-102	3	5	-516	-175	-12	9	612	-392	13
6	2141-1926	35		8	487	-525	-1	6	455	554	-4	10	-487	310	-19
7	1398-1289	13		9	-211	-112	-2	7	-420	135	-8	11	-577	-126	-20
8	1720-1710	1		10	414	375	1	8	703	-832	-8	12	668	-450	13
9	1444-1384	7		11	315	-70	4	9	340	-192	3	13	-66	401	-8
10	1799-1884	-13		12	520	-845	-19					14	1244	1347	-12
11	1236-1299	-7		13	-403	-91	-7	-5	k	14		15	670	-30	22
12	2028-2064	-5		14	501	115	10					16	968-1064	-8	
13	1011-920	7						0	3805-3845	-4		17	520	-416	4
14	1998-2008	-1		-5	k	9		1	1561	1542	1	18	474	538	-2
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16	1348-1282	6		2	481	25	9	3	1071	-931	11				
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	-5	k	5	4	-386	-152	-8	5	1028	895	10	0	1254	-973	19
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2	350	-480	-4	6	-425	-259	-11	7	1268-1279	-1		2	471	581	-6
3	567	61	15	7	164	38	1					3	498	-428	3
4	449	234	7	8	107	304	-3	-5	k	15		4	1143-1091	6	
5	414	-332	3	9	531	765	-13	1	-545	-208	-12	5	325	-398	-3
6	435	360	3	10	651	-333	14	2	-314	-2	-4	6	415	518	-5
7	128	-47	0	11	718	-147	21	3	723	601	6	7	1149	1239	-12
8	391	334	2	12	616	446	7					8	-398	-170	-11
9	-415	-208	-11	13	567	479	3	-4	k	1		9	236	264	0
10	612	-526	4	14	562	-524	1					10	859	913	-5
11	-423	-122	-9					1	77	6	0	11	604	-669	-4
12	600	357	11	-5	k	10		2	472	273	7	12	363	-452	-3
13	99	-179	-1	0	1901	2004	-8	3	919	855	7	13	421	-366	2
14	1188	-941	22	1	1175	-1097	6	4	725	-797	-7	14	487	120	11
15	703	644	3	2	2213-2175	5		5	667	-567	8	15	390	399	0
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4	660	661	0	10	1781-1749	4		13	1163	1095	7	3	337	428	-3
5	610	-356	12	11	1400	1231	17	14	301	126	3	4	588	-631	-3
6	1375-1315	7		12	1326	1445	-13	15	1191-1249	-6		5	225	718	-27
7	984	890	8	13	1146-1226	-7		16	456	-257	6	6	-296	27	-5
8	938	939	0					17	1490	1451	4	7	365	-353	0
9	816	-697	8	-5	k	11		18	423	517	-4	8	691	-662	2
10	1140	-966	17					19	1483-1470	1		9	-233	346	-10
11	1064	996	6	1	368	-307	1					10	670	462	13
12	713	985	-20	2	-174	500	-11	-4	k	2		11	-347	152	-8
13	703	-513	10	3	373	-304	2	0	1459	1504	-3	12	775	-726	3
14	421	-611	-8	4	125	242	-1	1	1948	1911	5	13	-588	-81	-18
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				6	254	229	0	3	2191-2235	-8		15	233	251	0
	-5	k	7	7	680	-698	-1	4	436	-316	5	16	544	-468	3
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3	243	-358	-3	10	-98	396	-7	7	2392-2371	4					
4	321	-276	1	11	966	-881	6	8	907	700	13	-4	k	6	
5	124	76	0	12	-192	-232	-3	9	2196	2190	1	0	816	-680	6
6	248	132	2					10	-162	299	-6	1	2130	2089	6
7	-551	-247	-18	-5	k	12		11	1804-1827	-3		2	631	524	6
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13	-504	-116	-12	5	2665	2687	-3	17	1595	1504	11	8	167	-466	-10
14	-413	-136	-8	6	1494-1420	8		18	562	46	14	9	1388	1370	2
								19	1477-1445	3					

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10	-256	0	-3	6	480	-634	-9	2	-521	-213	-14	6	1198	-1121	10	
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13	1104	924	18	9	3782	-3830	-8	5	300	-323	0	9	166	92	1	
14	-397	-64	-7	10	797	-649	10	6	-449	129	-9	10	676	-607	5	
15	1329	-1329	0	11	3320	3296	4	7	375	317	1	11	-309	143	-7	
16	330	-192	3	12	651	745	-6	8	-660	-230	-21	12	109	617	-21	
17	884	758	8	13	3040	-3046	-1	9	615	-708	-5	13	-556	186	-20	
18	456	-608	-7	14	1092	-988	9					14	-393	-163	-10	
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	-4	k	7						-4	k	16		16	757	573	13
					-4	k	11		0	1054	-1300	-12	17	409	-87	8
1	663	-501	9	1	290	-321	0	1	900	-924	-1	18	706	-418	15	
2	-622	-186	-22	2	633	718	-5	2	1629	1665	-4	19	360	-26	6	
3	589	-388	11	3	385	-464	-3	3	522	626	-5	20	644	488	7	
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5	432	-373	2	5	771	591	12	5	1076	-1157	-7		-3	k	4	
6	479	577	-5	6	532	-549	0	6	1072	1063	0	0	884	1005	-7	
7	375	514	-6	7	1134	-1177	-4					1	1620	-1558	9	
8	-453	150	-13	8	-438	106	-10		-4	k	17	2	1399	-1183	29	
9	714	-597	8	9	1210	1292	-9	1	395	-123	5	3	1716	1595	20	
10	384	353	1	10	594	-40	17	2	-543	122	-12	4	1609	1516	15	
11	776	723	4	11	1341	-1378	-4					5	2153	-2173	-4	
12	368	-641	-14	12	210	373	-4		-3	k	1	6	1325	-1101	22	
13	499	-438	2	13	1418	1344	8	1	1061	-956	10	7	1591	1527	11	
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## Di-iridium complex.

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Di-iridium complex.										Page	15				
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## Di-iridium complex.

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